

Alexander

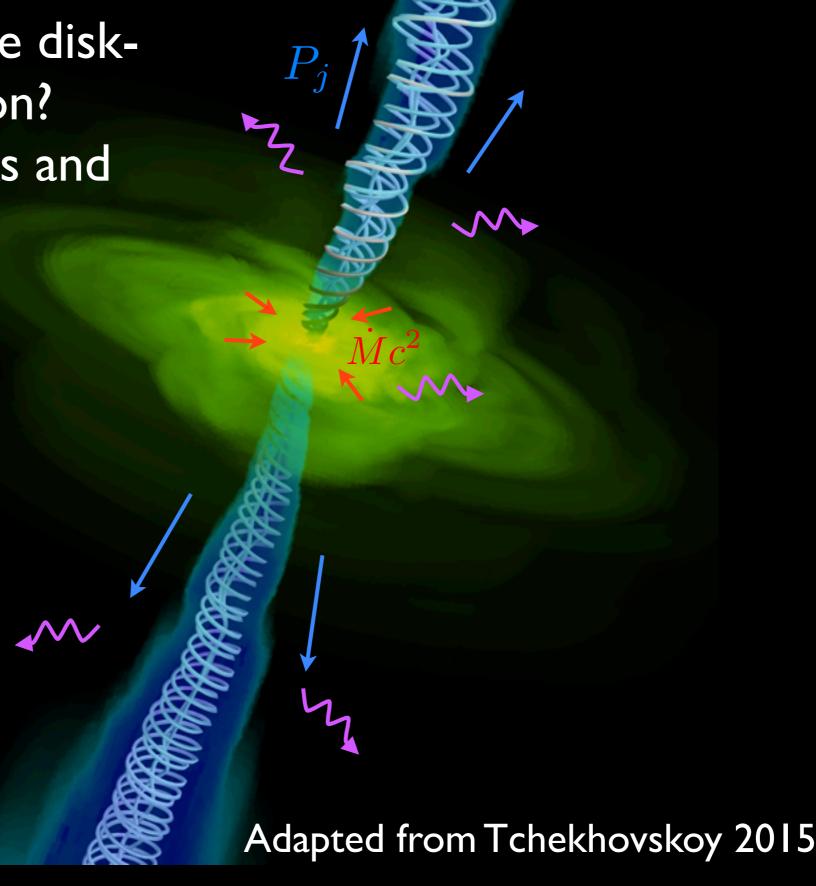
(Sasha) Tchekhovskoy

Einstein Fellow UC Berkeley

Interpreting Black Hole Observations What sets the diskjet connection? Adapted from Tchekhovskoy 2015

Interpreting Black Hole Observations What sets the diskjet connection? How do disks and

jets emit?

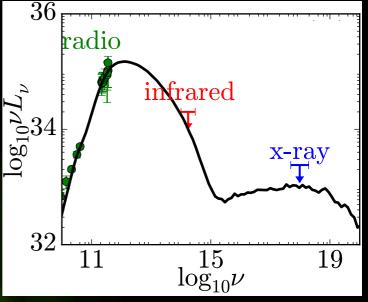


What sets the diskjet connection?

- How do disks and jets emit?
- What can we learn from
  - spectra

Simulated ALMA, VLA, Chandra, NuSTAR spectra

Observations





What sets the diskjet connection?

How do disks and jets emit?

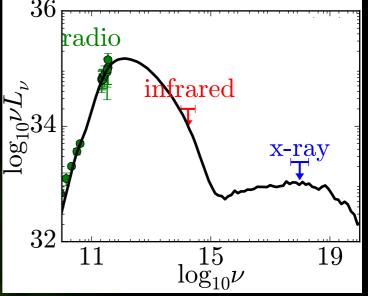
What can we learn from

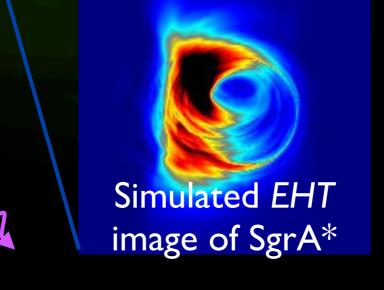
spectra

imageson small and

Simulated ALMA, VLA, Chandra, NuSTAR spectra

Observations





What sets the diskjet connection?

How do disks and jets emit?

 What can we learn from

spectra

images on small and large scales? Simulated ALMA, VLA, Chandra, NuSTAR spectra

11

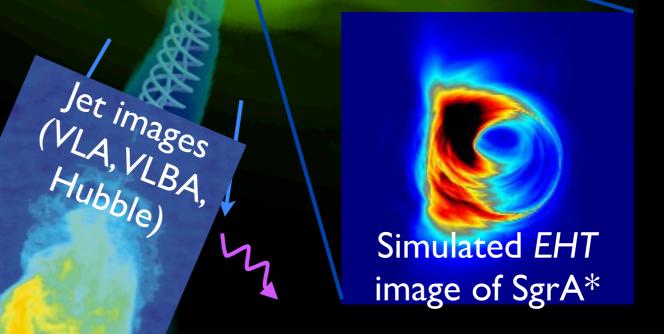
Observations

x-ray

19

15

 $\log_{10} \nu$ 



Jet images (VLA, VLBA,

What sets the diskjet connection?

How do disks and jets emit?

What can we learn from

spectra

images on small and large scales? What is a
Healthy Jet Diet?

• Disk?

• Disk size/thickness/
collimation?

• Net vertical

• Net vertical magentic flux/ dynamo?

• Ambient medium?

Observations

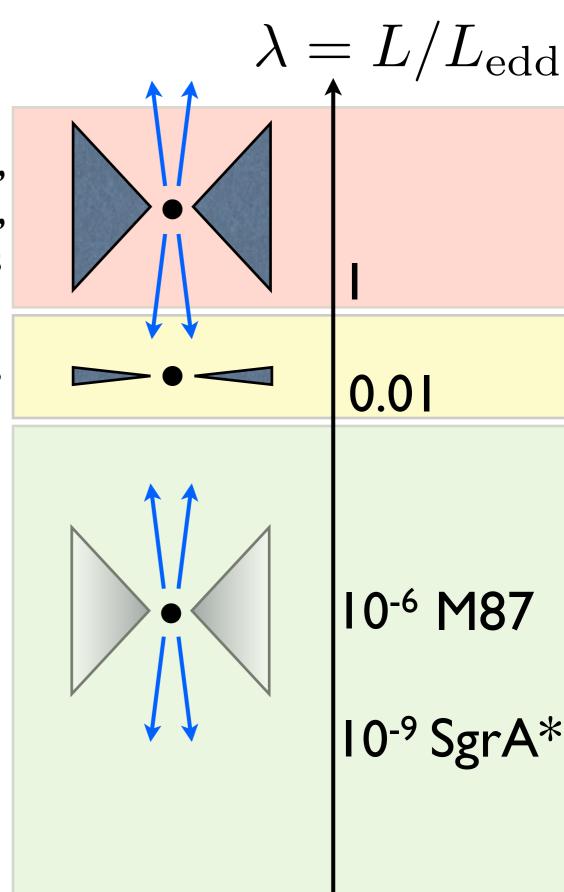
Simulated EHT image of SgrA\*

#### When are Jets Produced?

Tidal disruptions (TDEs), ultra-luminous X-ray sources, gamma-ray bursts

Quasars, X-ray binaries, TDEs

Low-luminosity active galactic nuclei (LLAGN), X-ray binaries

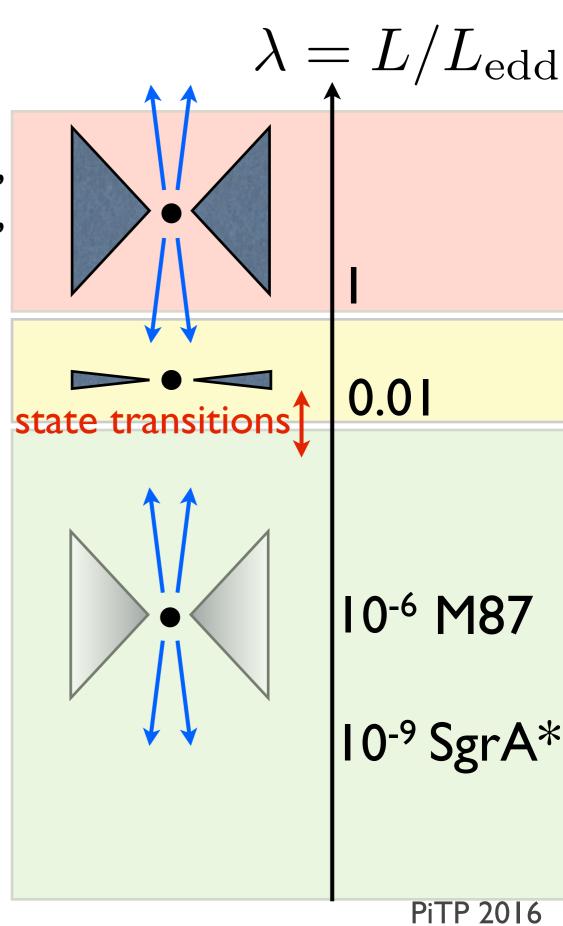


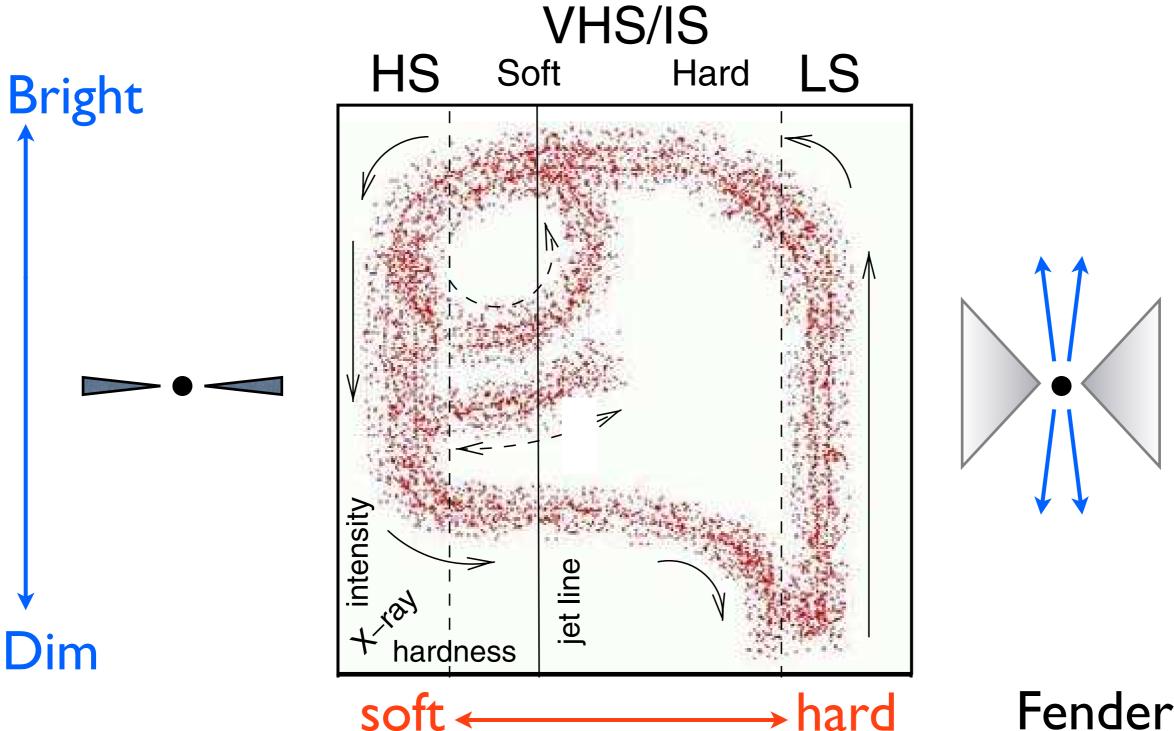
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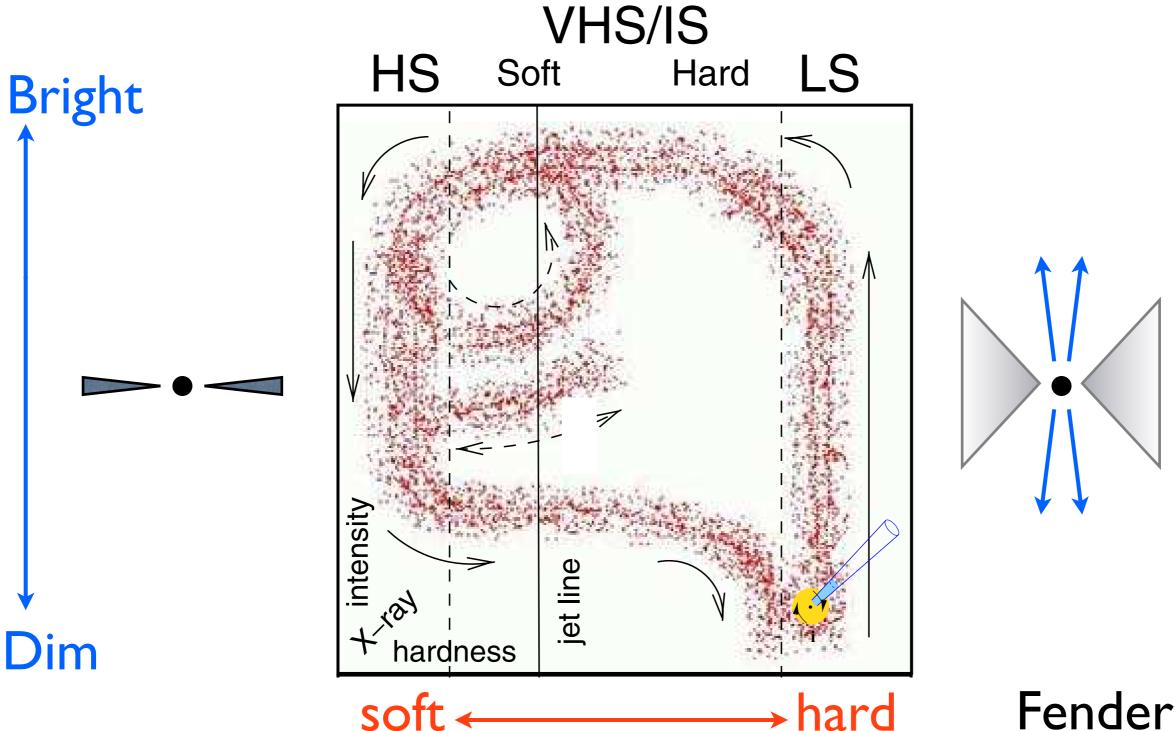
Tidal disruptions (TDEs), ultra-luminous X-ray sources, gamma-ray bursts

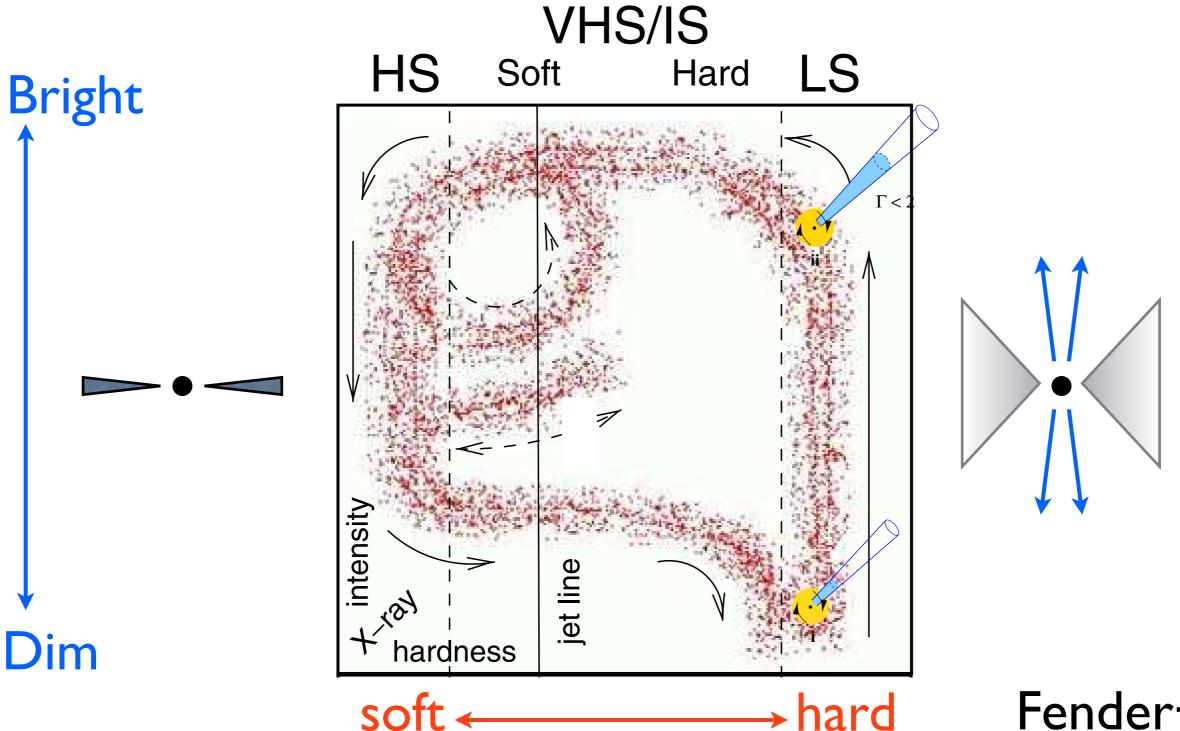
Quasars, X-ray binaries, TDEs

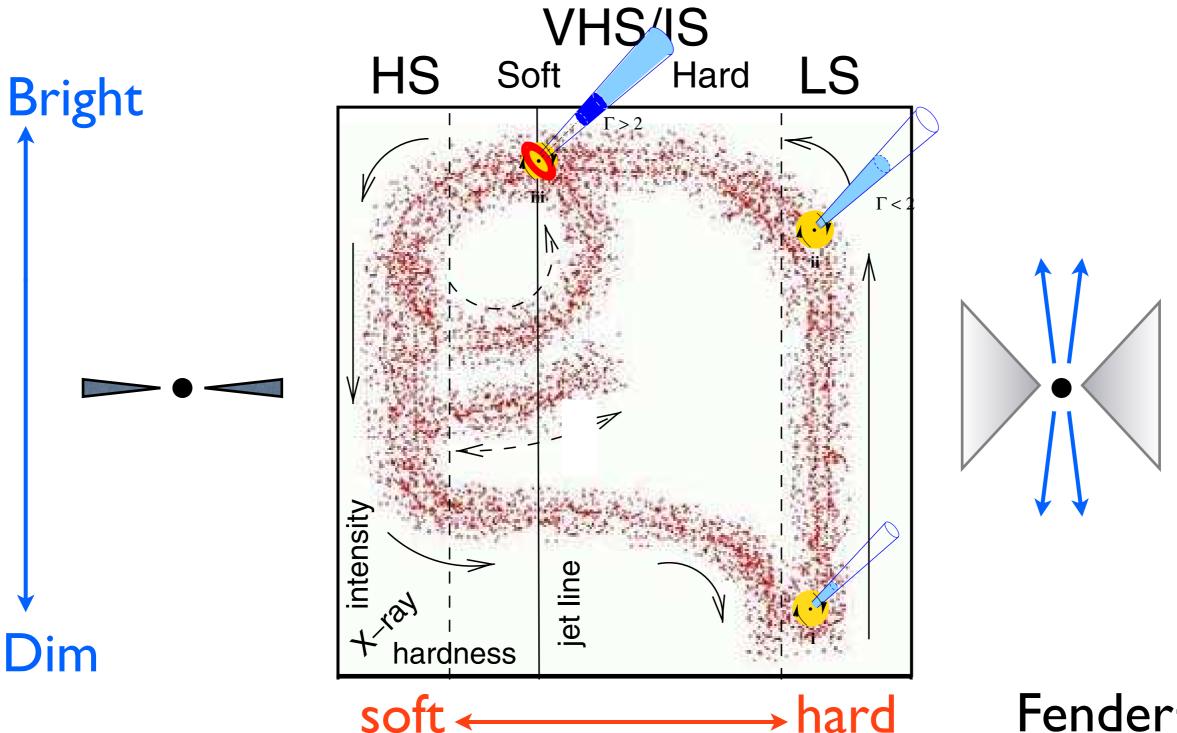
Low-luminosity active galactic nuclei (LLAGN), X-ray binaries

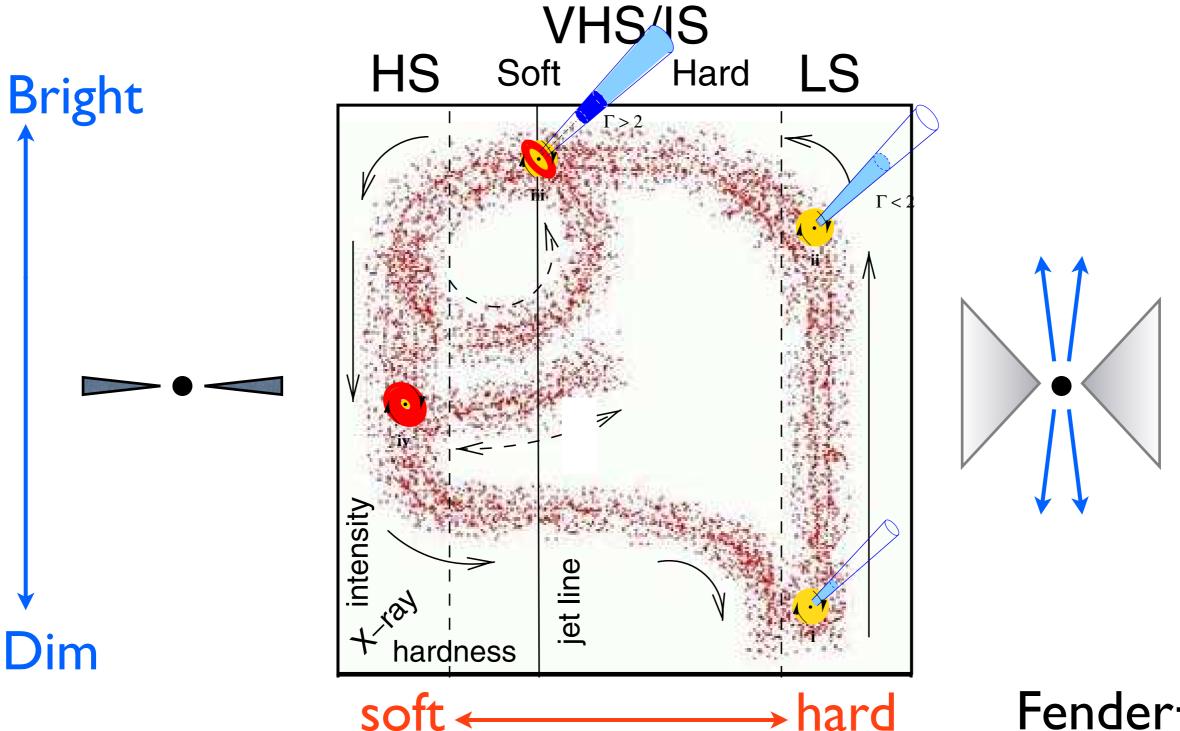






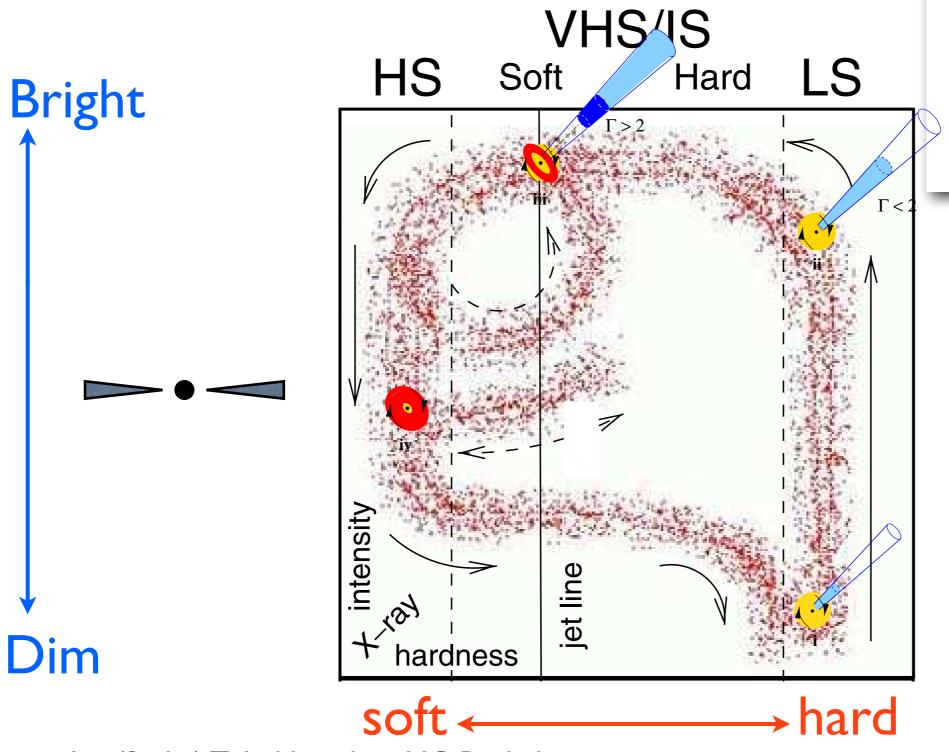




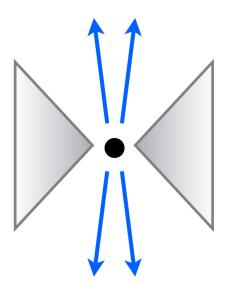


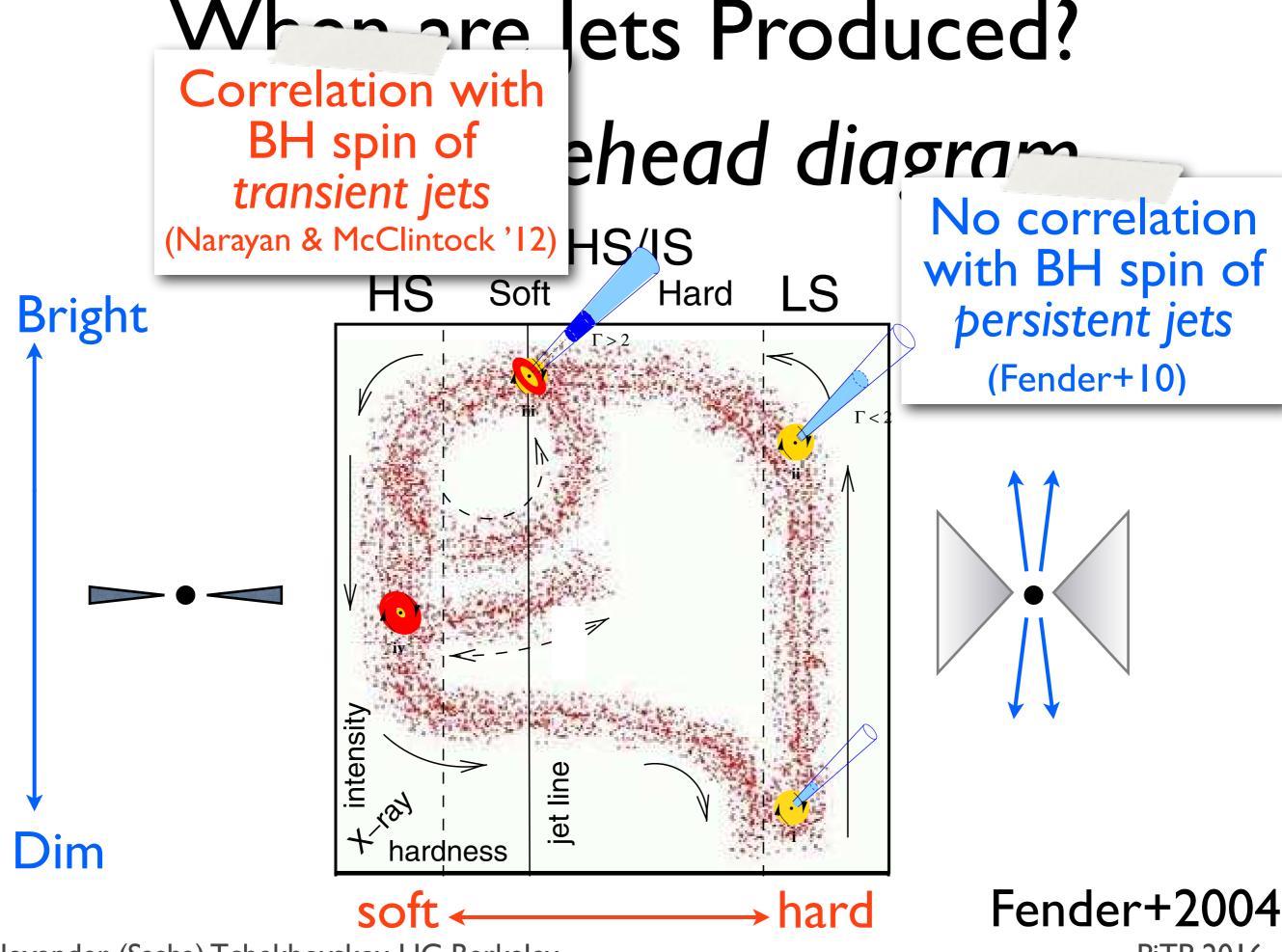
#### When are Jets Produced?

"q" or turtlehead diagram



No correlation with BH spin of persistent jets (Fender+10)





Alexander (Sasha) Tchekhovskoy, UC Berkeley

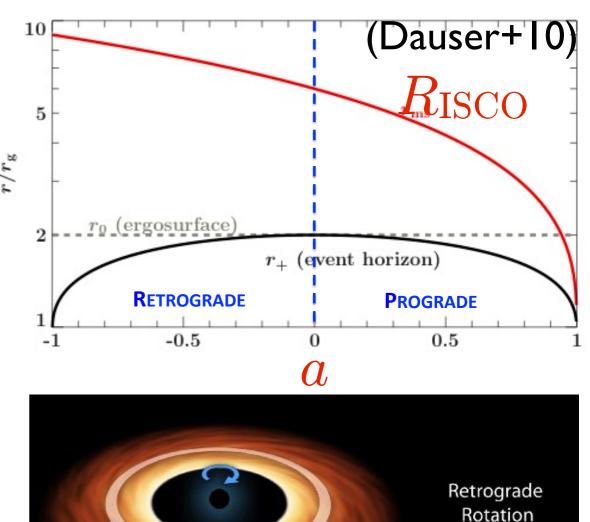
PiTP 2016

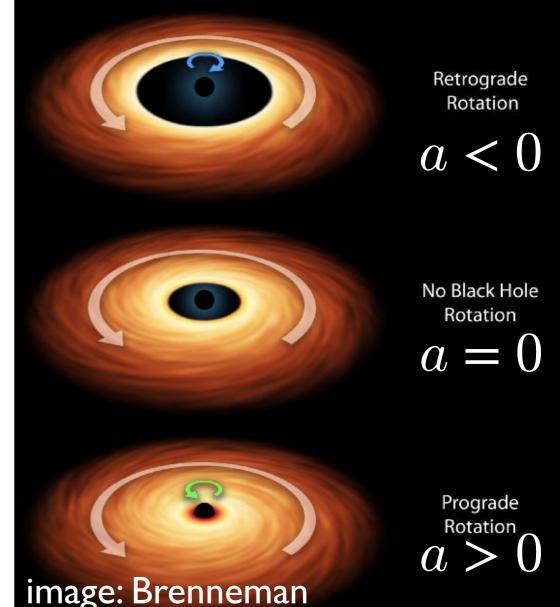
# Two major methods of measuring BH spin

- Both methods rely on measuring the size of the 'hole' in the disk
- 'Hole' size = Radius of Innermost Stable Circular Orbit (ISCO), Risco
- Continuum fitting: via blackbody-like spectrum (McClintock, Narayan, Steiner, ...)
- Iron line: via the shape of the Fe line

(Brenneman, Fabian, Reynolds, Russell, ...)

Alexander (Sasha) Tchekhovskoy, UC Berkeley



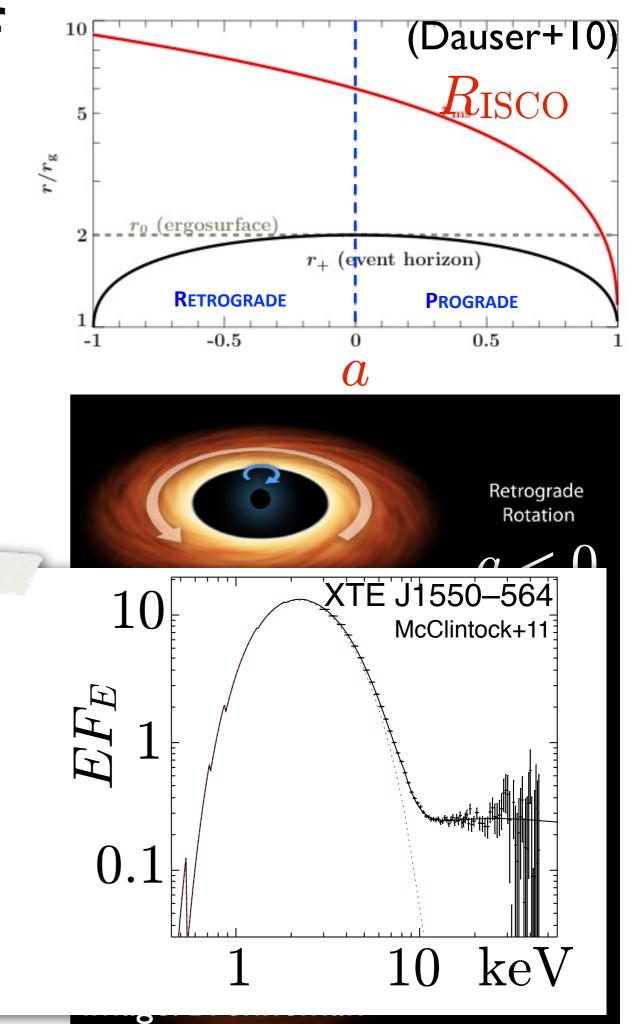


# Two major methods of measuring BH spin

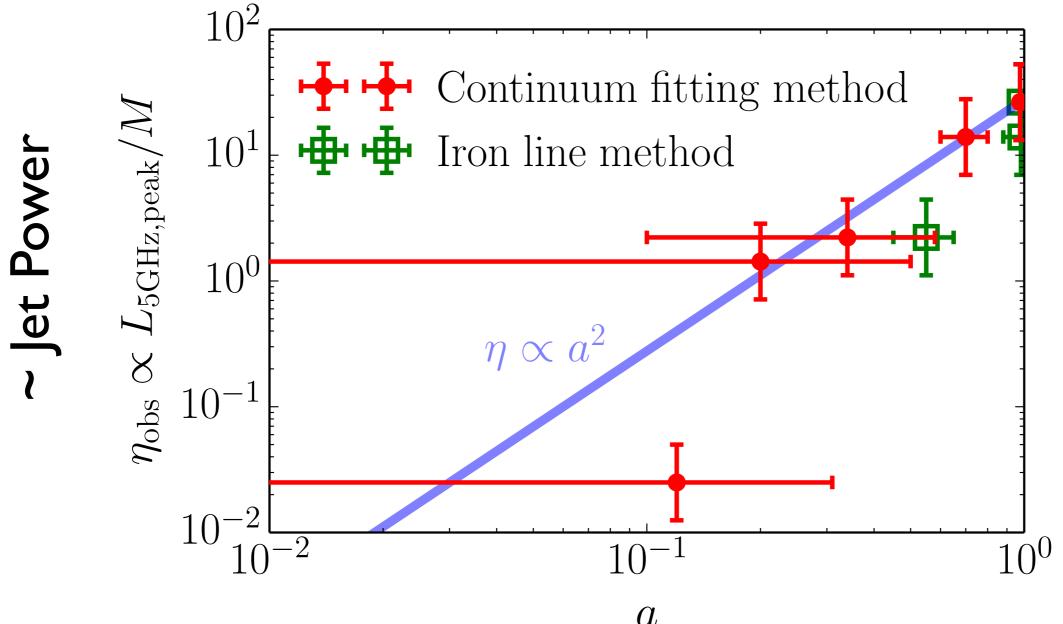
- Both methods rely on measuring the size of the 'hole' in the disk
- 'Hole' size = Radius of Innermost Stable Circular Orbit (ISCO), Risco

#### Continuum fitting method:

$$L \sim \pi R_{\rm ISCO}{}^2 \sigma T^4$$
 
$$\uparrow \qquad \uparrow$$
 normalization shape of the spectrum



# Are Transient Jets Powered by BH spin?



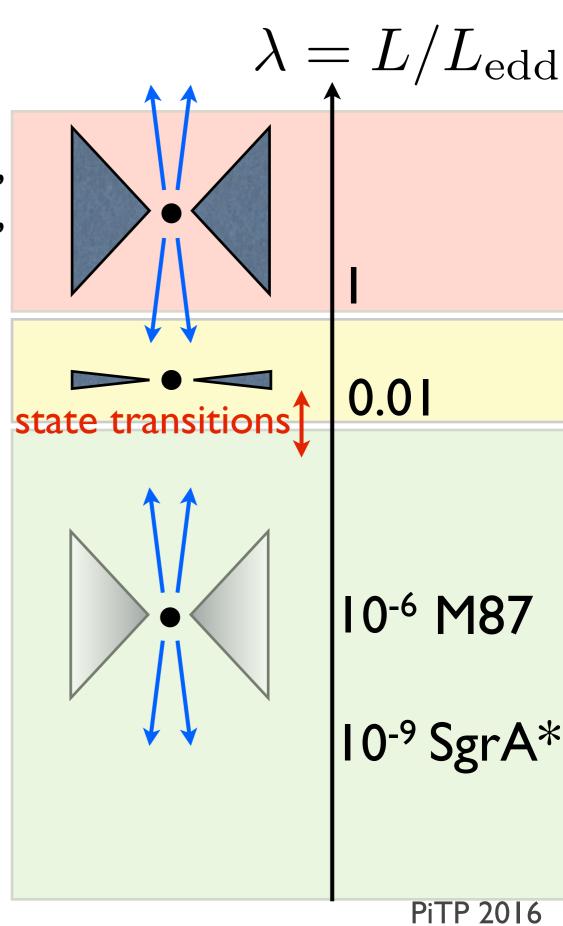
AT15; data Steiner+13, Narayan & McClintock+12 Alexander (Sasha) Tchekhovskoy, UC Berkeley

#### When are Jets Produced?

Tidal disruptions (TDEs), ultra-luminous X-ray sources, gamma-ray bursts

Quasars, X-ray binaries, TDEs

Low-luminosity active galactic nuclei (LLAGN), X-ray binaries



#### When are Jets Produced?

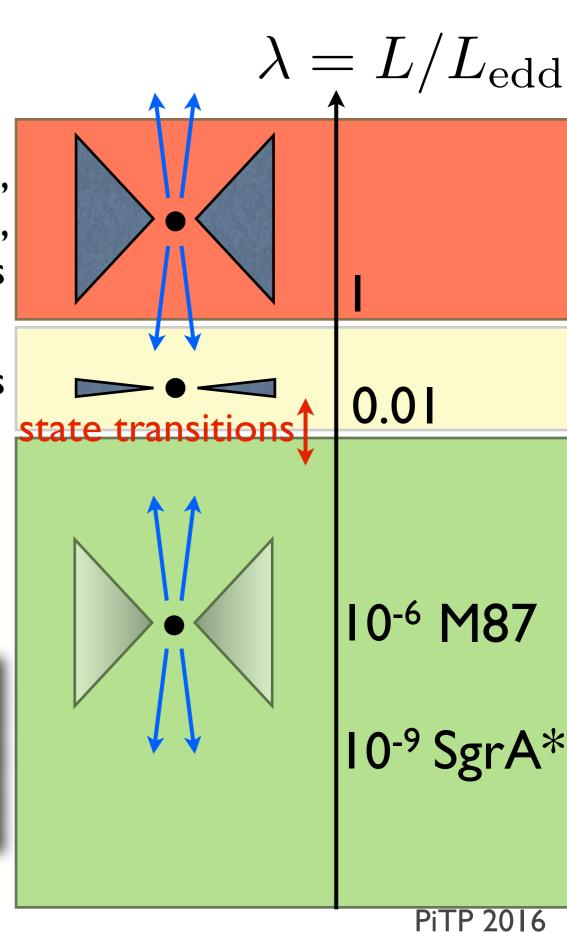
Tidal disruptions (TDEs), ultra-luminous X-ray sources, gamma-ray bursts

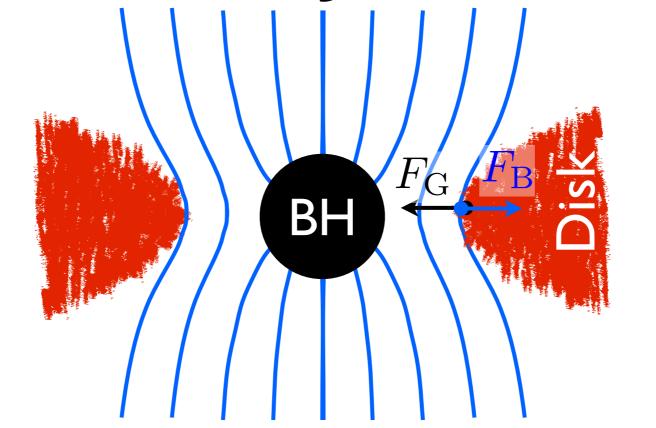
Quasars, X-ray binaries, TDEs

Low-luminosity active galactic nuclei (LLAGN), X-ray binaries

Both high- and low-luminosity disks are *radiatively inefficient*.

Neglect radiation and simulate.





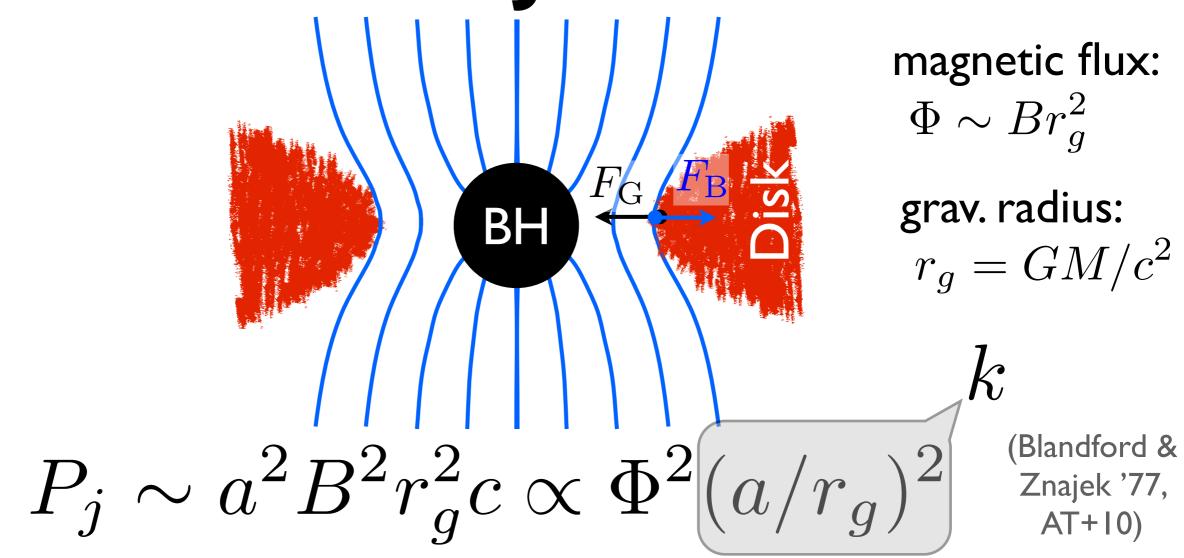
magnetic flux:

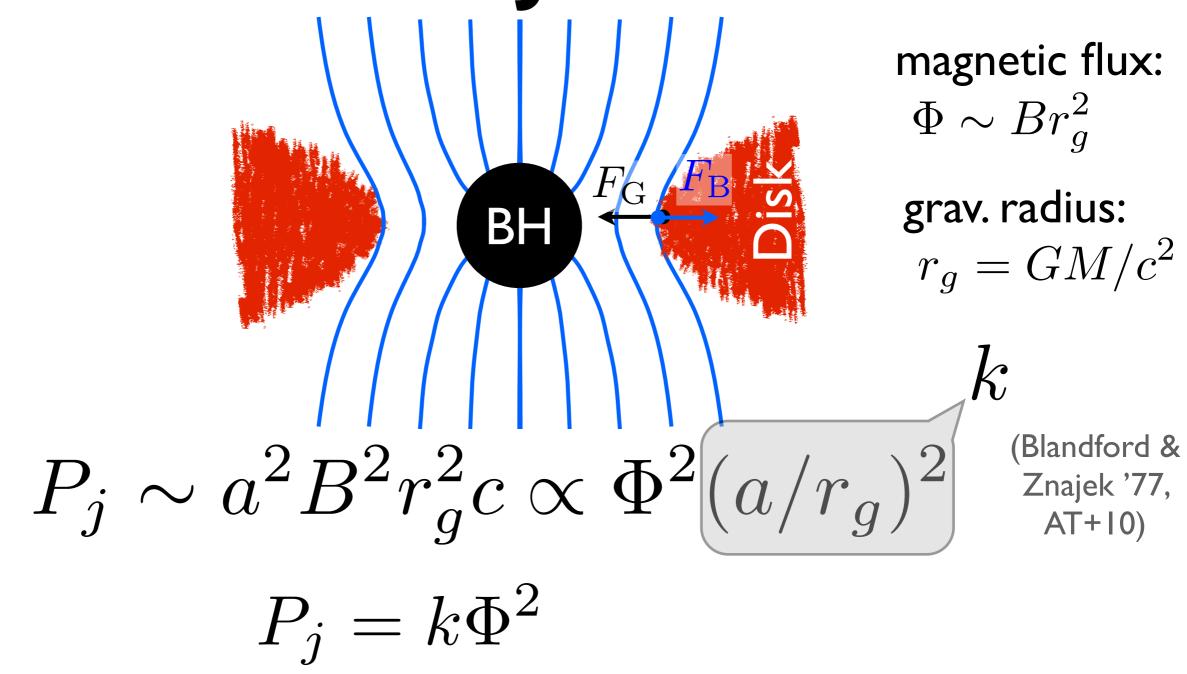
$$\Phi \sim Br_g^2$$

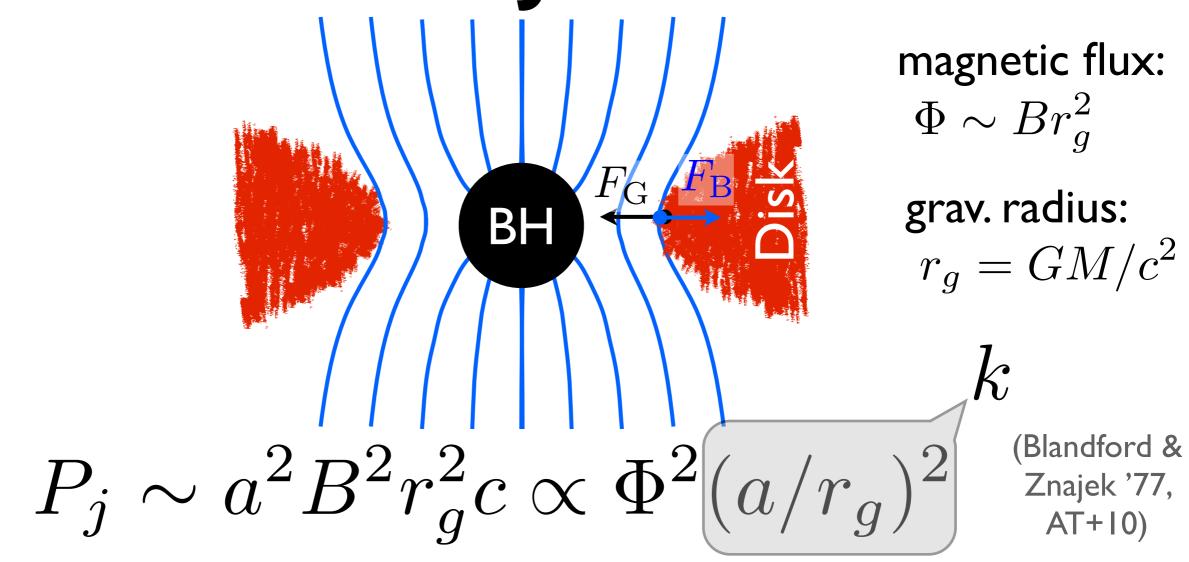
grav. radius:

$$r_g = GM/c^2$$

$$P_j \sim a^2 B^2 r_g^2 c \propto \Phi^2 (a/r_g)^2 \stackrel{\text{(Blandford \& Znajek '77, AT+10)}}{}$$





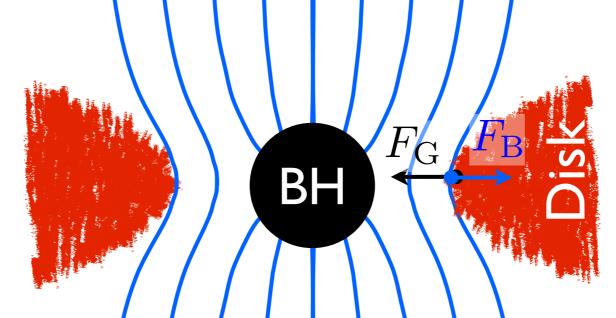


B subdominant

$$0 \le P_j = k\Phi^2$$

$$\Phi = 0$$





magnetic flux:

$$\Phi \sim B r_g^2$$

grav. radius:

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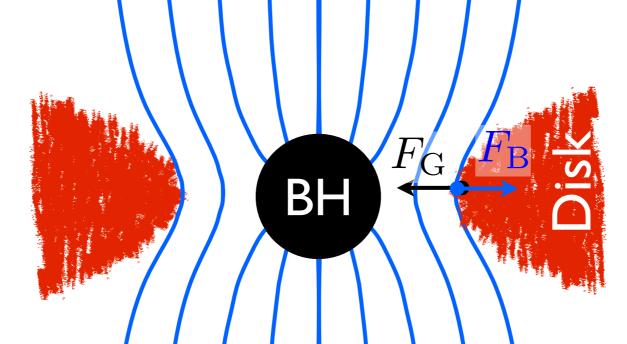
$$P_j \sim a^2 B^2 r_g^2 c \propto \Phi^2 (a/r_g)^2$$
 (Blandford & Znajek '77, AT+10)

$$0 \le P_j = k\Phi^2 \lesssim \dot{M}c^2$$

$$\Phi = 0 \qquad \Phi = \dot{\Phi}_{MAX}$$

**B** dominant





magnetic flux:

$$\Phi \sim B r_g^2$$

grav. radius:

$$r_g = GM/c^2$$

$$P_j \sim a^2 B^2 r_g^2 c \propto \Phi^2 (a/r_g)^2$$
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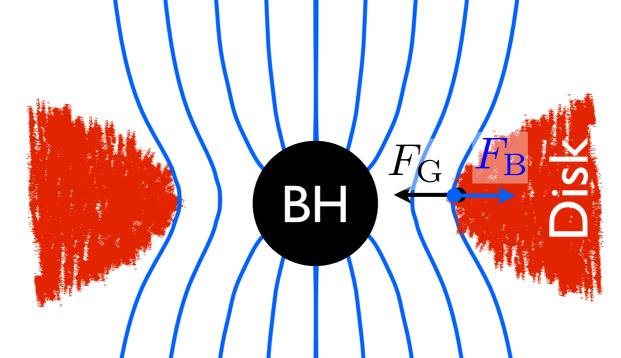
B sub-dominant

$$\begin{array}{ccc}
0 \leq P_j = k\Phi^2 \lesssim \dot{M}c^2 \\
\Phi = 0 & \Phi = \Phi_{\text{MAX}}
\end{array}$$

B dominant Magnetically-Arrested Disk (MAD)

(Narayan+ 2003, AT+ 2011)





magnetic flux:

$$\Phi \sim B r_g^2$$

grav. radius:

$$r_g = GM/c^2$$

$$P_j \sim a^2 B^2 r_g^2 c \propto \Phi^2 (a/r_g)^2$$
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B sub-dominant

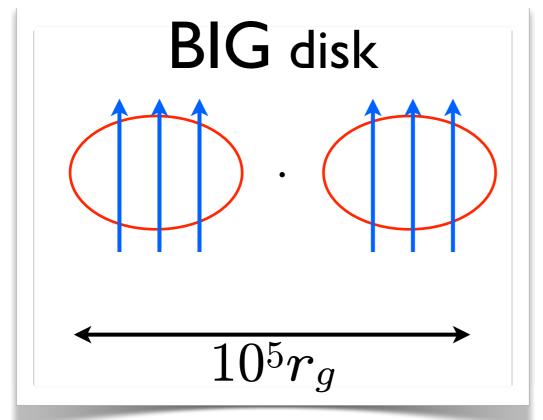
ont 
$$\Phi = 0$$
  $\Phi^2 \lesssim \dot{M}c^2$  ant  $\Phi = 0$   $\Phi = \dot{\Phi}_{\rm MAX}$ 

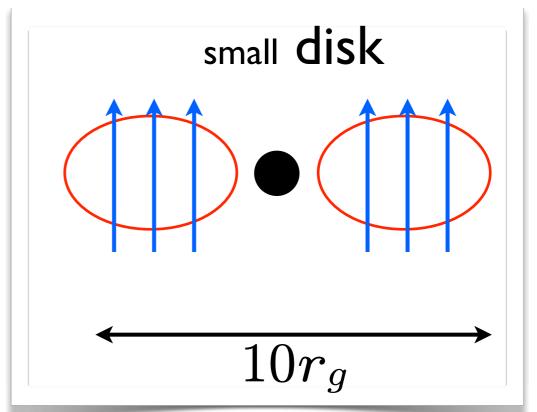
How strong are the jets?  $p_j = P_j / \dot{M}c^2$ 

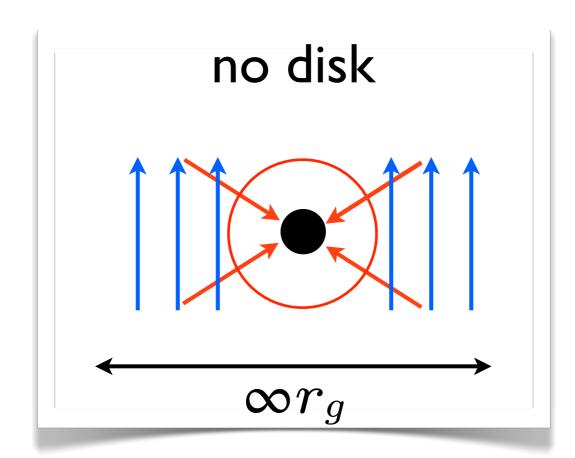
B dominant Magnetically-Arrested Disk (MAD)

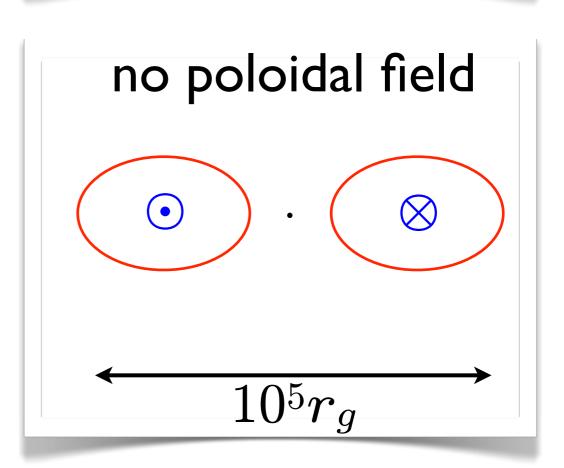
(Narayan+ 2003, AT+ 2011)

### What is a Healthy Jet Diet?

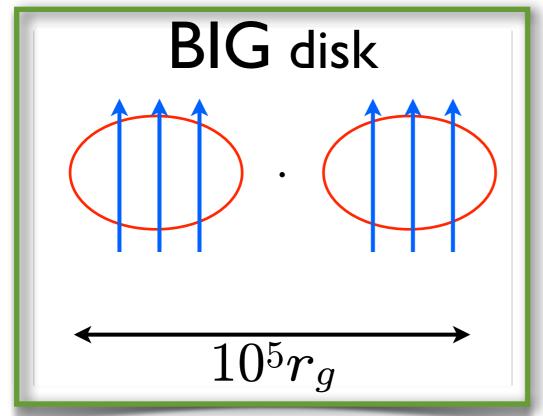


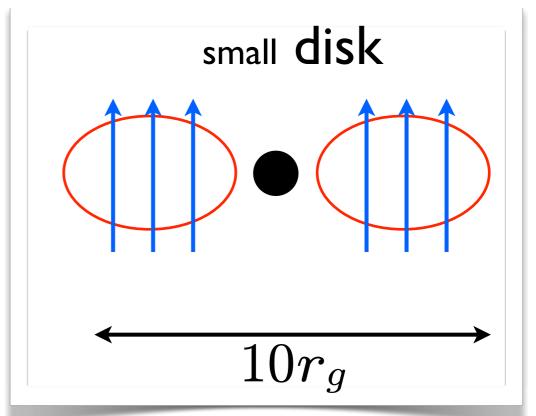


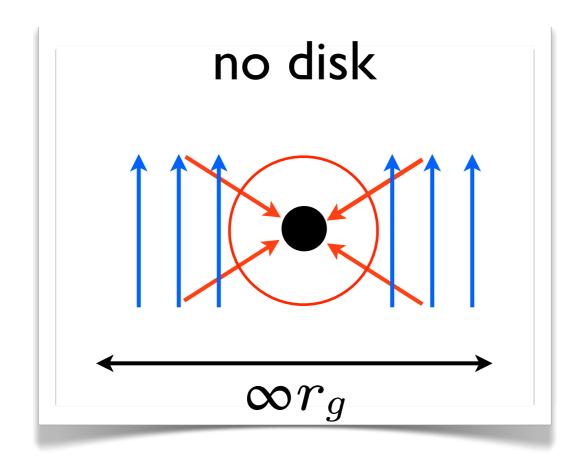


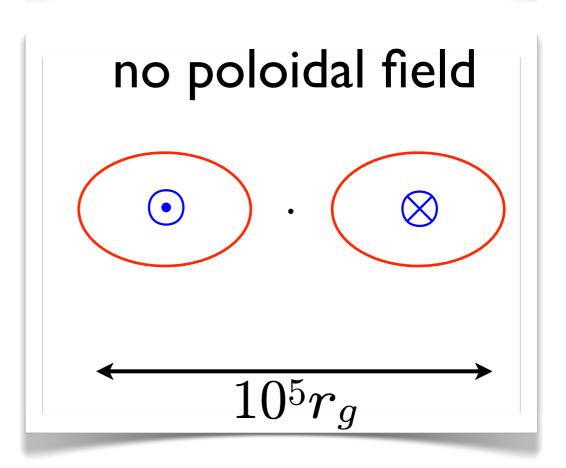


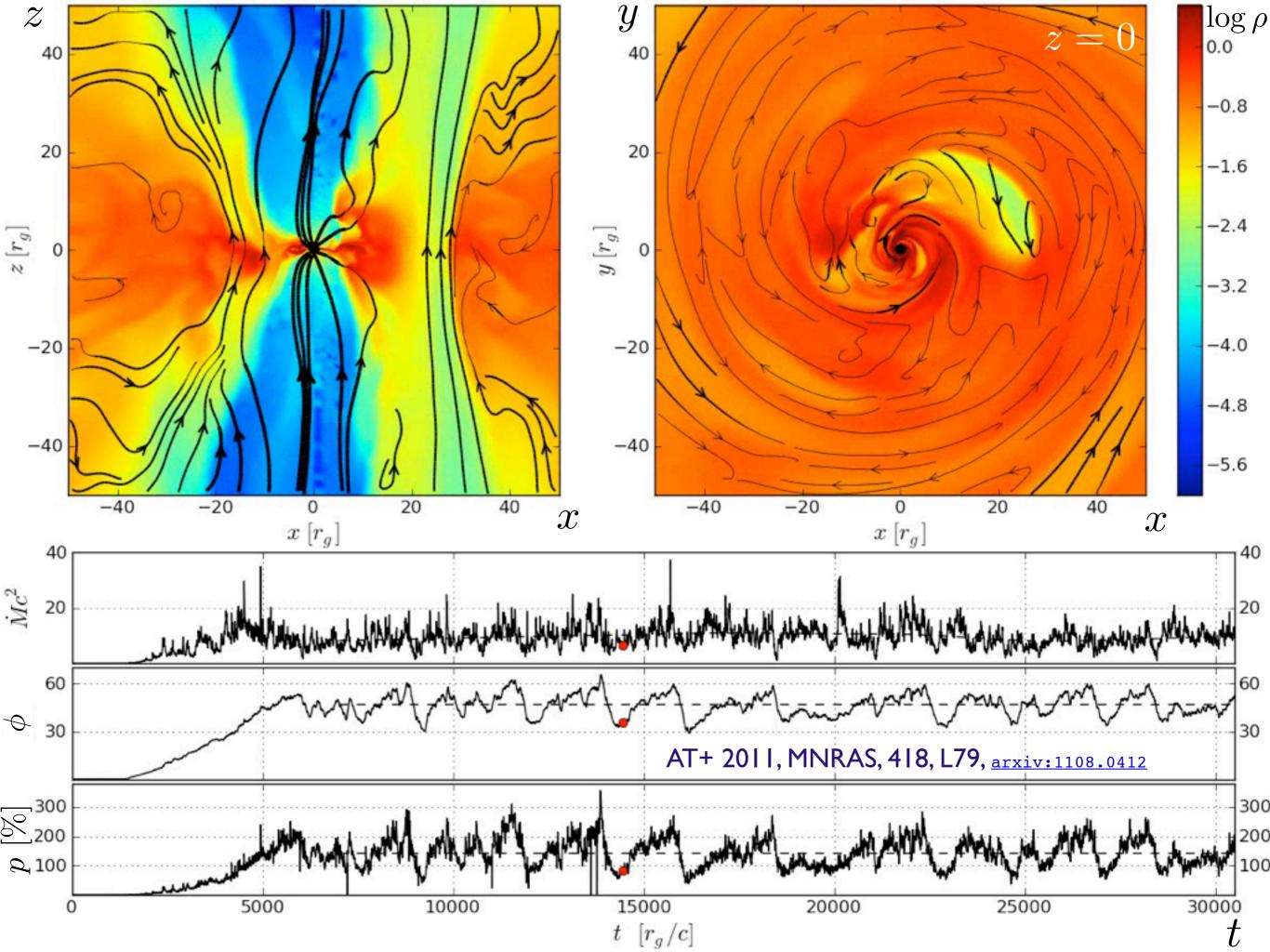
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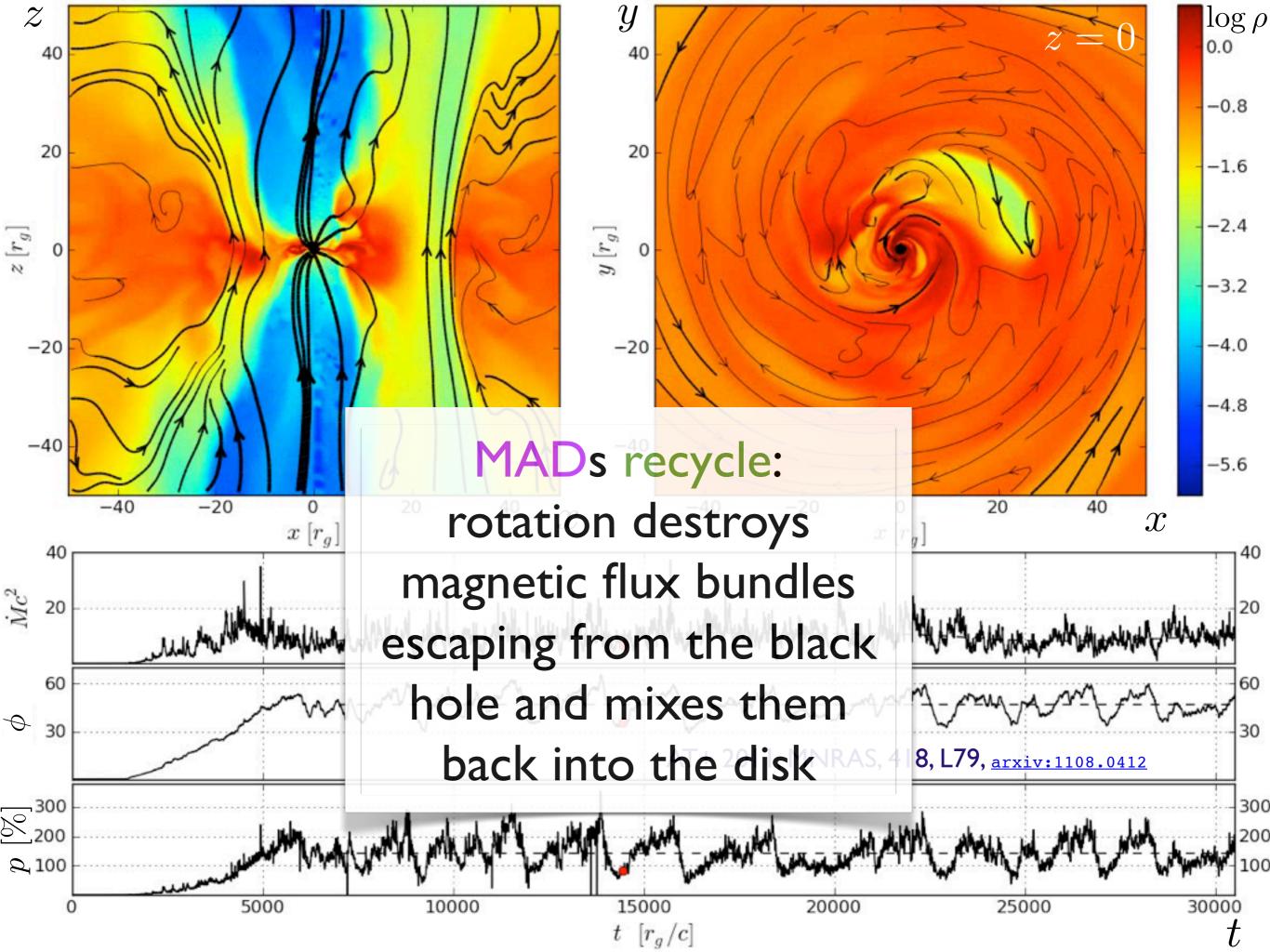










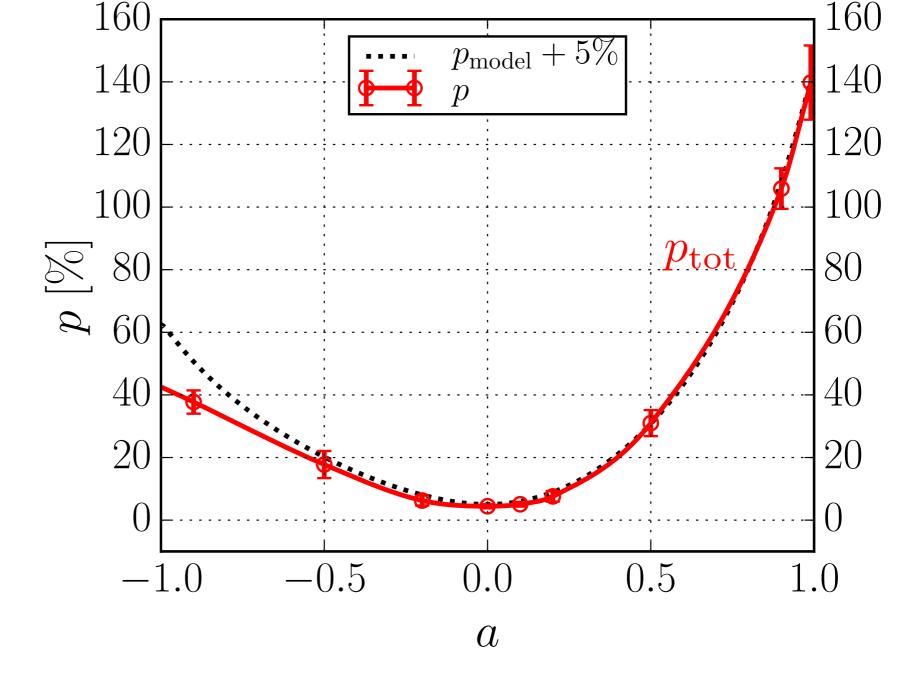


# Upper Envelope of Jet Power vs. Spin

 $(h/r \sim 0.3)$ 

(Tchekhovskoy+ 11;
Tchekhovskoy, McKinney 12;
McKinney, Tchekhovskoy,
Blandford 12;
Tchekhovskoy 15)

Quantify feedback due to black hole jet, disk wind from first principles



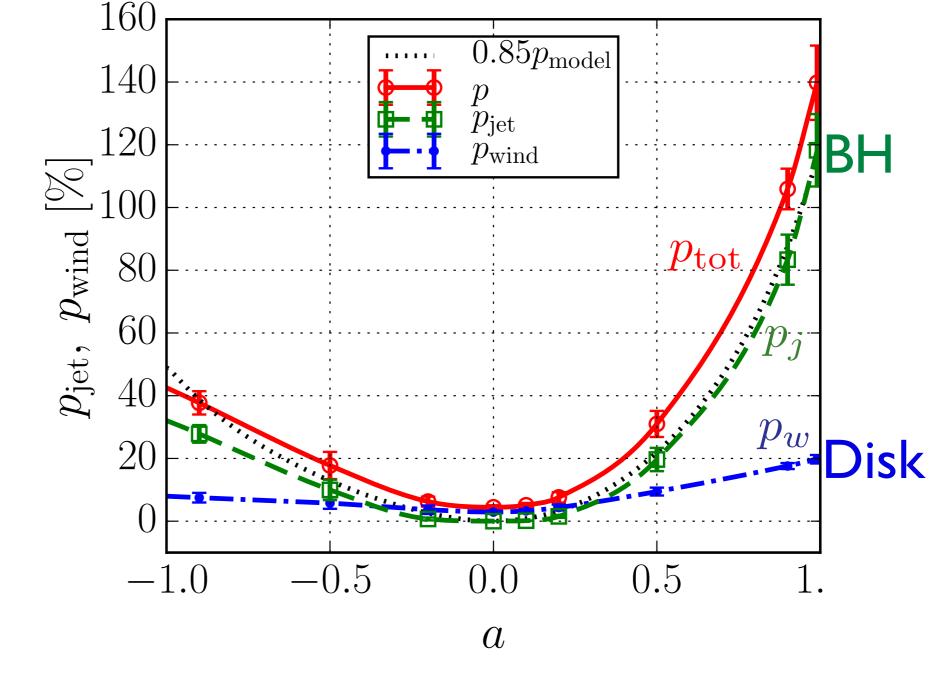
p>100% means net energy is extracted from the BH

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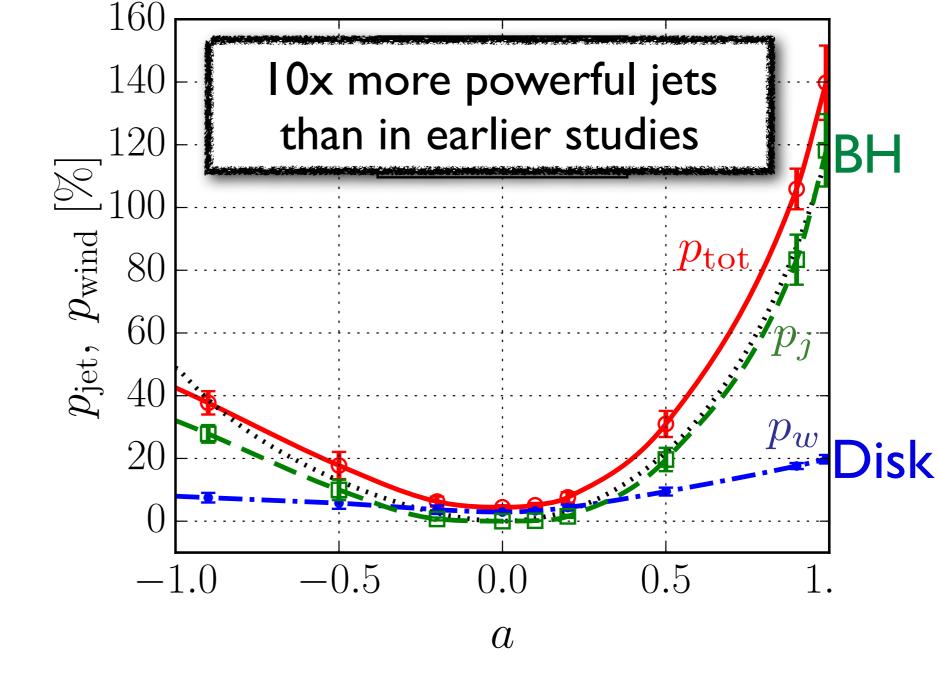
Jet = 85% of Blandford-Znajek power
Wind = BP = 15% of BZ power + 5%
Disk wind is powered by a combination of BH
spin and disk rotation

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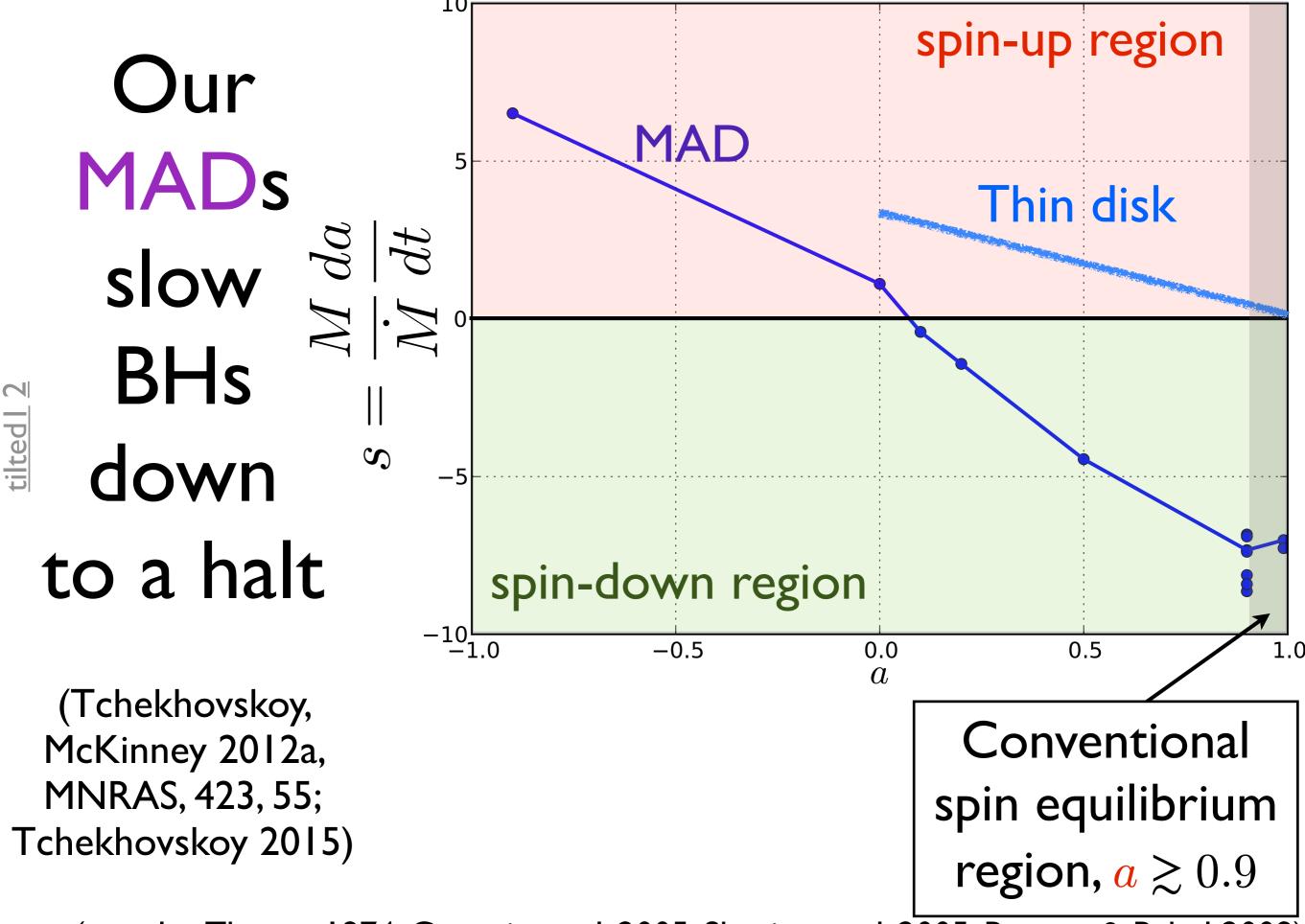
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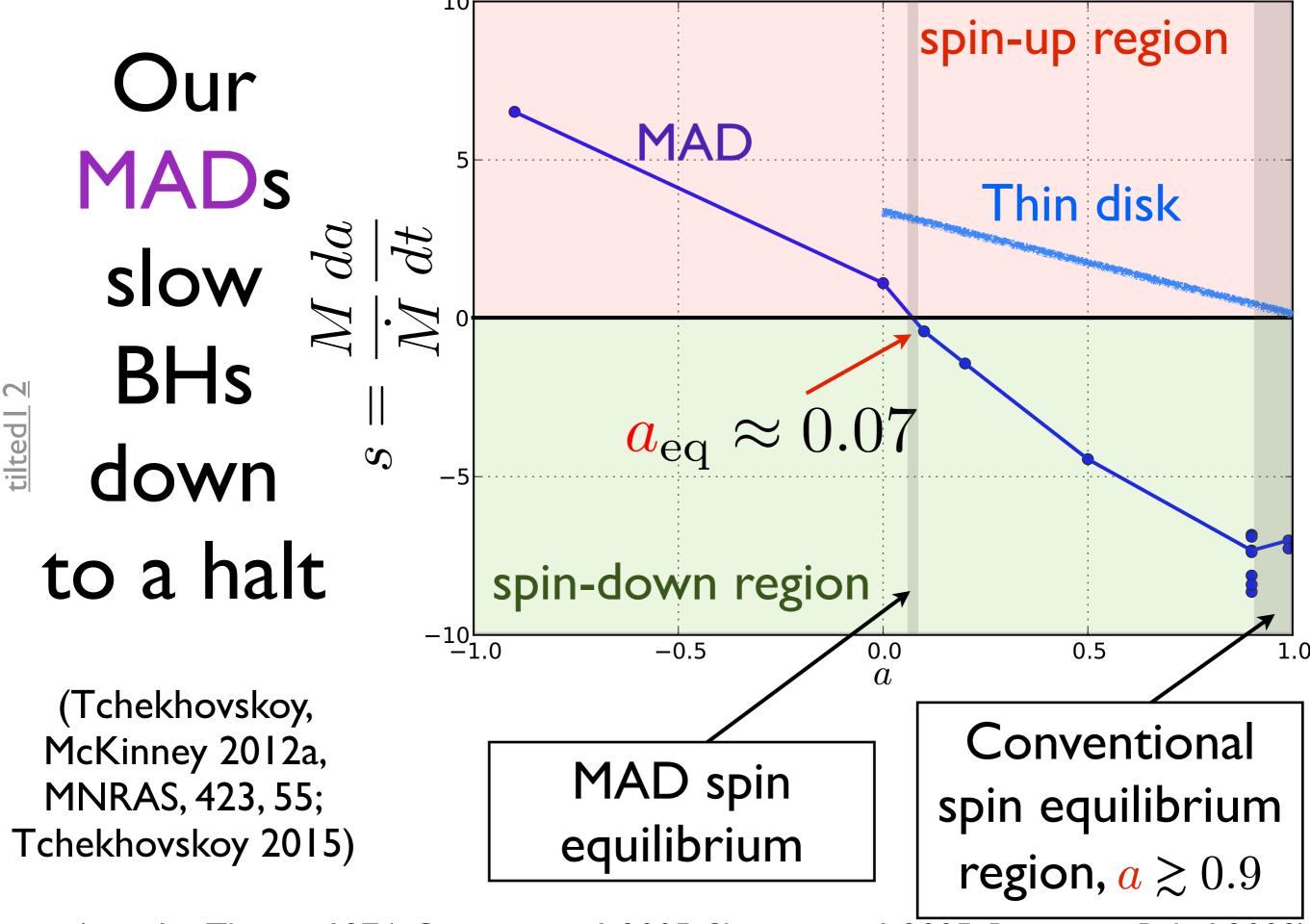
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(see also Thorne 1974, Gammie et al. 2005, Shapiro et al. 2005, Benson & Babul 2009)



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#### When are Jets Produced?

Dynamically important magnetic fields:

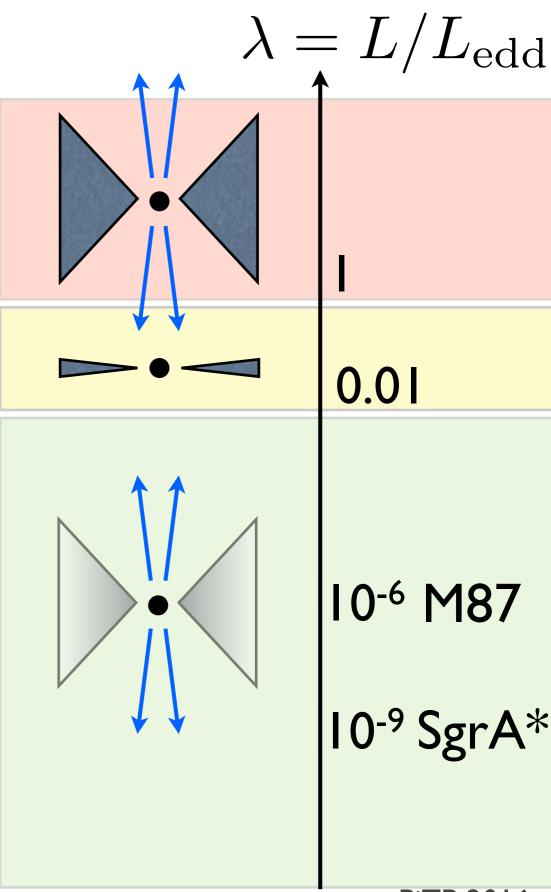
(AT+13, Tidal disruptions (TDEs), AT & Giannios 15) ultra-luminous X-ray sources, gamma-ray bursts

(Zamaninasab ++AT 14, Ghisellini+14)

Blazars, X-ray binaries, TDEs

(Nemmen & AT 14)

Low-luminosity active galactic nuclei (LLAGN), X-ray binaries



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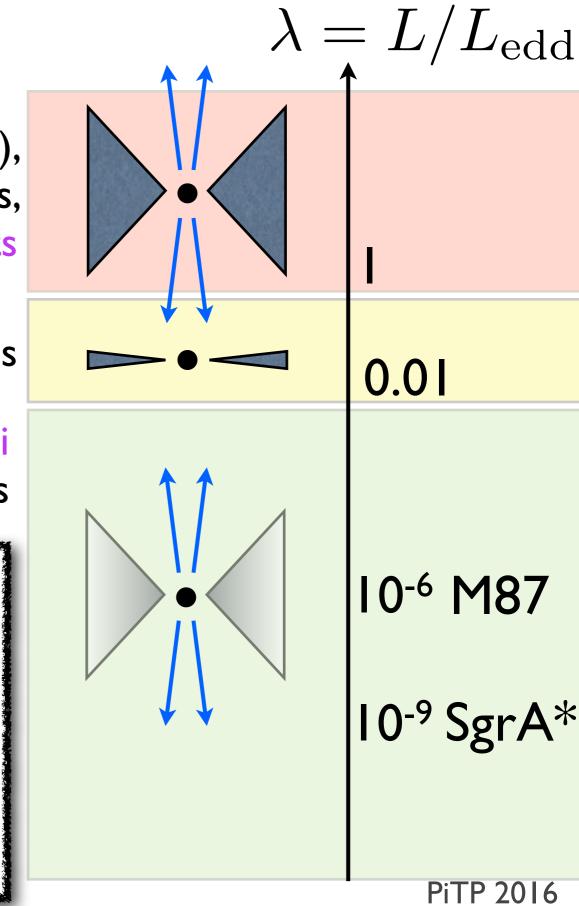
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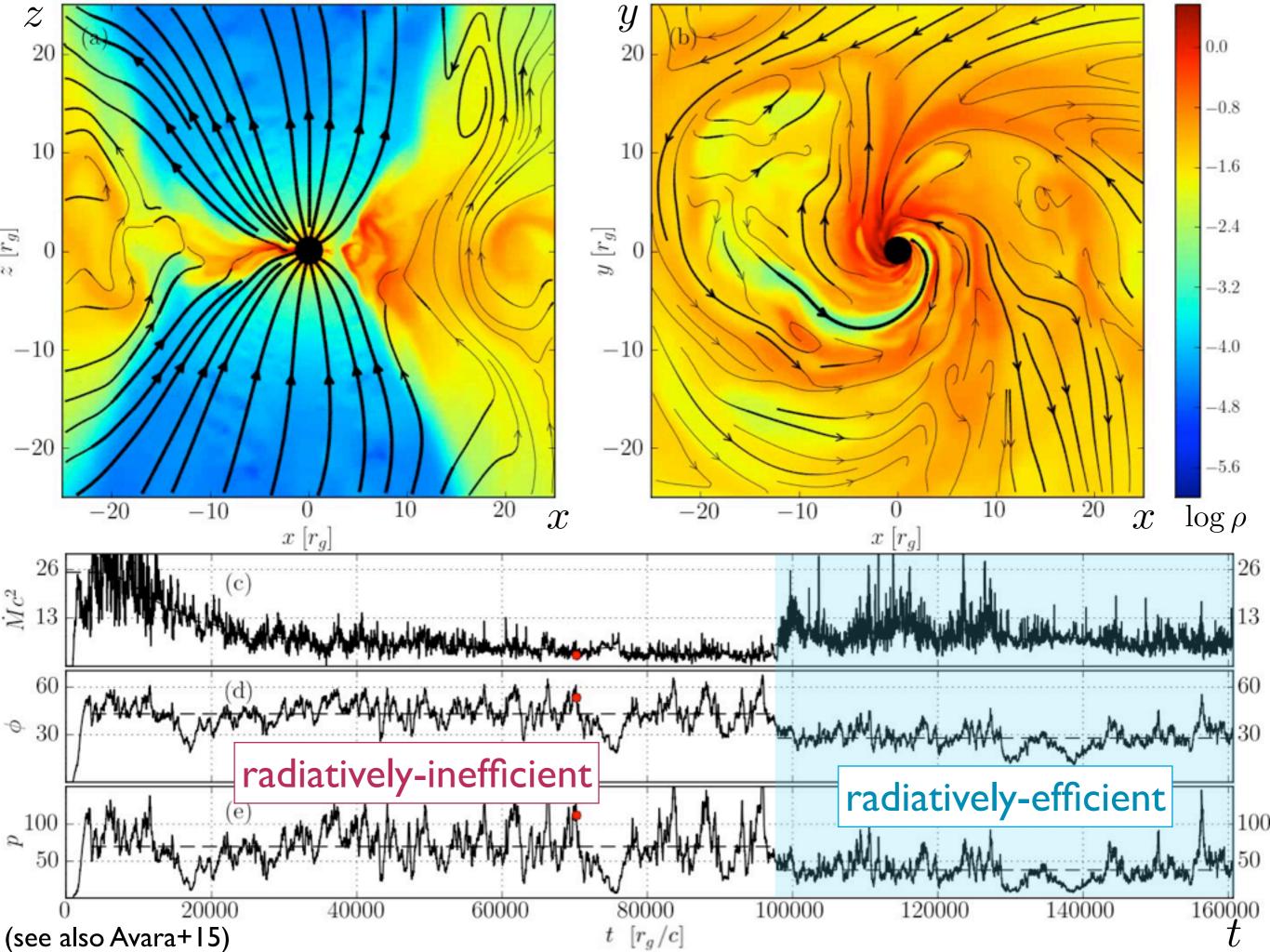
(Nemmen & AT 14)

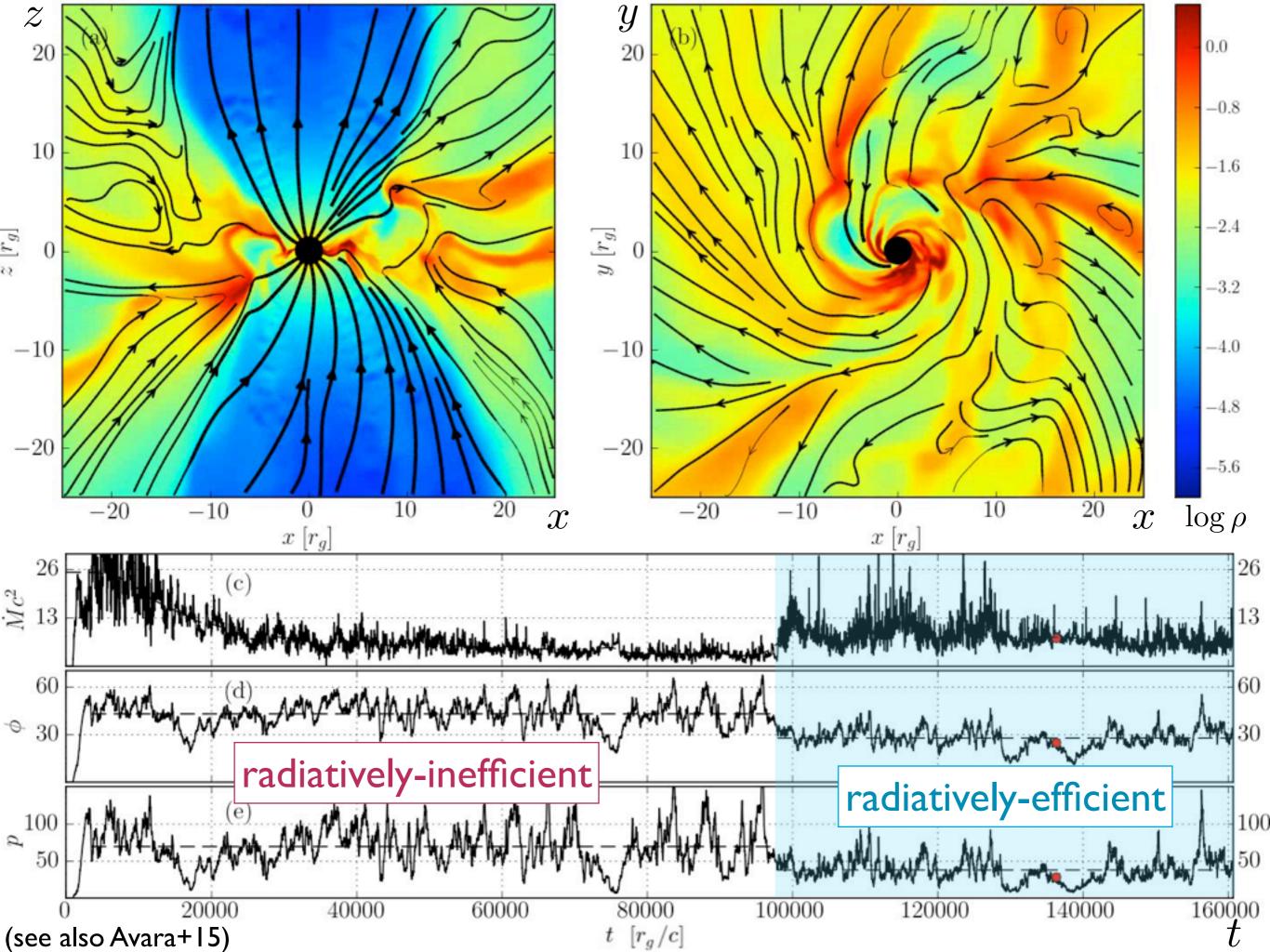
Low-luminosity active galactic nuclei (LLAGN), X-ray binaries

- Can radiatively-efficient disks drag large-scale magnetic flux in and make powerful jets?
- Analytical studies: does not seem so! (Lubow et al 1994, Guilet & Oglivie 2013a,b)
- But then, how do quasars make jets?



Alex





### How strong is B in Blazars?

 Radio jet core is where jet becomes transparent to its own synchrotron radiation:

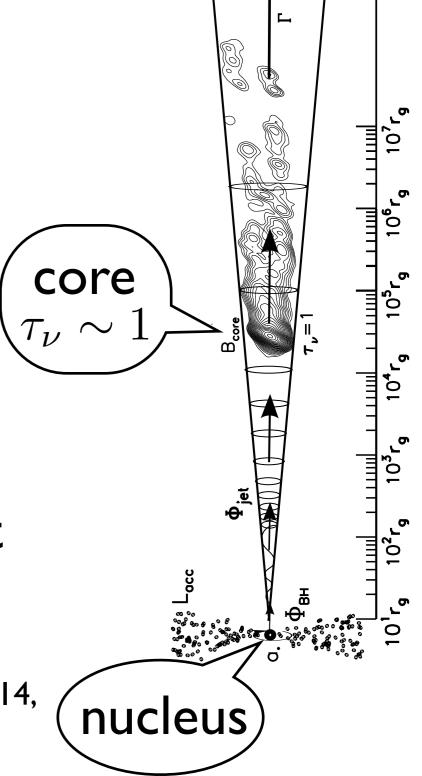
$$\tau_{\nu} \sim 1$$

• At higher  $\nu$ , the core shifts inward

$$B \propto (dr_{\rm core})^{3/4}$$

- ullet Can use this to measure B in the jet
- Magnetic flux  $\Phi \approx B\pi r_{\rm core}^2 \theta_j^2$

(Zamaninasab, Clausen-Brown, Savolainen, Tchekhovskoy, Nature, 2014, Zdziarski, Sikora, Pjanka, Tchekhovskoy, MNRAS, 2015)

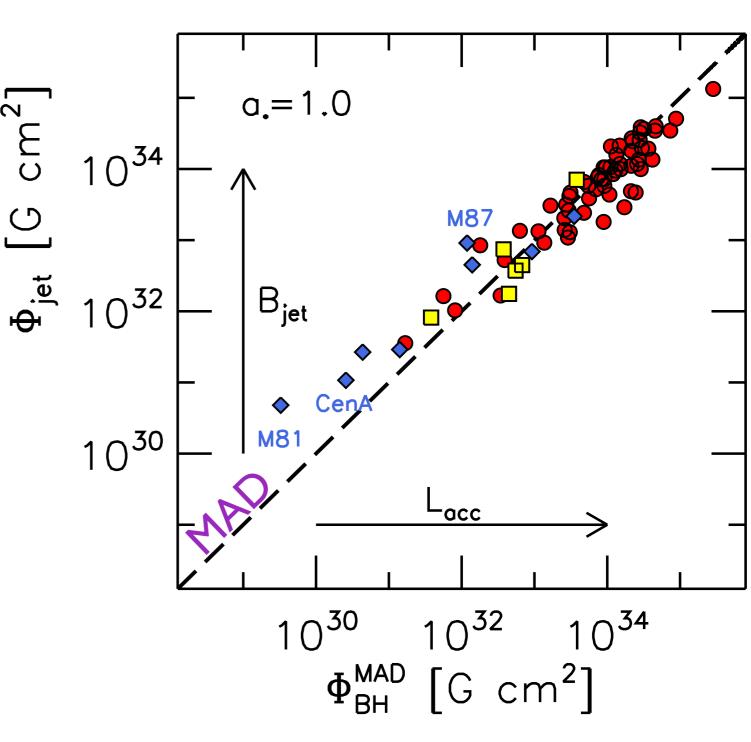


#### MADs in Blazars?

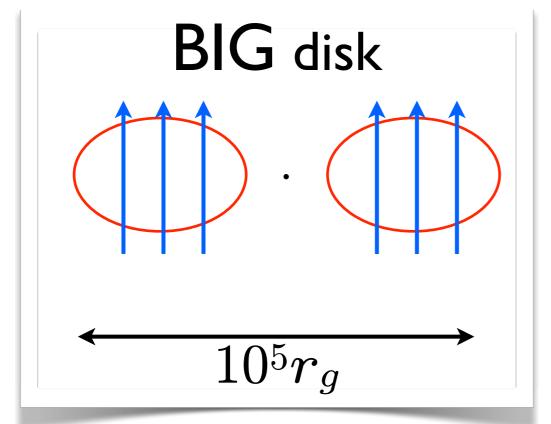
Observed scaling:

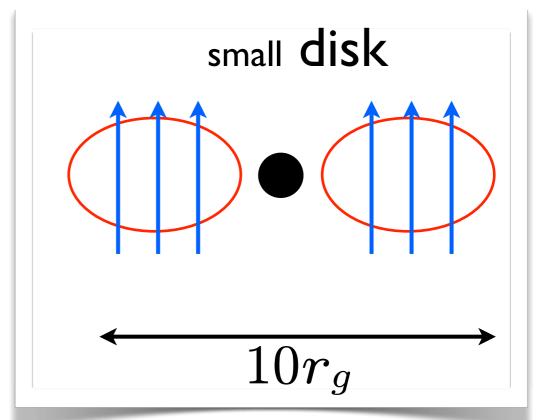
$$B_{
m jet} \propto L_{
m acc}^{1/2}$$

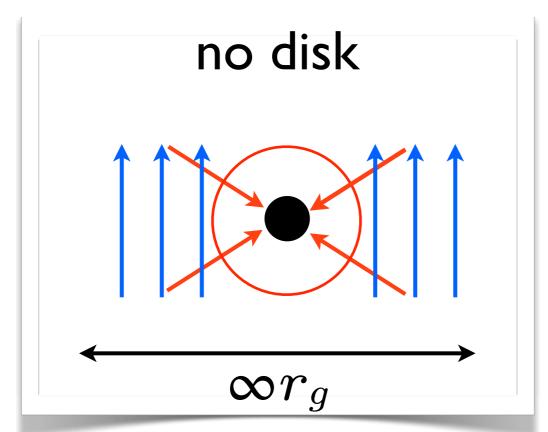
- Magnitude of magnetic flux in radio-loud AGN is consistent with MAD expectation (Tchekhovskoy, McKinney, Narayan, MNRAS, 2011)
- Many AGN are MAD
  - their central BHs are surrounded by dynamically important magnetic field

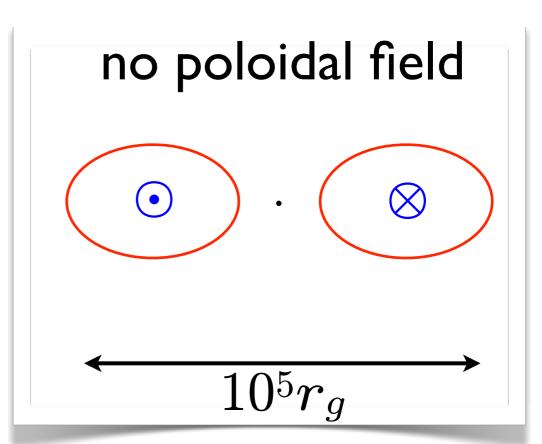


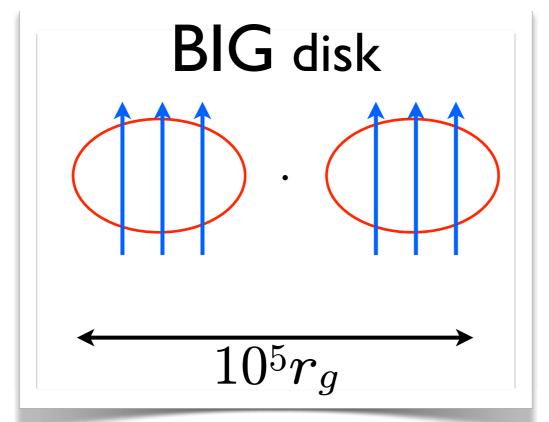
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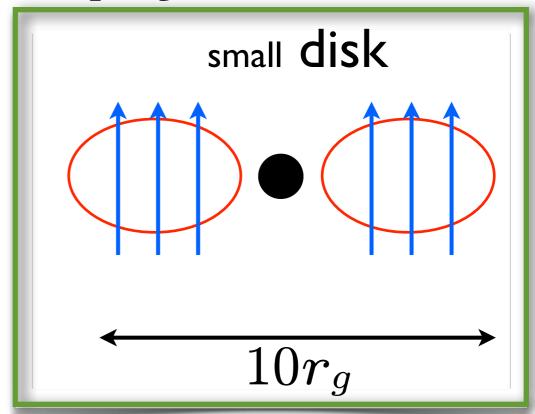


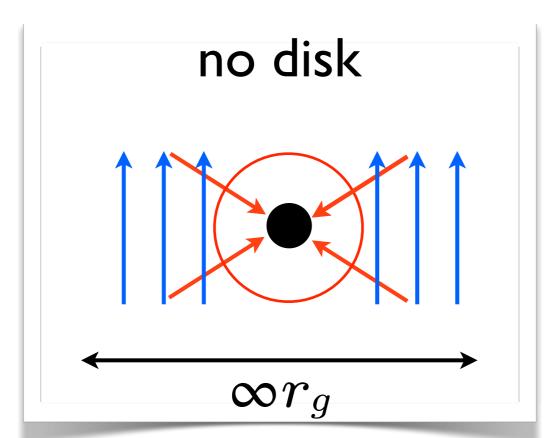


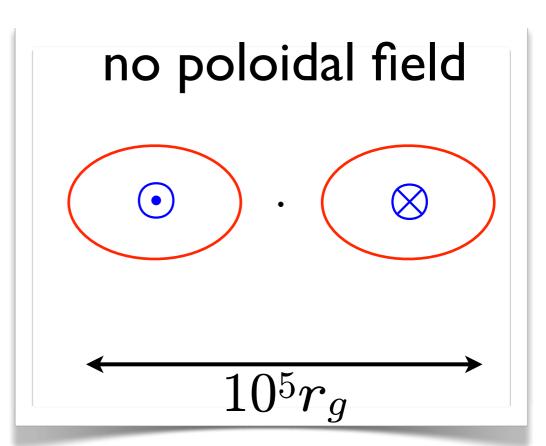




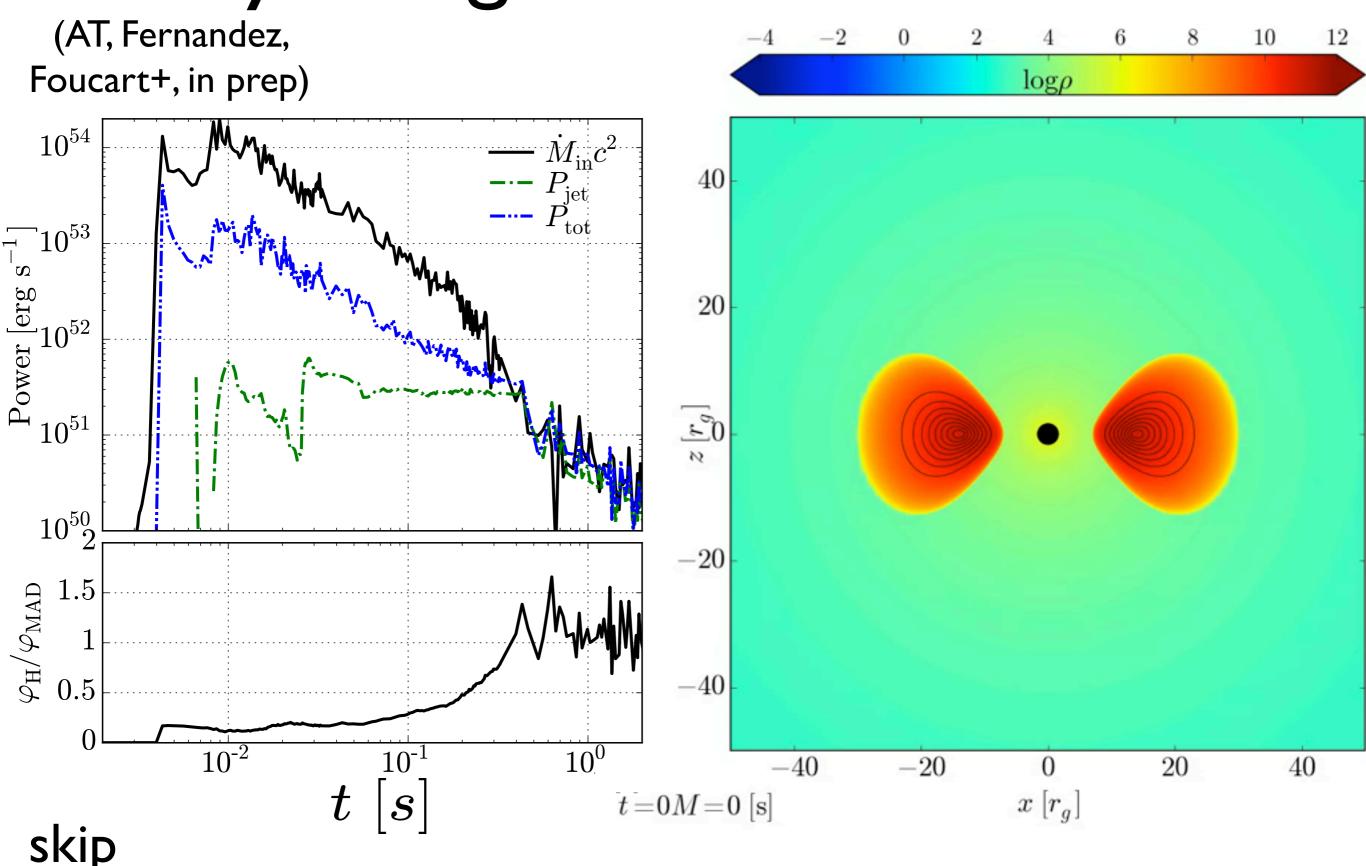






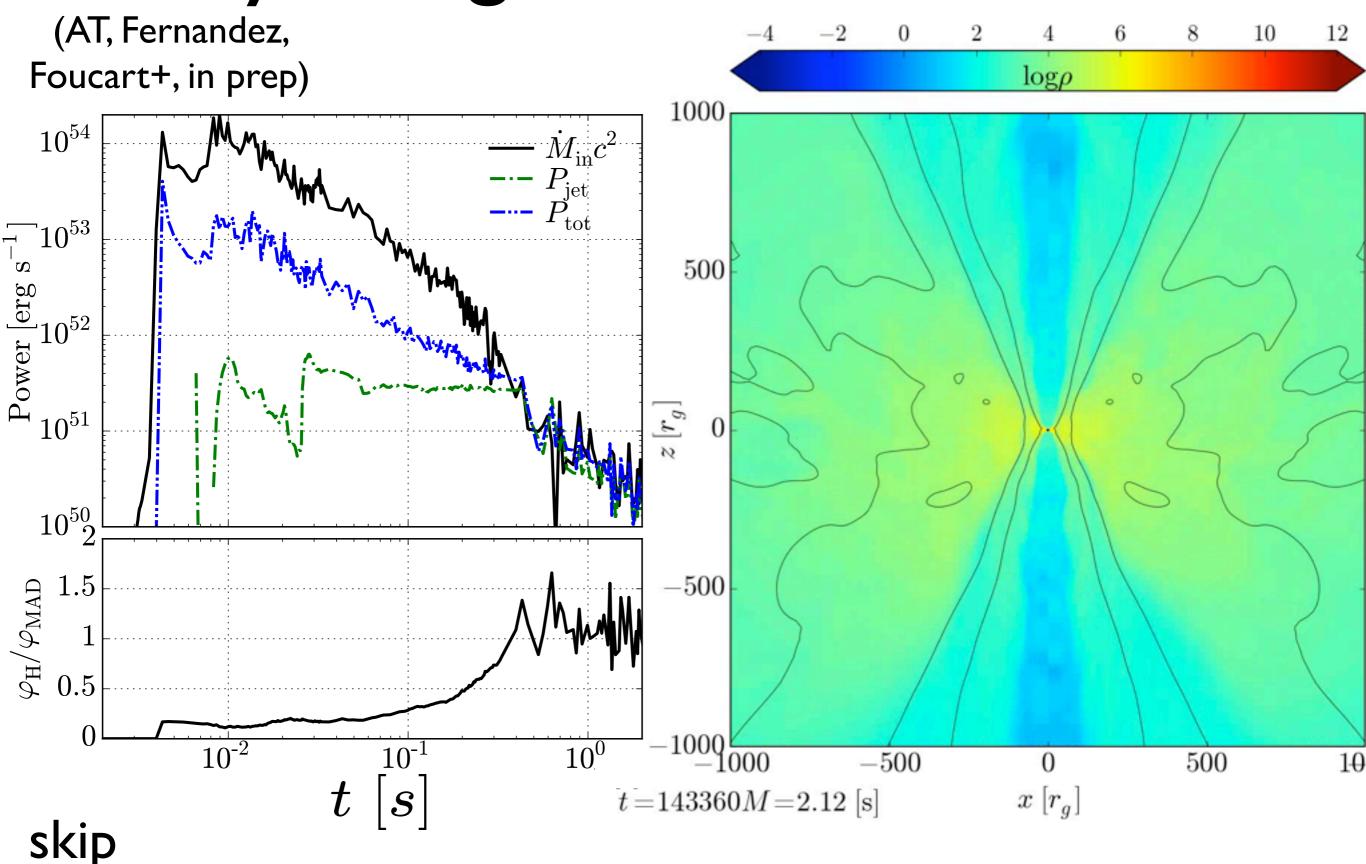


### Binary Merger Disks Gone MAD

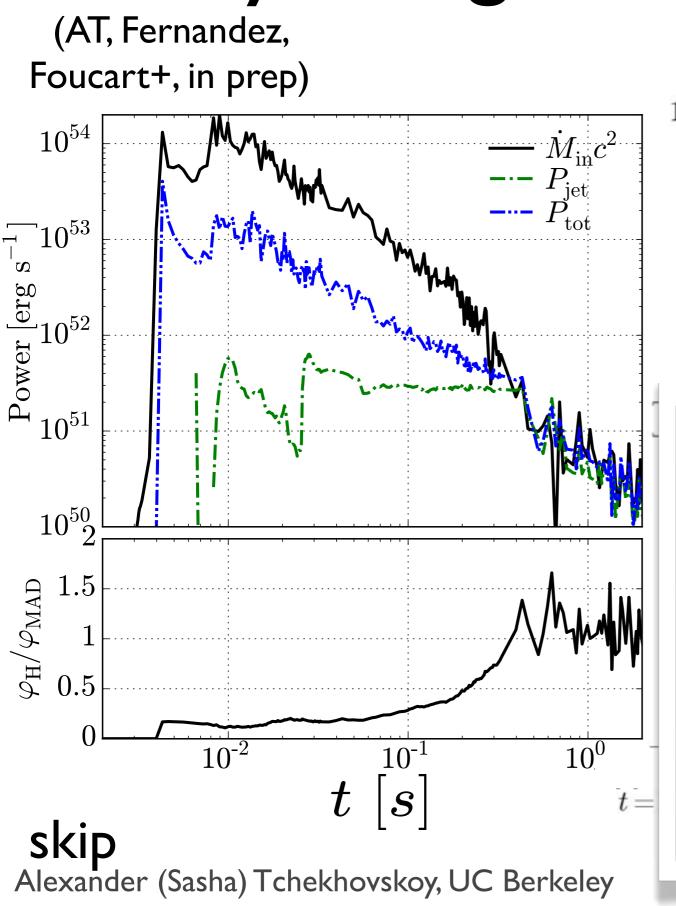


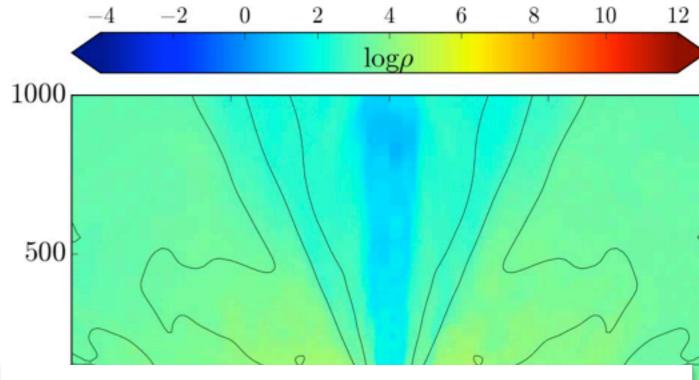
Alexander (Sasha) Tchekhovskoy, UC Berkeley

### Binary Merger Disks Gone MAD



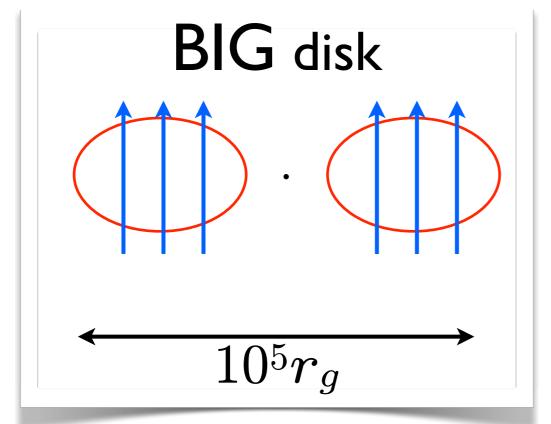
#### Binary Merger Disks Gone MAD

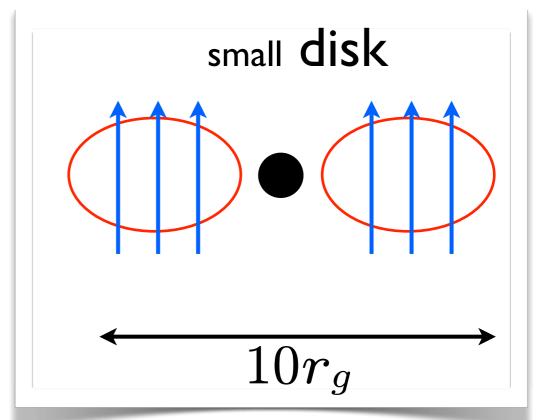


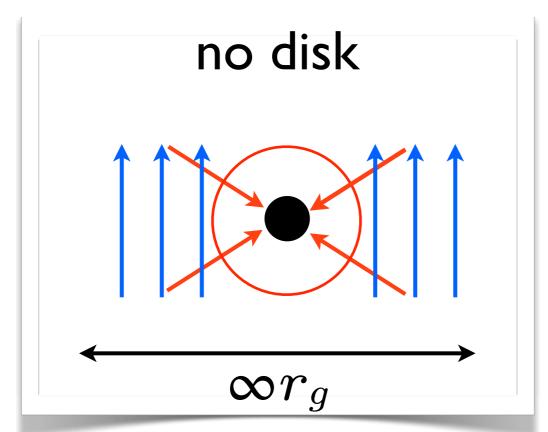


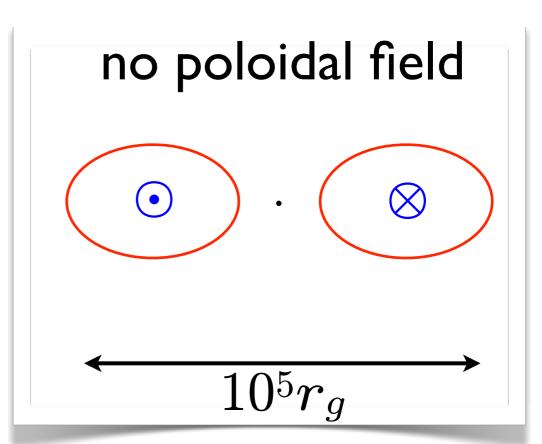
Jet power shows no trends before abruptly switching off at MAD onset (see also AT & Giannios 2015)

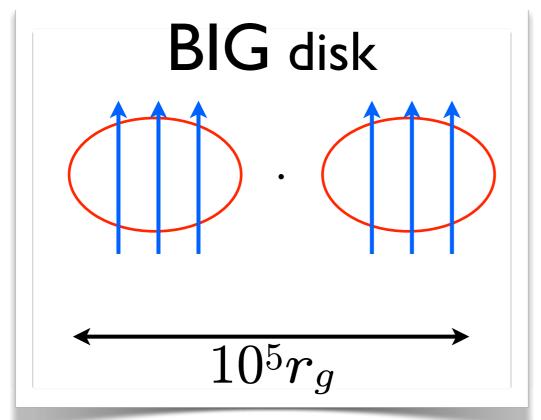
Jet opening angle ~ 0.2 rad agrees with observations (Fong+2015)

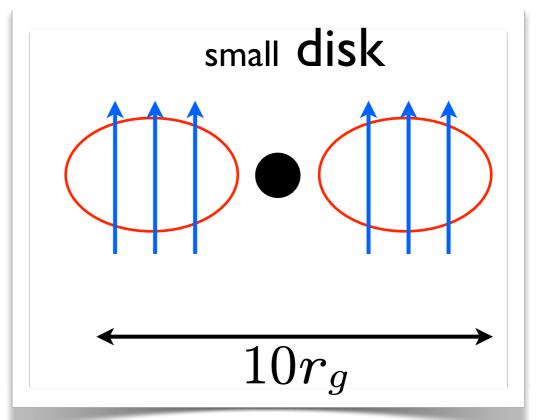


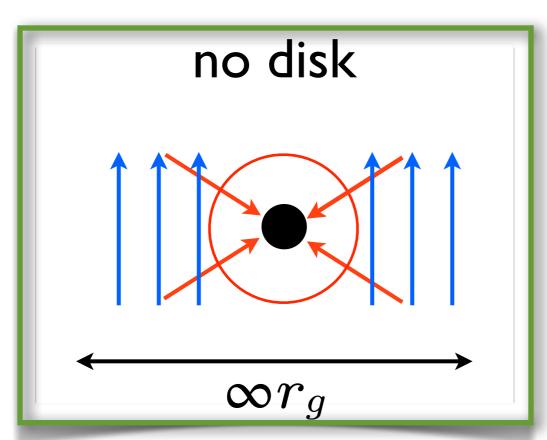


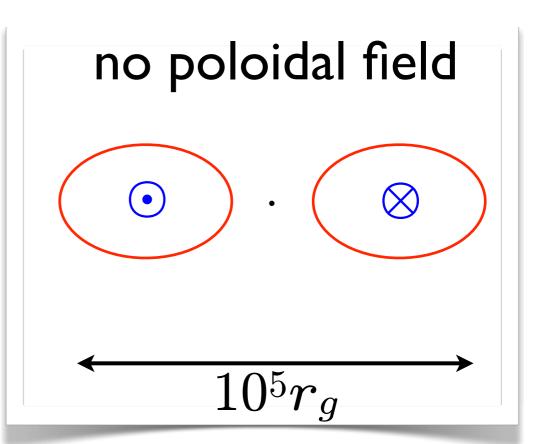




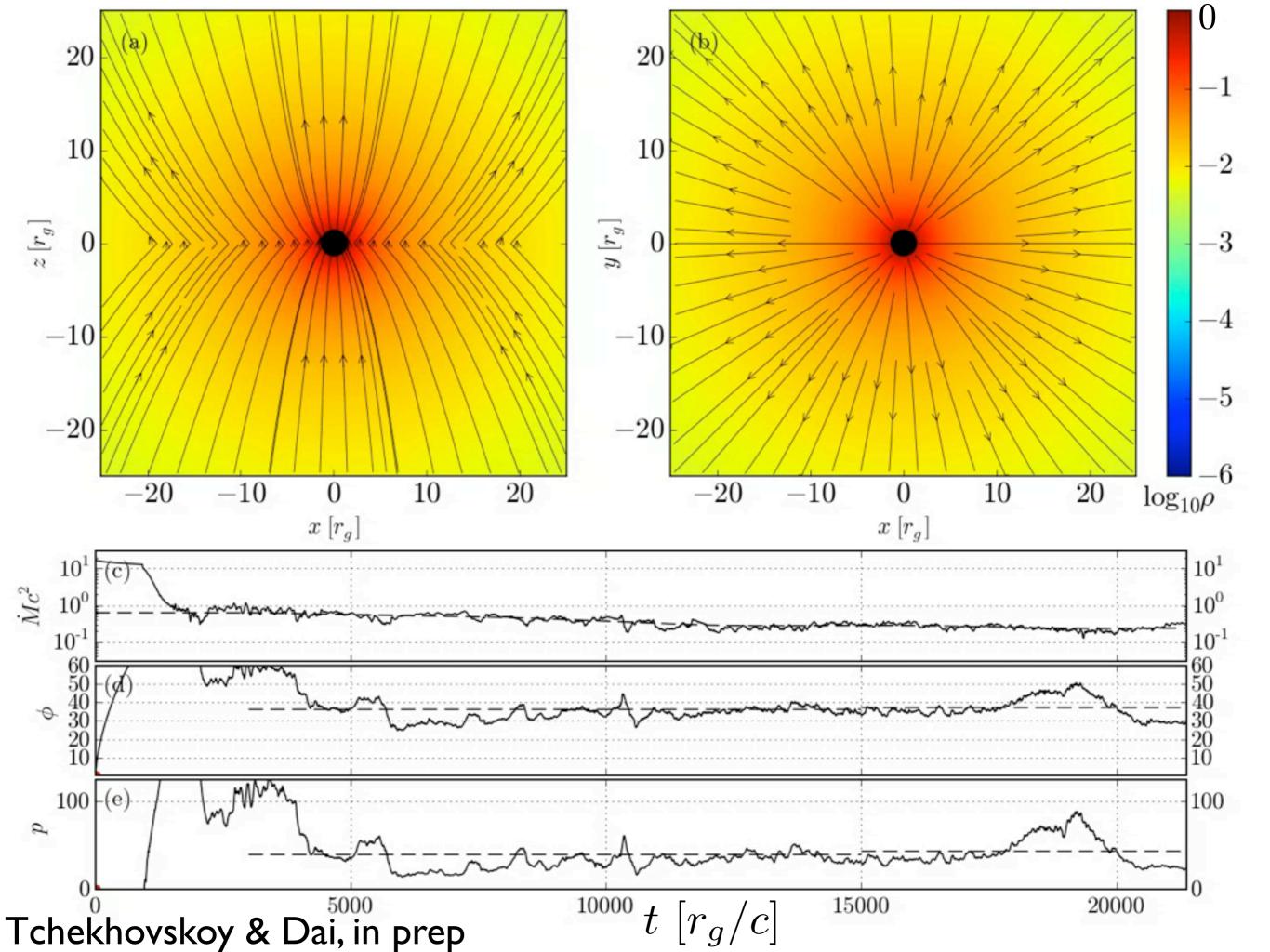


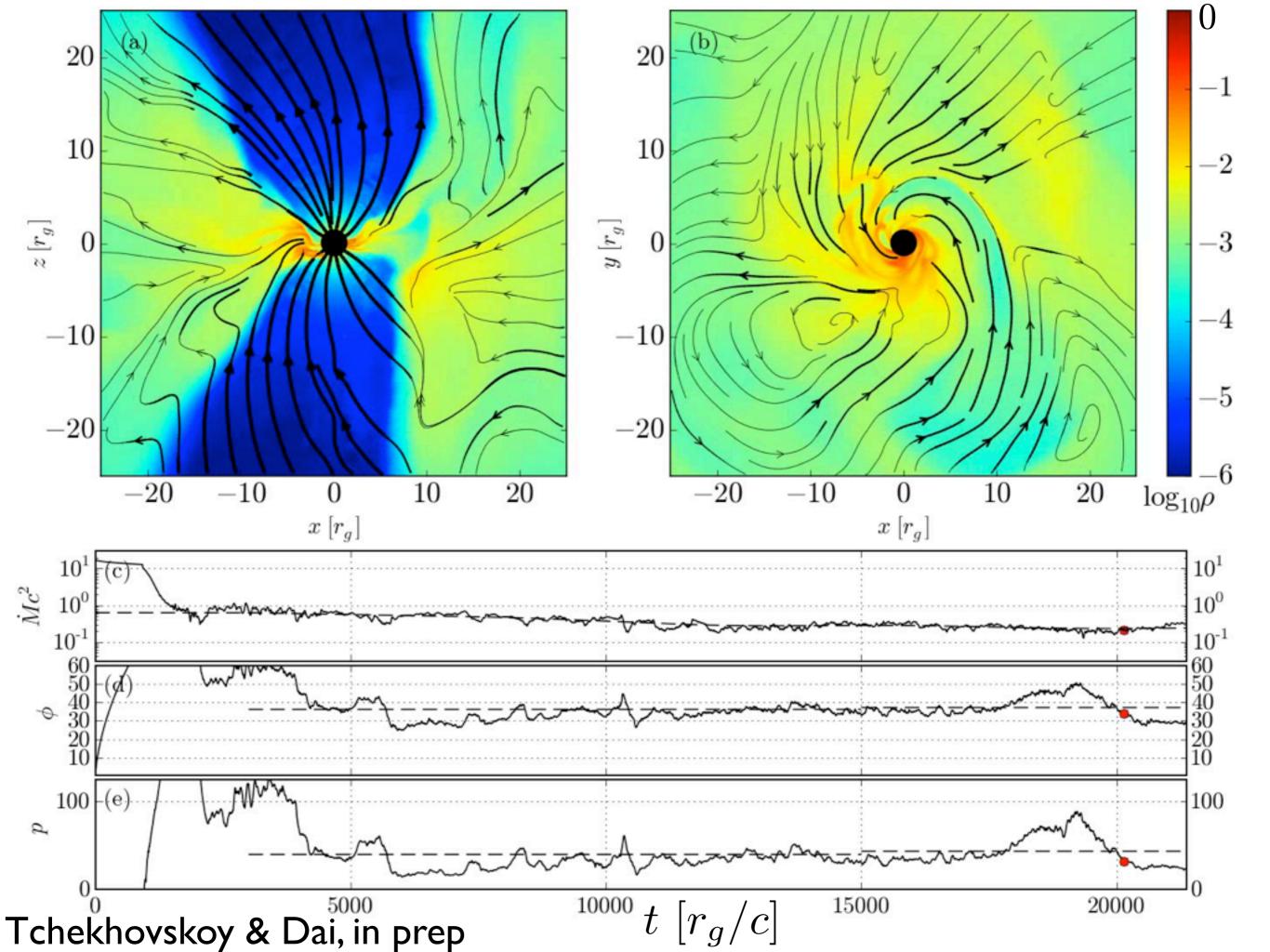


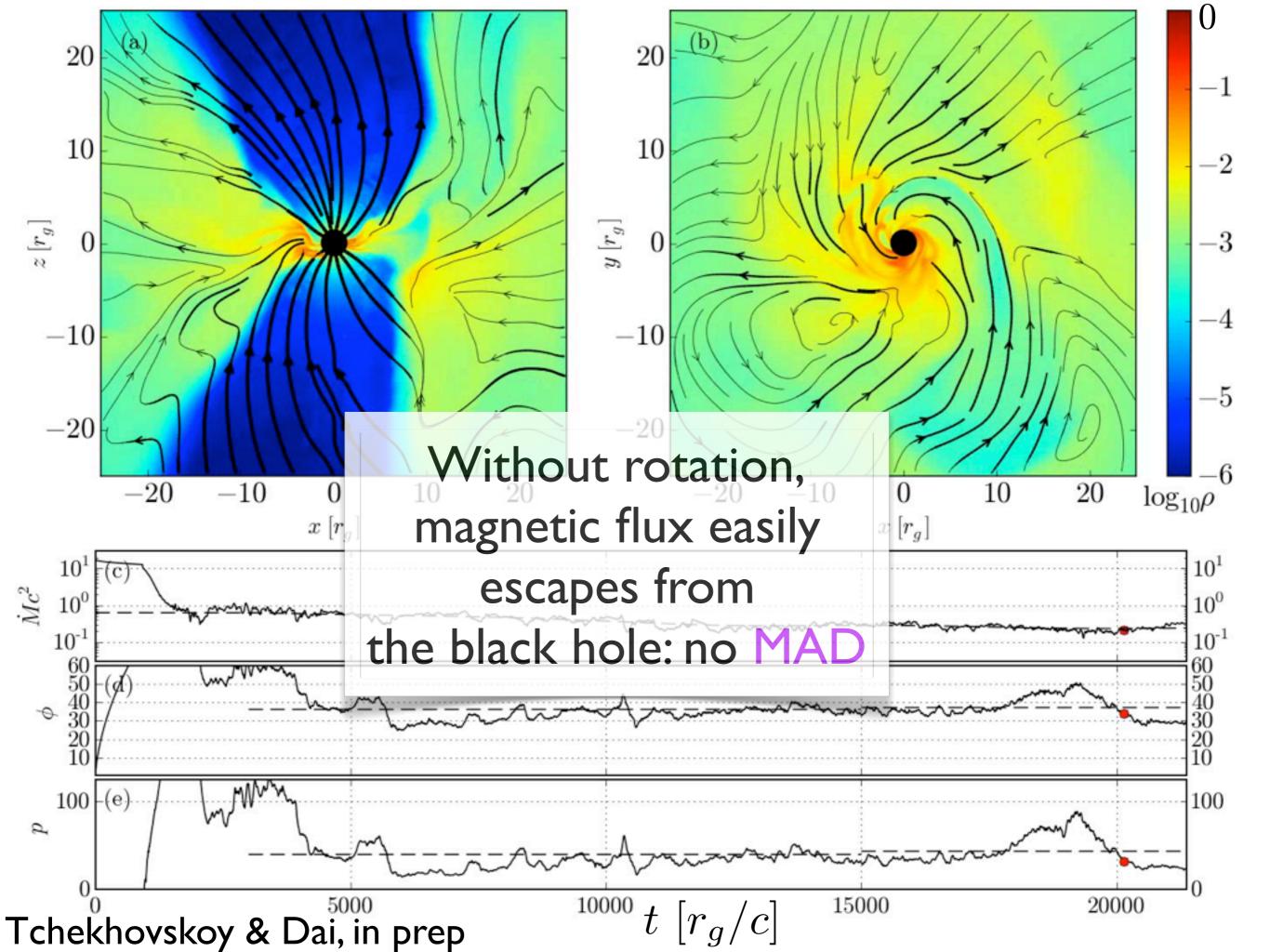




$$a=0.9$$
 spherical accretion no rotation

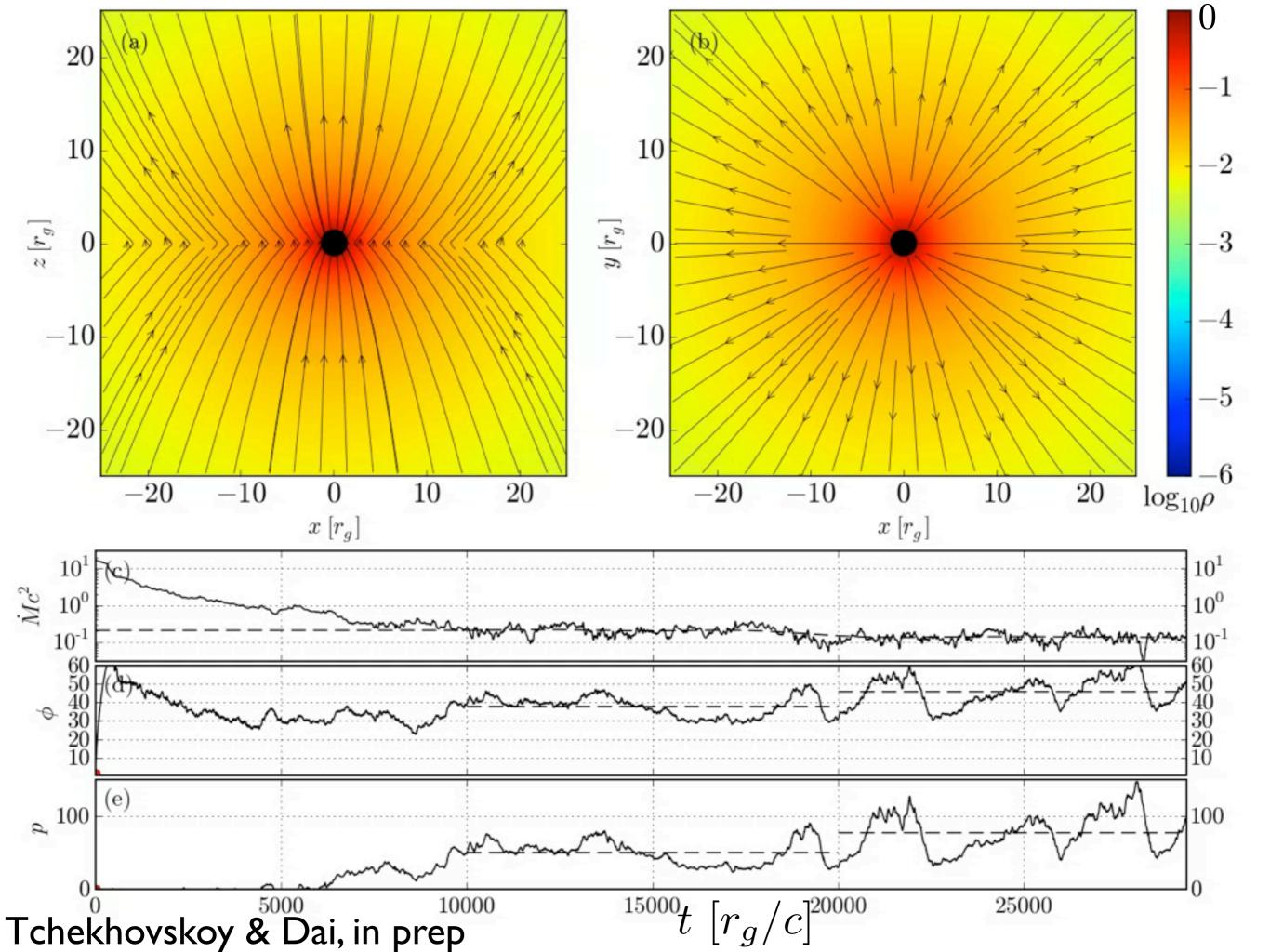


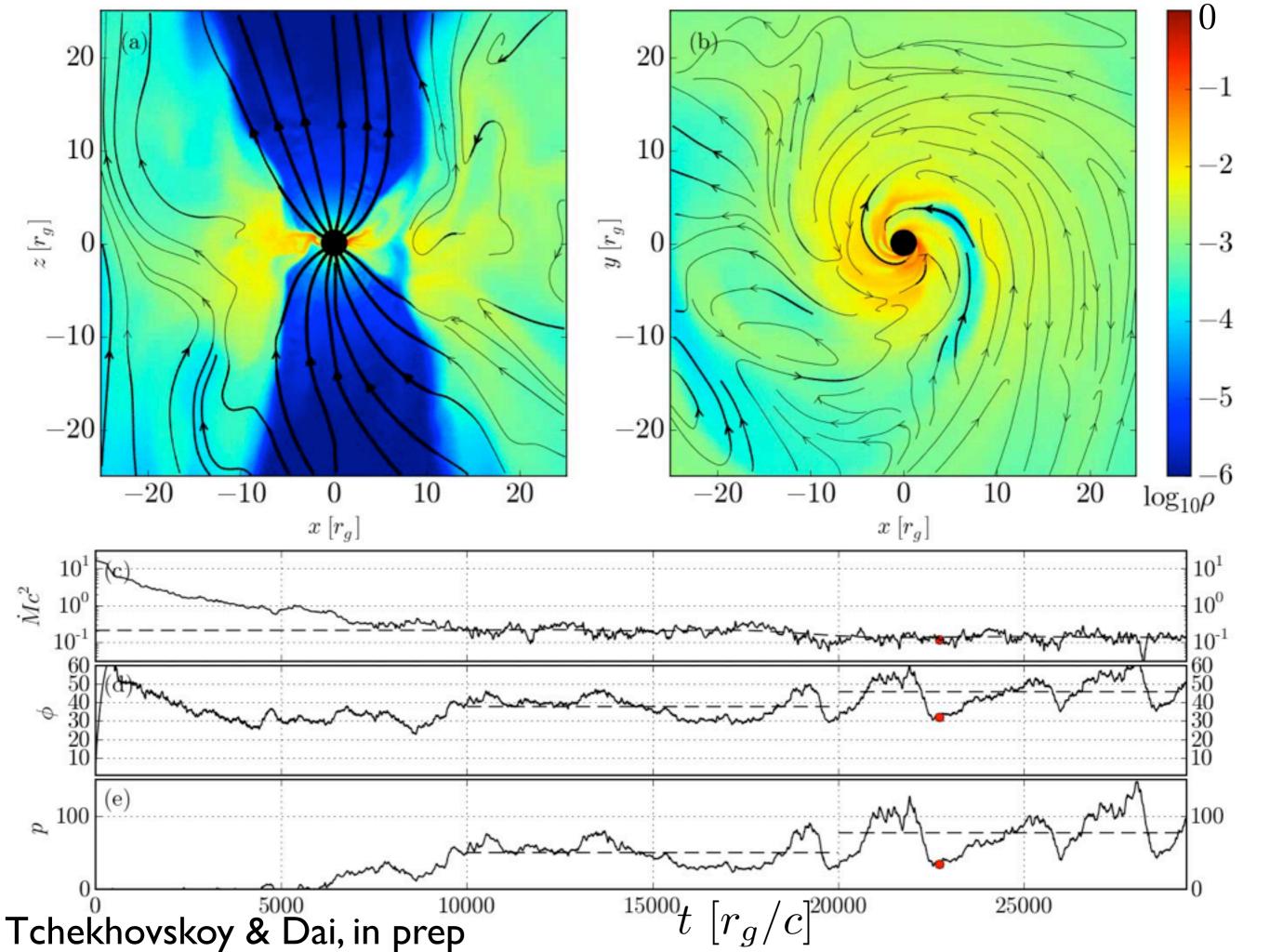


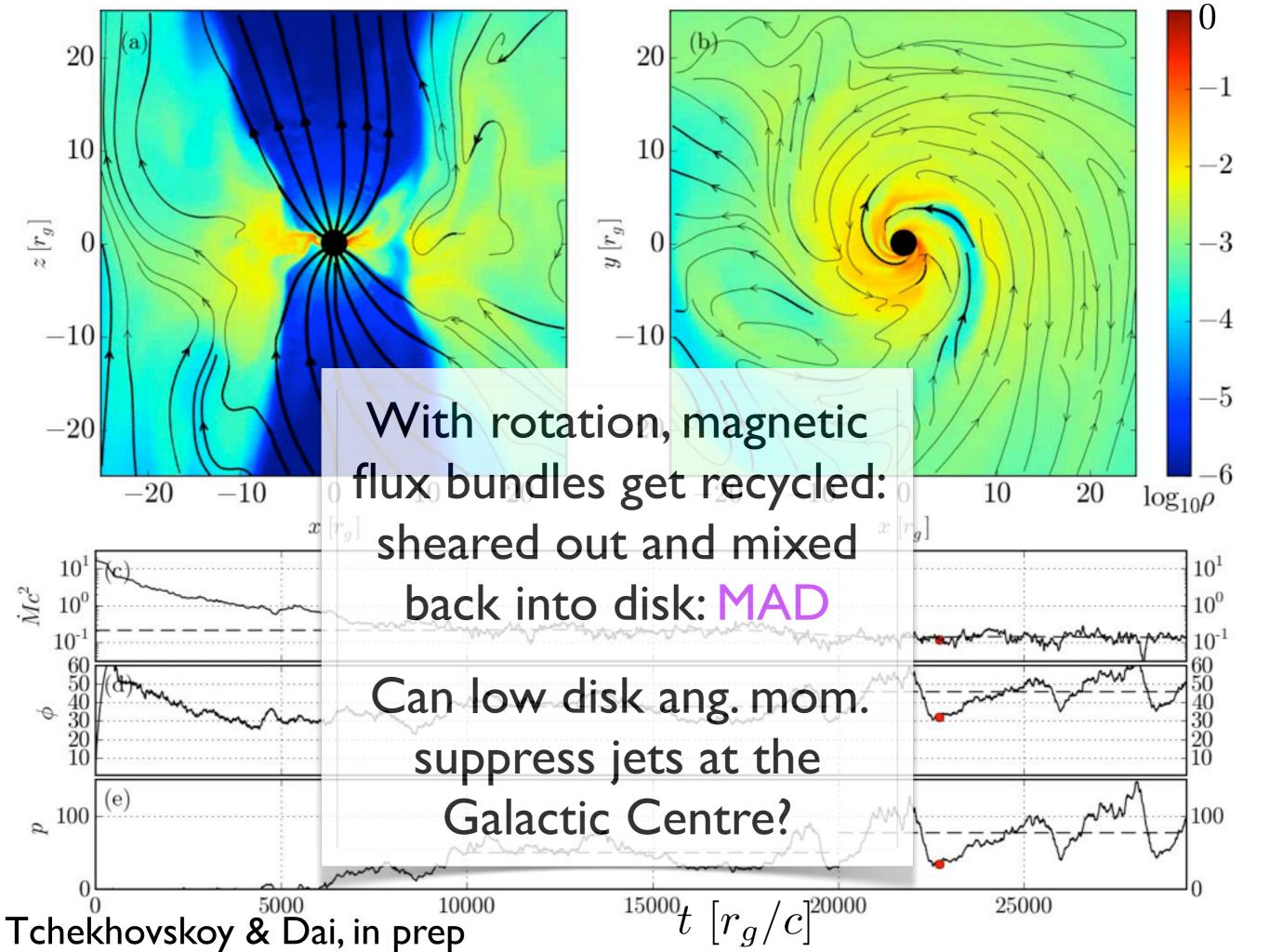


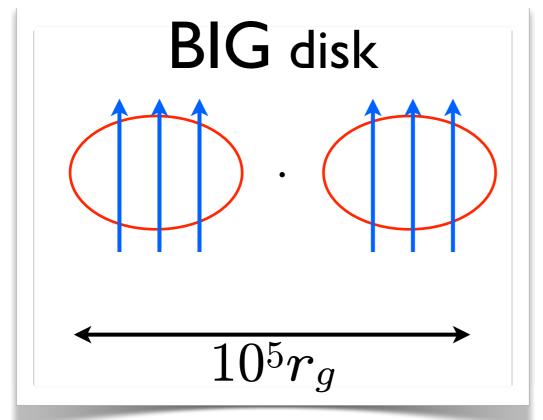
$$a=0.9$$
 spherical accretion rotation:  $R_{\rm circ}=50r_g$ 

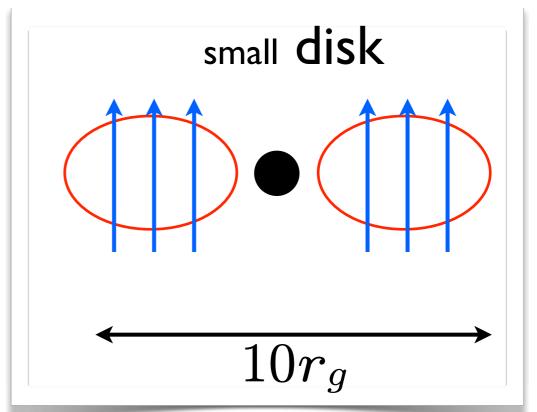
 $t \ [r_g/c]$ 

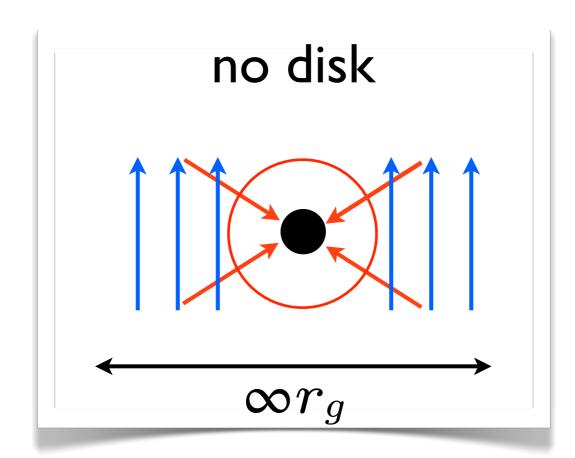


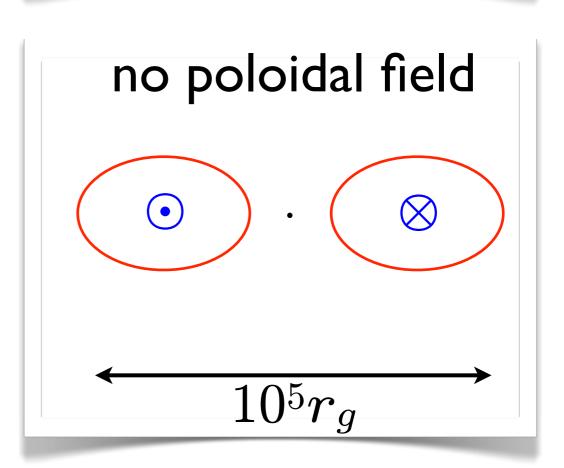


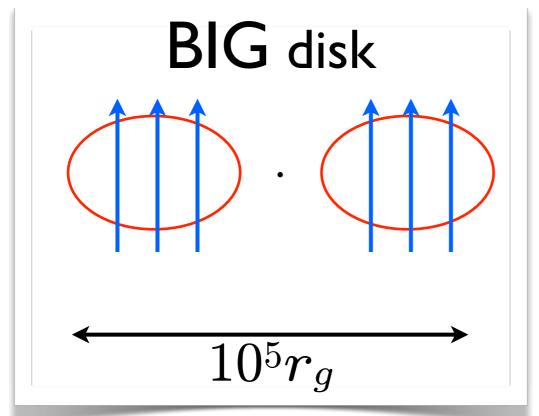


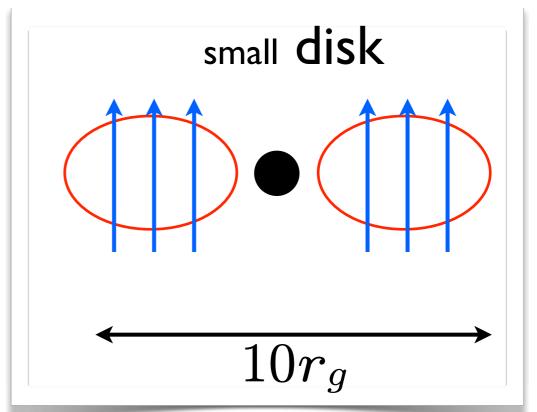


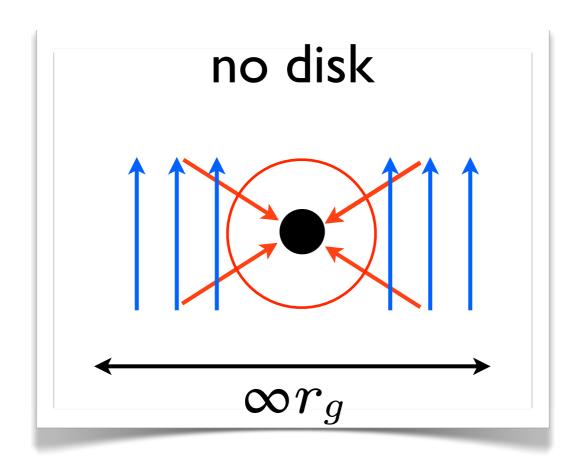


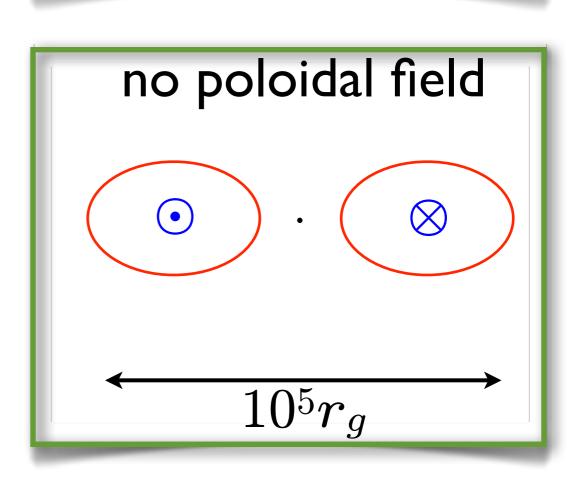












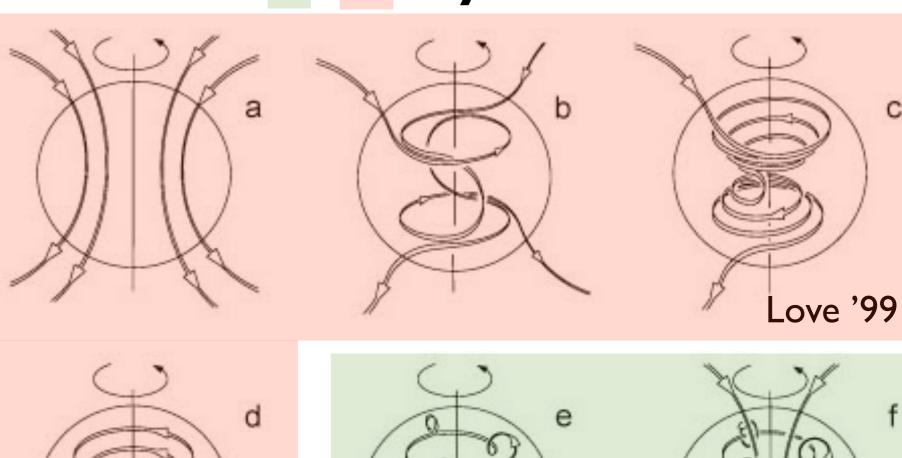
### Can Toroidal Fields Make Jets?

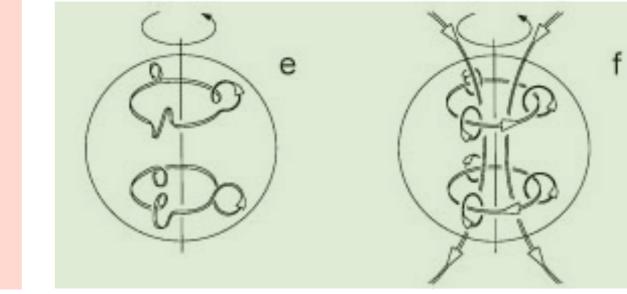
- Unlikely: healthy
  jets need
  poloidal field (e.g.,
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- Possible mechanism for jets without  $B_p$ ?
  - Large-scale α-ω dynamo (Moffatt '78; Parker '79)
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 $\alpha$ - $\omega$  dynamo:

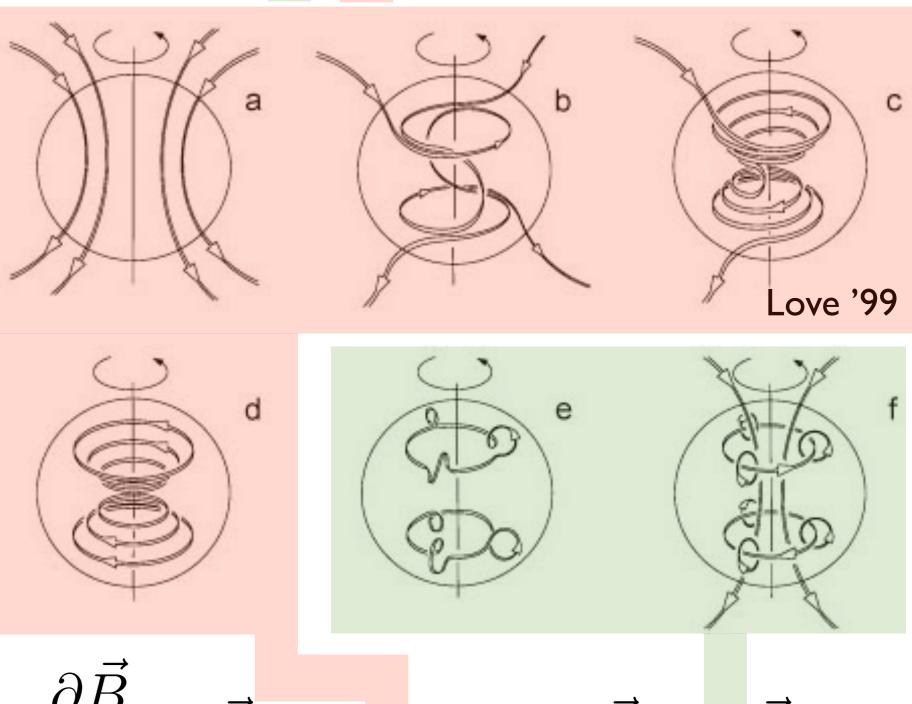




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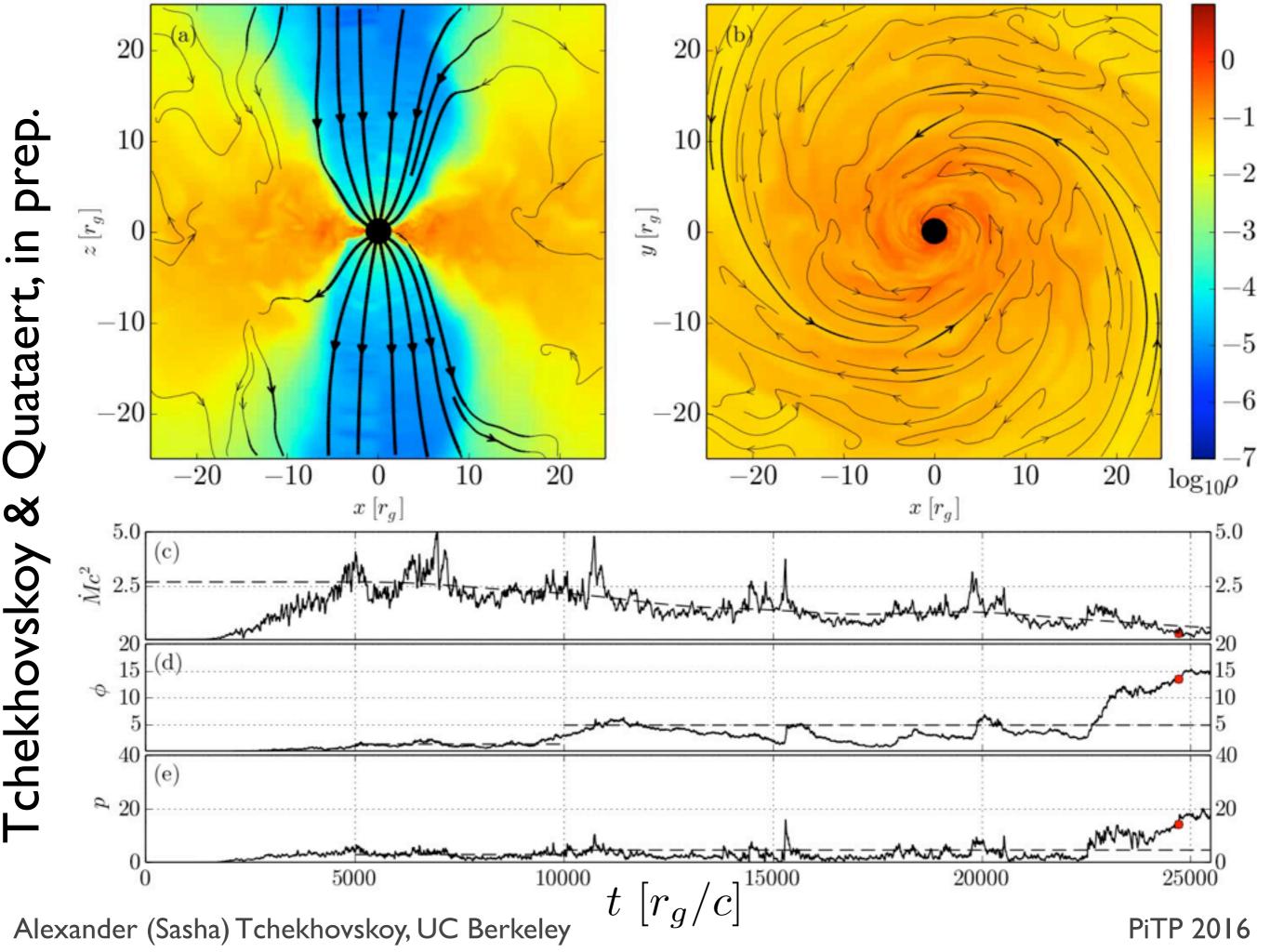
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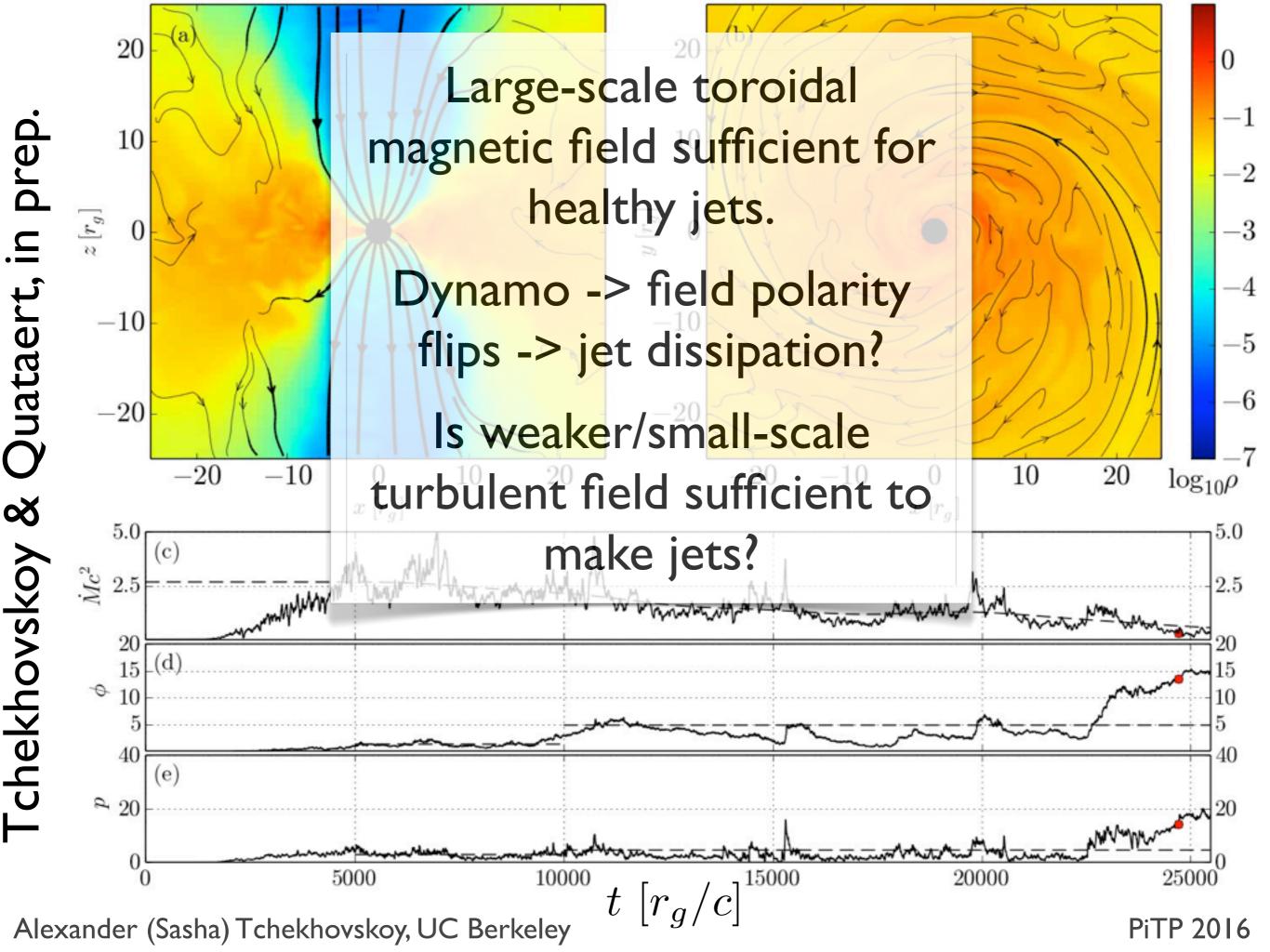
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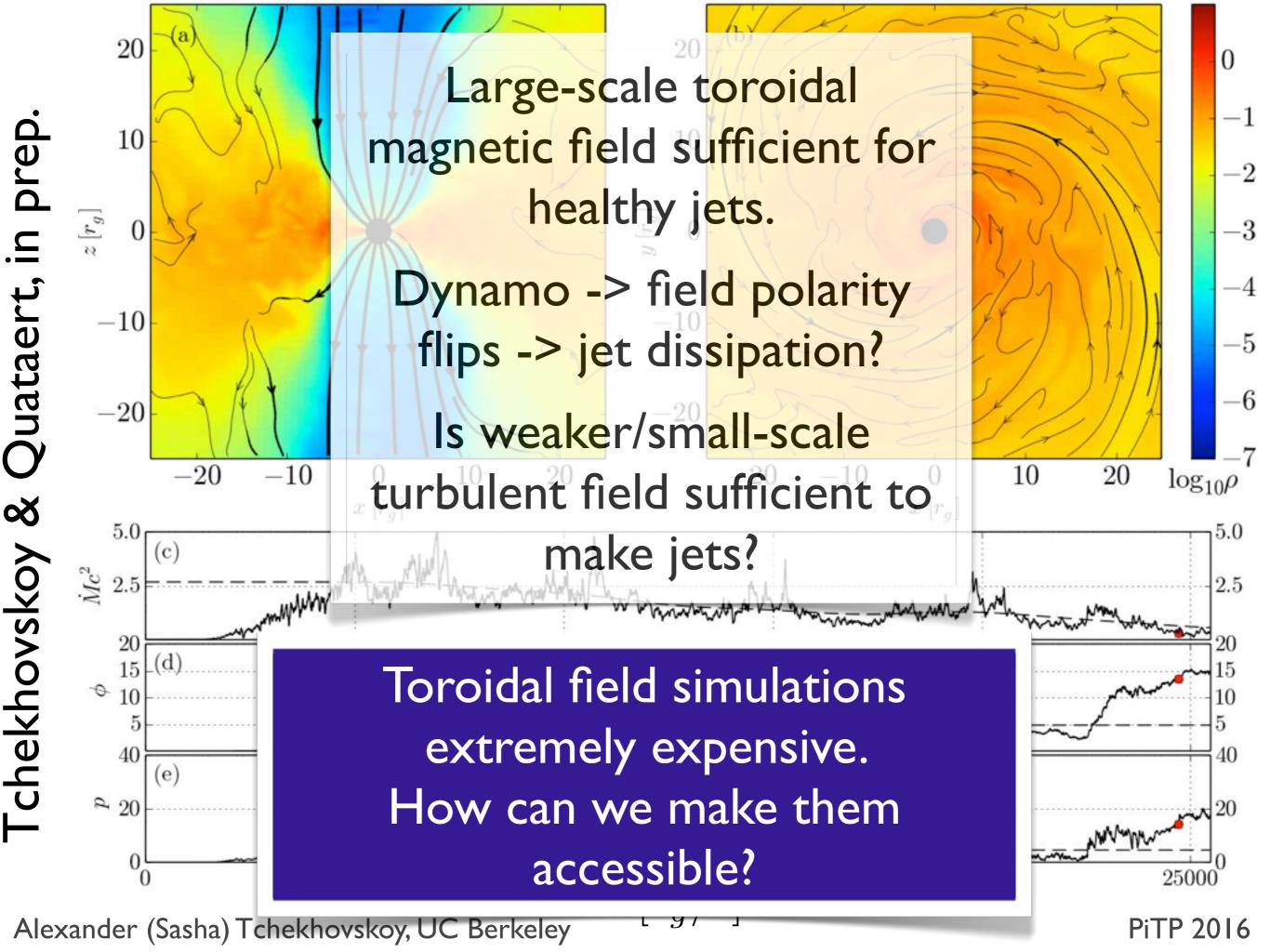


$$\frac{\partial \vec{B}}{\partial t} = \vec{\nabla} \times [(\vec{\omega} \times \vec{r}) \times \vec{B} + \alpha \vec{B}]$$
 Moffatt '78

$$a=0.9$$
 toroidal field,  $\beta=5$  large torus







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  - cutting edge of modern supercomputing



Matthew Liska (U of Amsterdam)

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- Multi-GPU 3D H-AMR ("hammer"):
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    - Excellent scaling to >= 4096 GPUs.
    - 3D, staggered fields, AMR
    - Hierarchical time stepping
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    - Stanford XStream (production now)
    - ORNL Summit (production in '18)



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    - Stanford XStream (production now)
    - ORNL Summit (production in '18)
- Whole slew of important applications:
  - Long-term disk evolution
  - Tilted thin disks
  - Etc.



Matthew Liska (U of Amsterdam)

#### When are Jets Produced?

#### MADs:

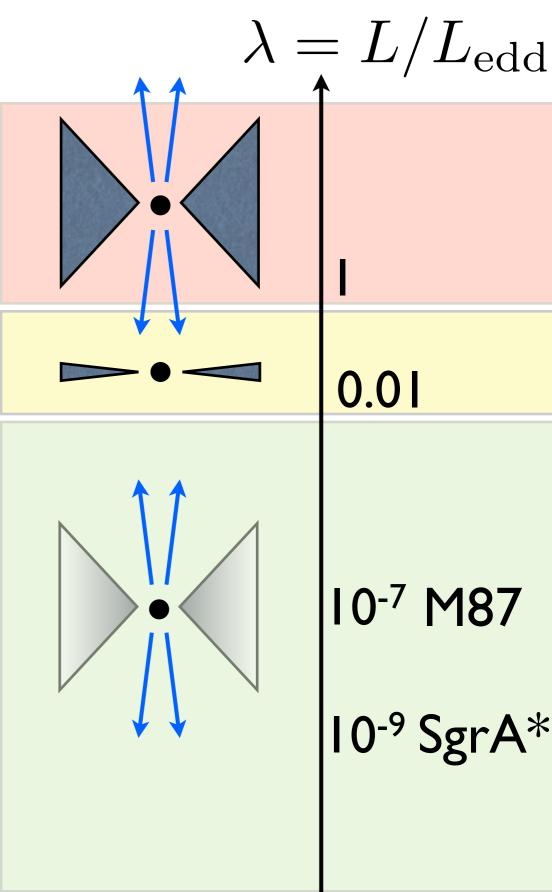
(AT+13, Tidal disruptions (TDEs), AT & Giannios 15) ultra-luminous X-ray sources, gamma-ray bursts

(Zamaninasab ++AT 14, Ghisellini+14)

Quasars, X-ray binaries, TDEs

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Low-luminosity active galactic nuclei (LLAGN), X-ray binaries



Alexander (Sasha) Tchekhovskoy, UC Berkeley

PiTP 2016

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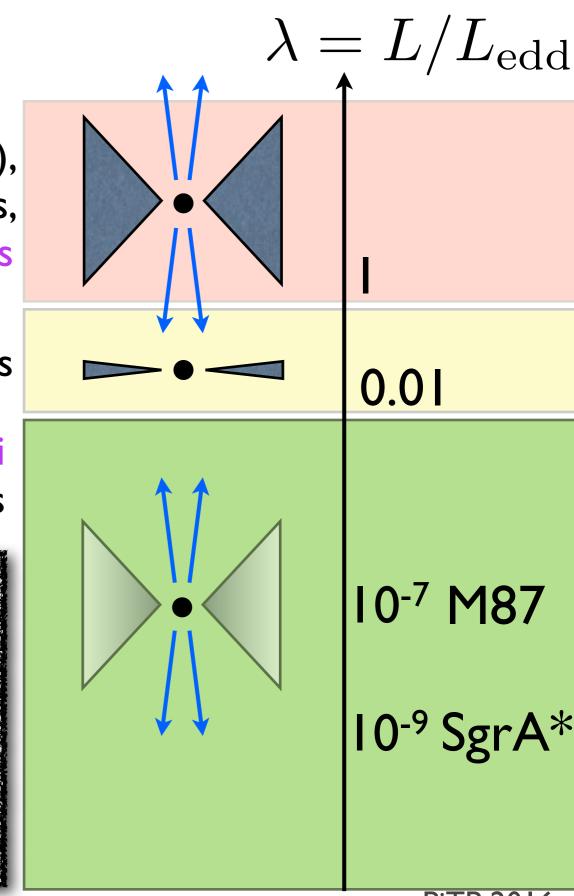
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Low-luminosity active galactic nuclei (LLAGN), X-ray binaries

Disk radiative properties are most uncertain at low luminosities.

In the next several years, *EHT* will resolve the shadows of two black holes that accrete in this regime.

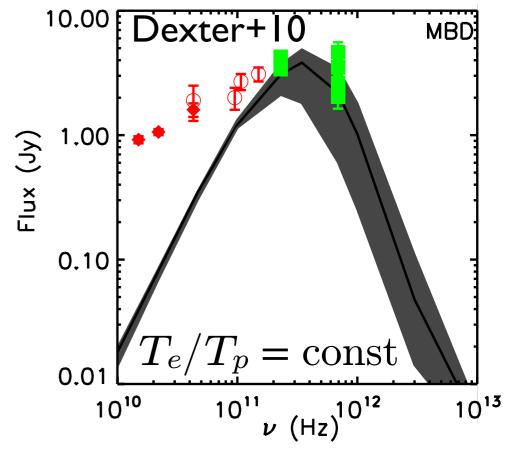


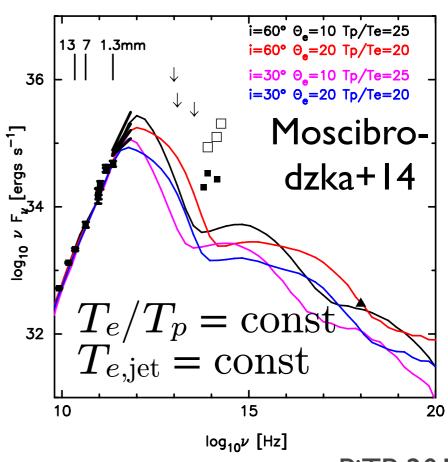
Alexander (Sasha) Tchekhovskoy, UC Berkeley

PiTP 2016

#### Electron Micro-physics is Key to SgrA\* Observations

- Plasma is collisionless, so electron and proton temperatures decouple
  - but,  $T_e (\neq T_p)$  is poorly known!
- Dissipation predominantly heats protons, whereas electrons radiate
- So,  $T_e$  is usually "painted" on top of simulations:
  - Usual assumption (eg Dexter+10):  $T_e/T_p = \mathrm{const.} < 1$
  - To reproduce flat radio spectrum, need to "paint" polar regions with hot  $T_e{=}10^{11}~{\rm K}$  electrons (Moscibrodzka et al. 2014)
- Is there a way to eliminate the free function,  $T_e(T_p, ...)$ ?





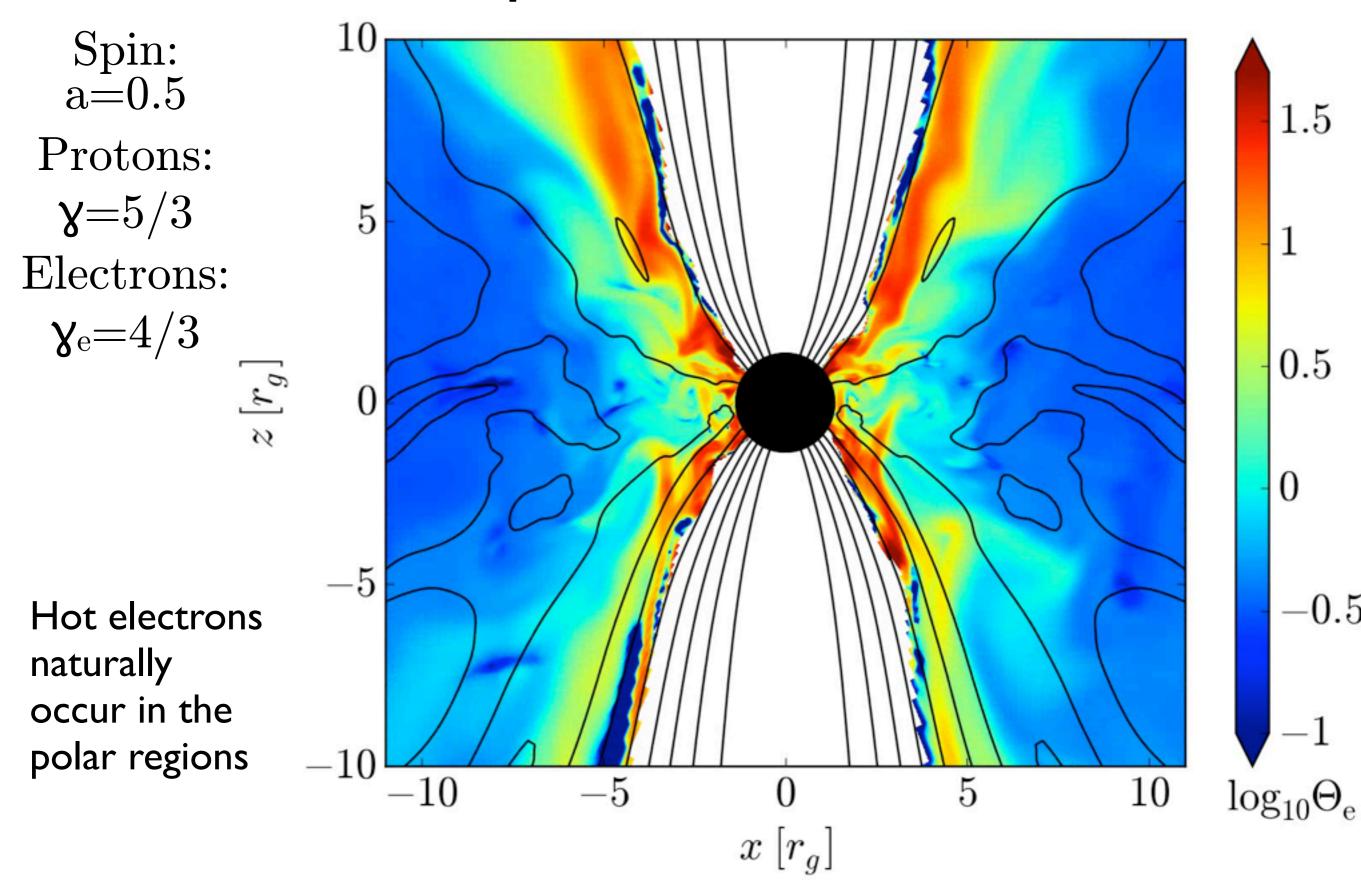
#### Electron Micro-physics is Key to SgrA\* Observations

- Our improved approach:
  - Evolve electrons as a second fluid
  - Electrons receive a fraction  $f_e(T_e, T_p, \beta)$  of dissipated heat (Howes 2010)
    - stronger heating in highly magnetized regions
  - Include thermal conduction along field lines (Chandra et al. 2015)
  - Neglect back-reaction of electrons on the flow
- Simulations with e-HARMPI, new parallel, 3D general relativistic MHD code that includes electrons as a separate fluid (Ressler, AT et al., 2015, 2016)



Sean Ressler (UC Berkeley)

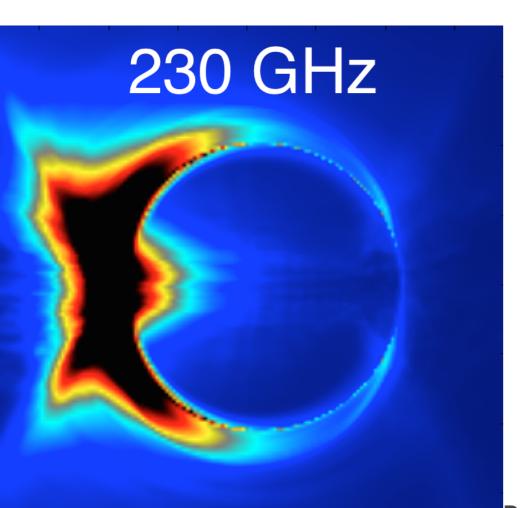
#### Electron Temperature in Simulations

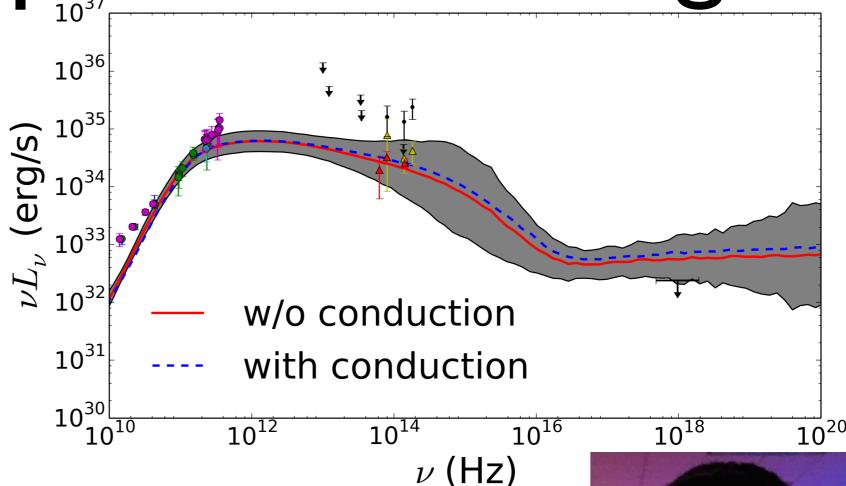


Alexander (Sasha) Tchekhovskoy, UC Berkelsyler, AT et al, 2016, in prep)

Predicted Spectra and Images

Electron micro-physics naturally reproduces broadband spectrum and high-E variability but falls short at very low frequencies. More to come!





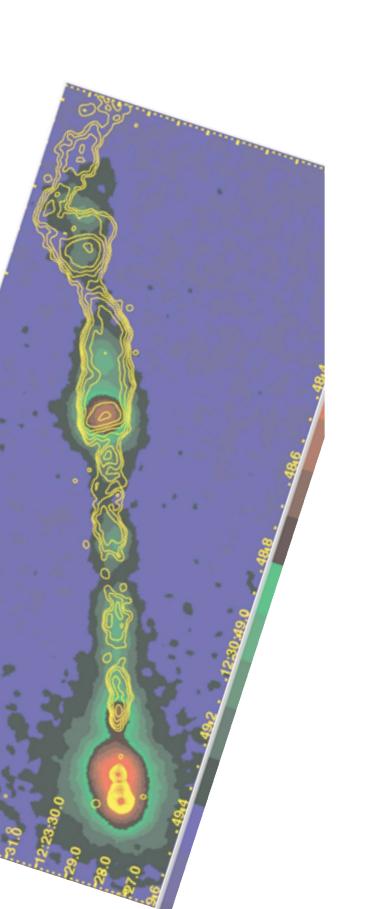
Predicted EHT mmwavelength image of SgrA\* black hole shadow

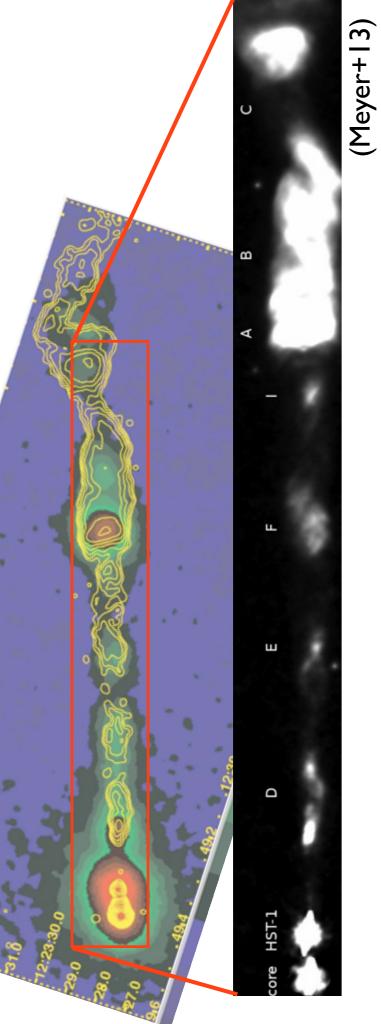
(Ressler, AT et al, 2016, in prep)

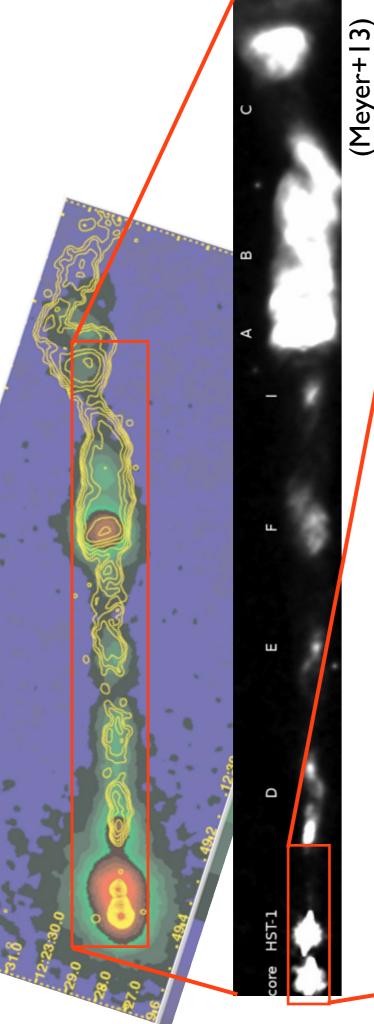


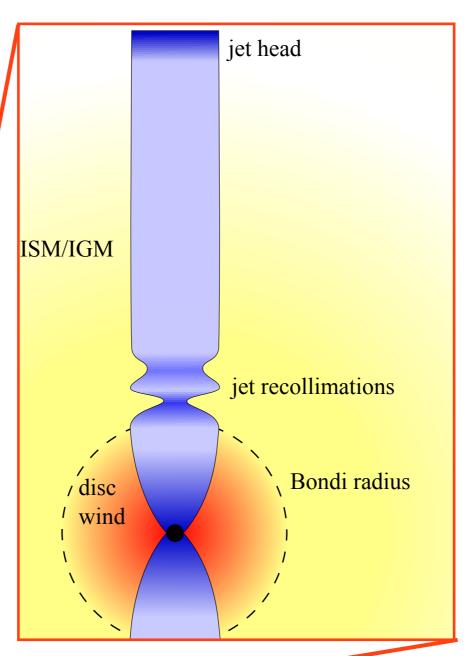
Sean Ressler (UC Berkeley)

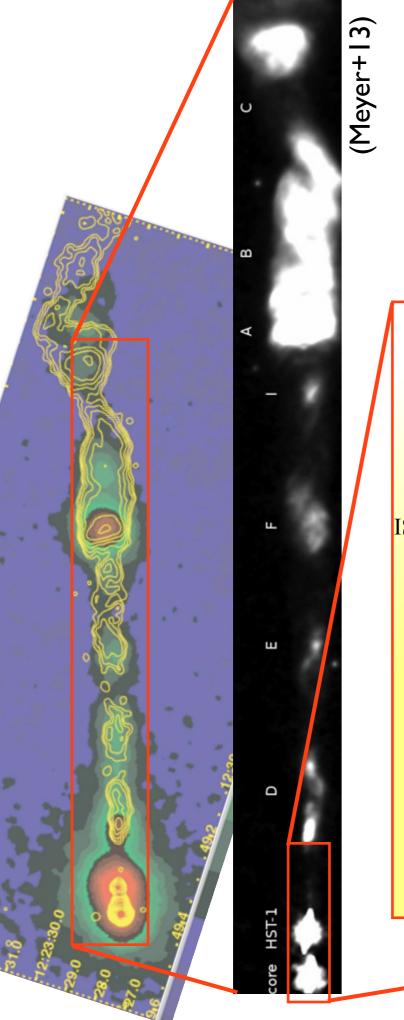
Berkeley PiTP 2016







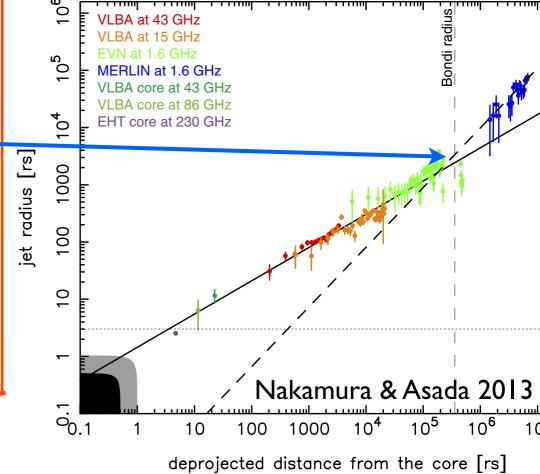




jet head ISM/IGM jet recollimations Bondi radius / disc wind

Are we seeing jet-ISM interaction?

(see also Meier 2012)



### Internal Kink Makes Jets Hot

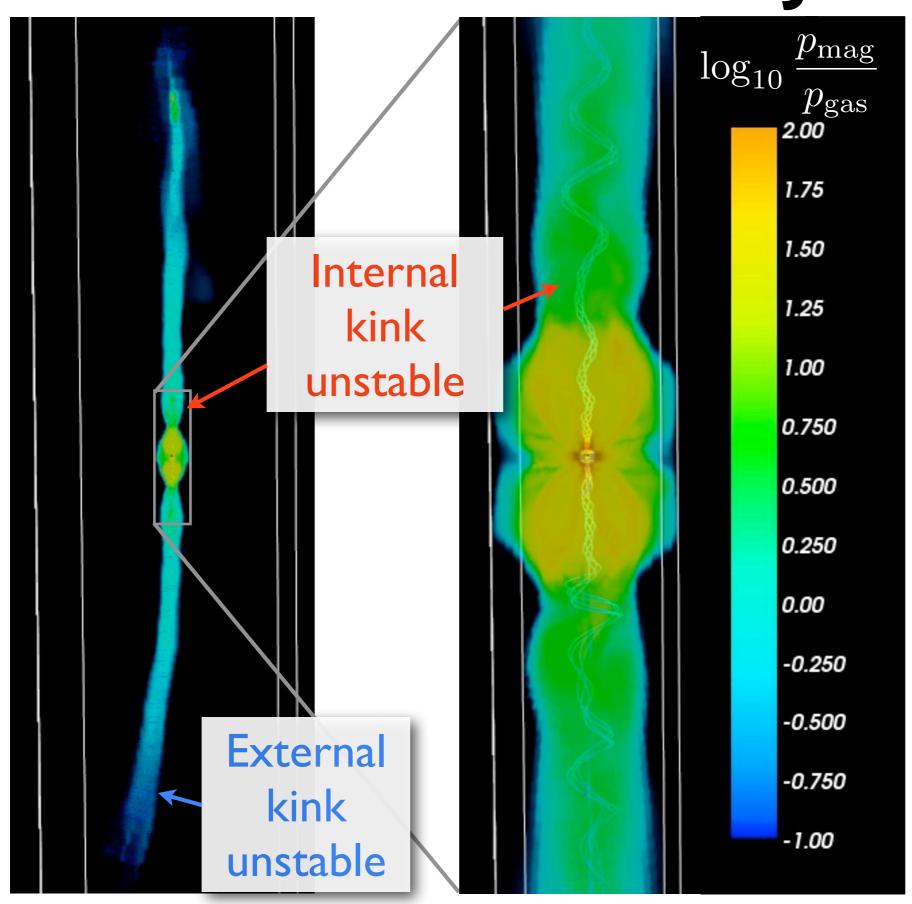
External kink unstable

rg and Tchekhovskoy, 2016, S, 456, 1739; figures/movies courtesy Bromberg

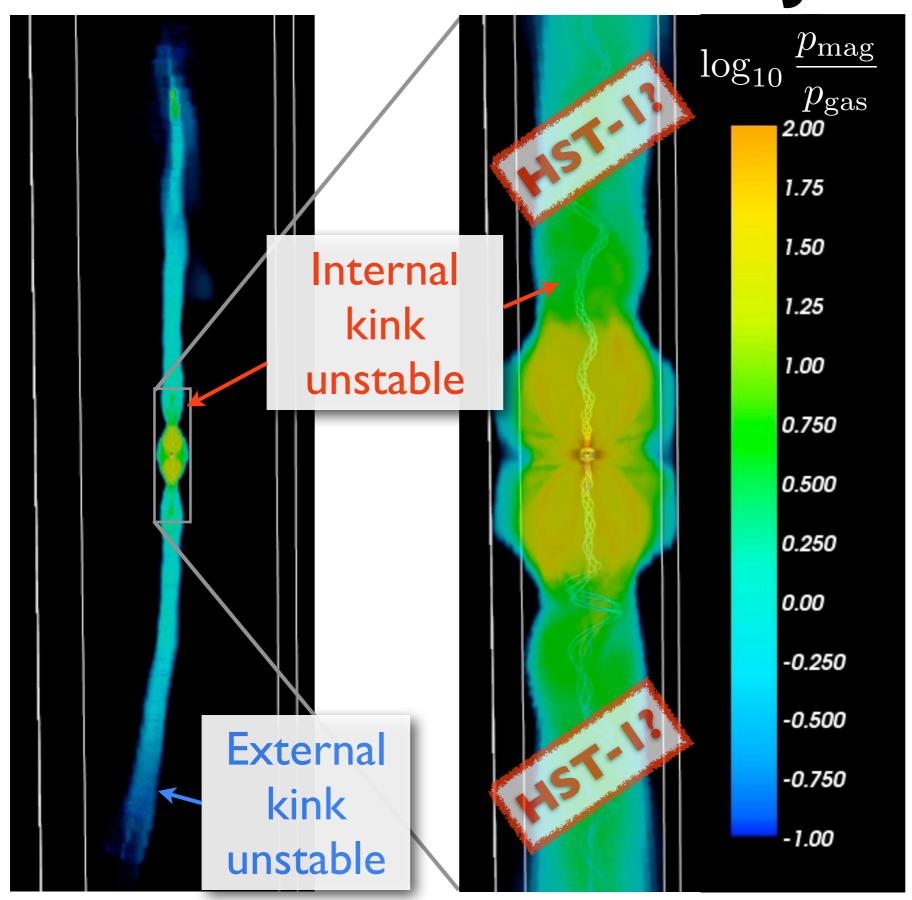
Bromberg and MNRAS, 456,

# Bromberg and Tchekhovskoy, 2016, MNRAS, 456, 1739; figures/movies courtesy Bromberg

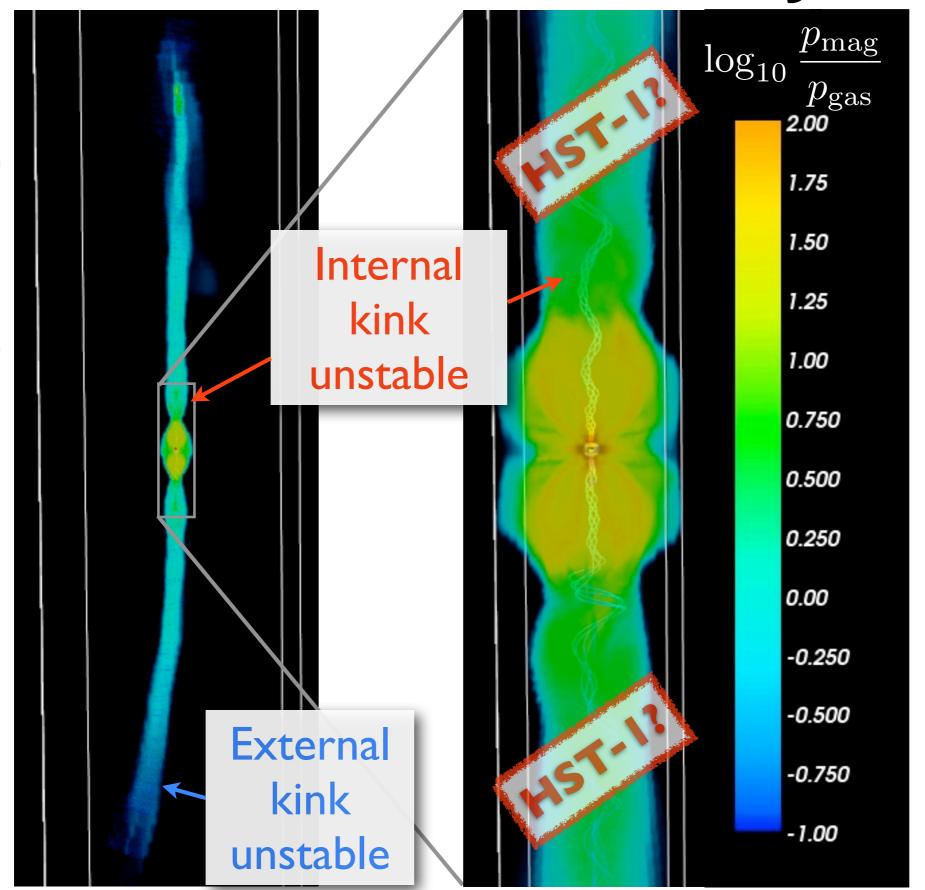
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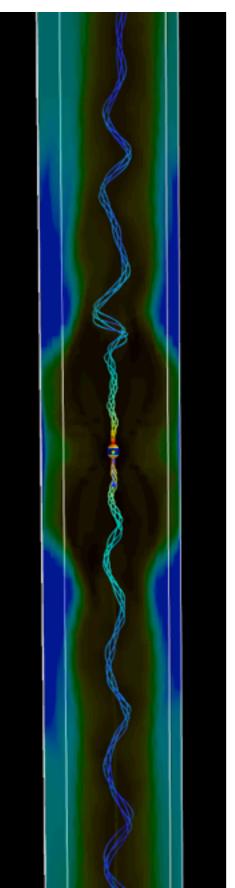


## Internal Kink Makes Jets Hot



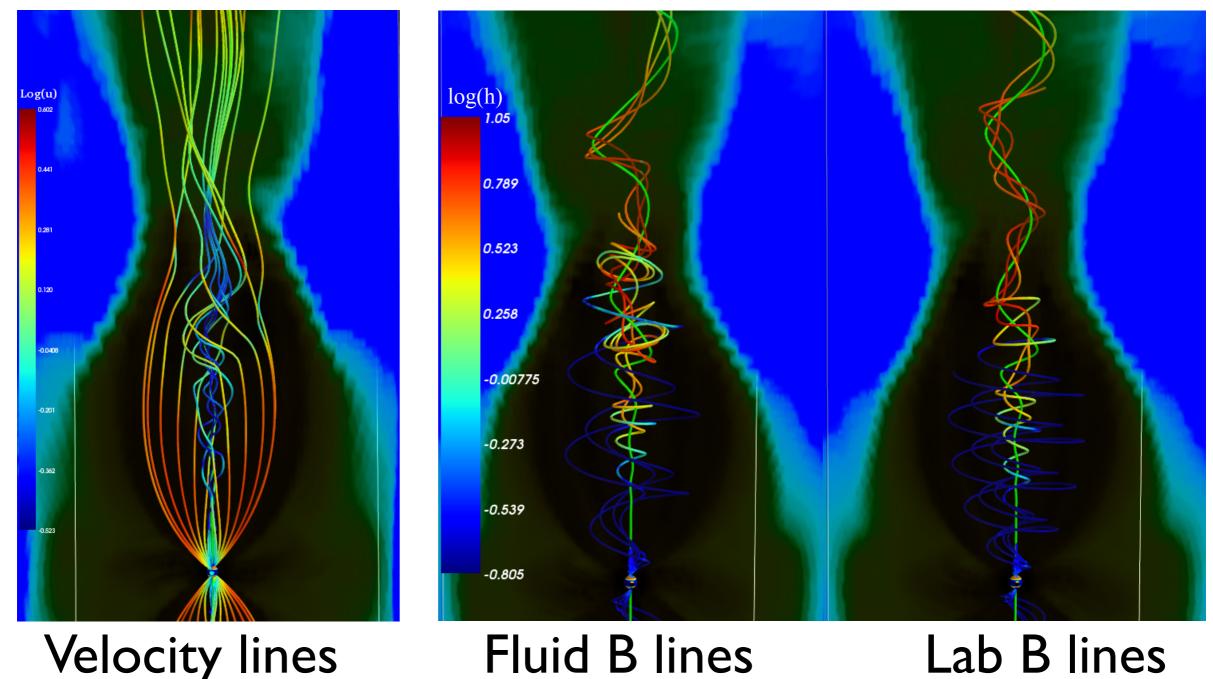
## Internal Kink Makes Jets Hot





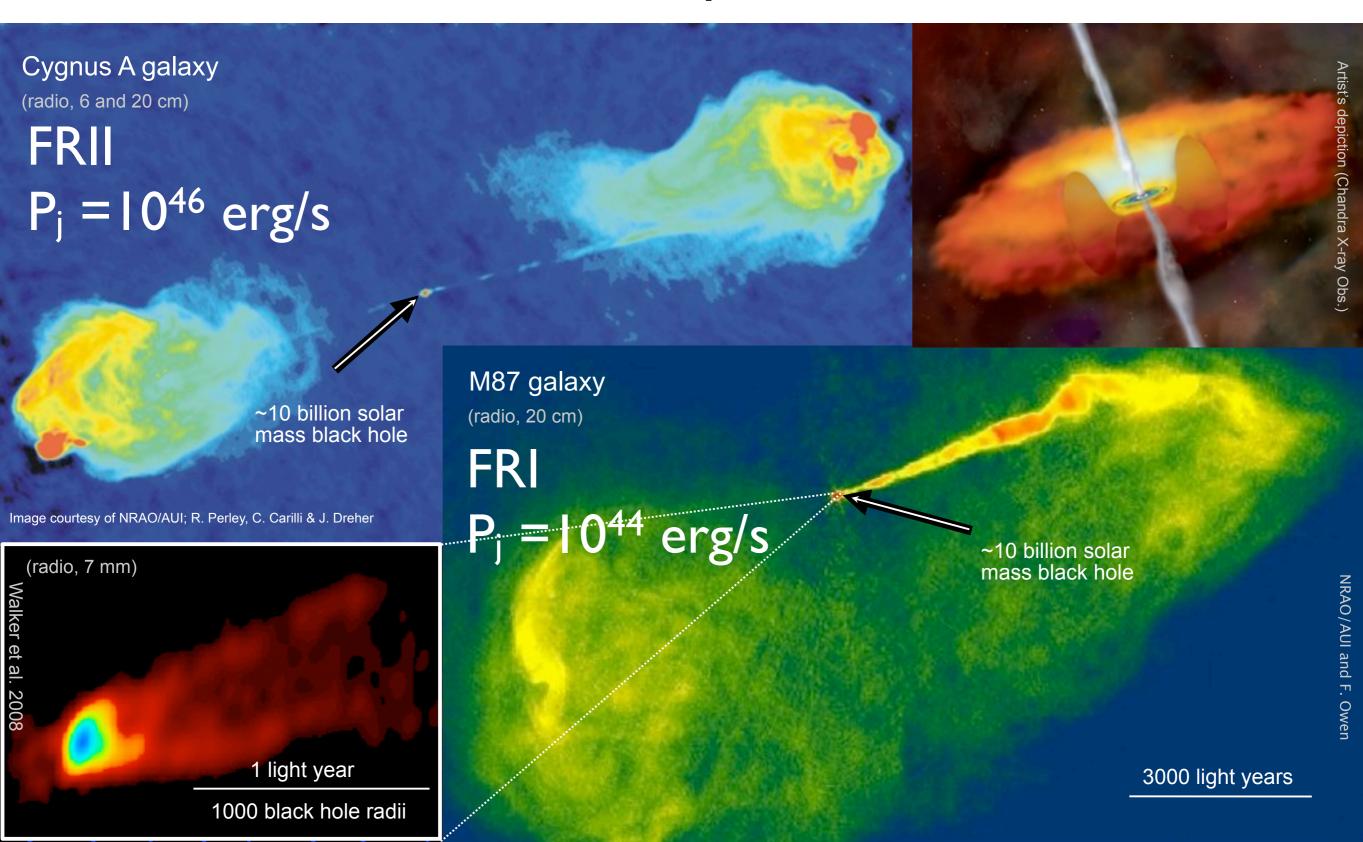
(see also Nakamura+07,08; O'Neill+12; Porth & Komissarov 14)

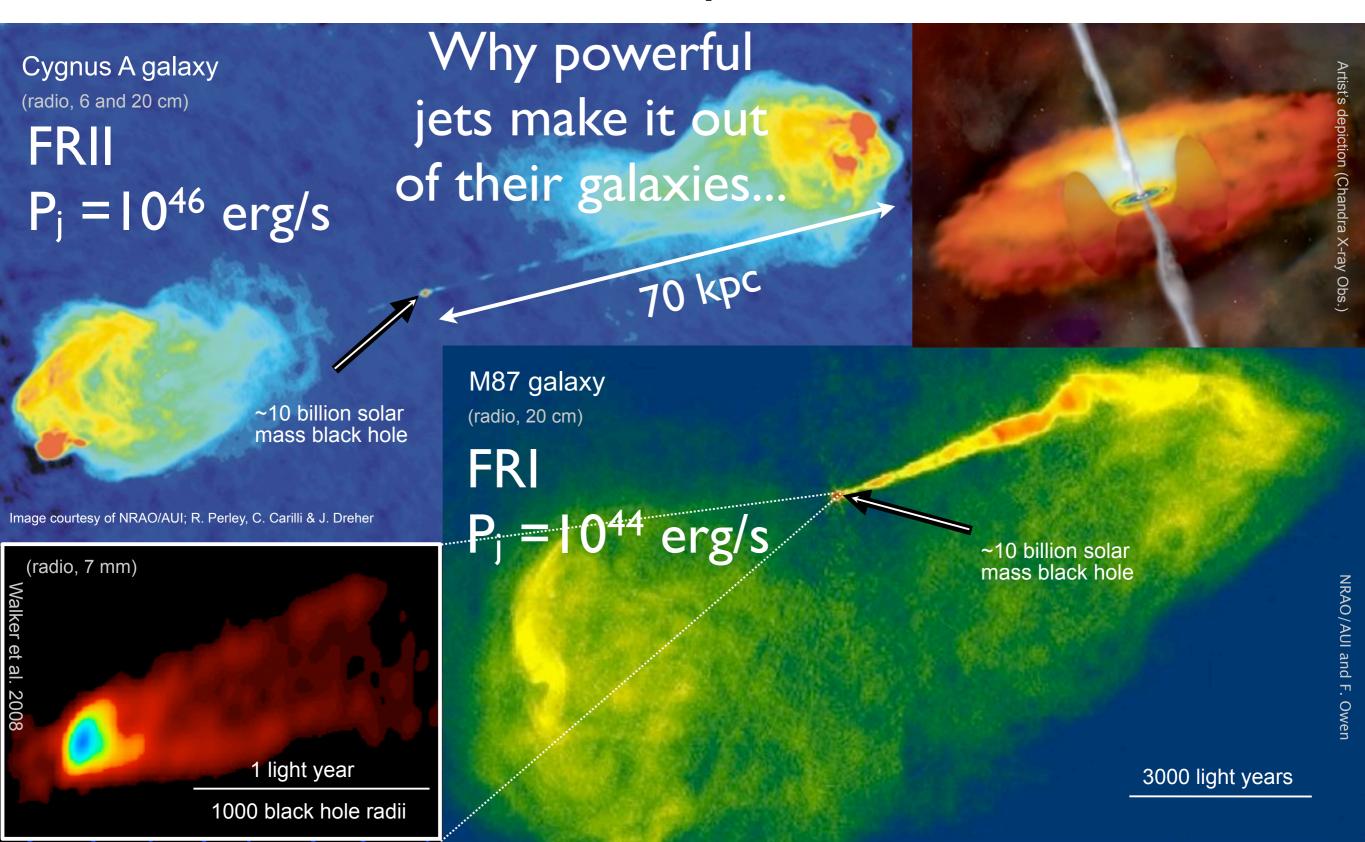
# How does Jet Heating Work?

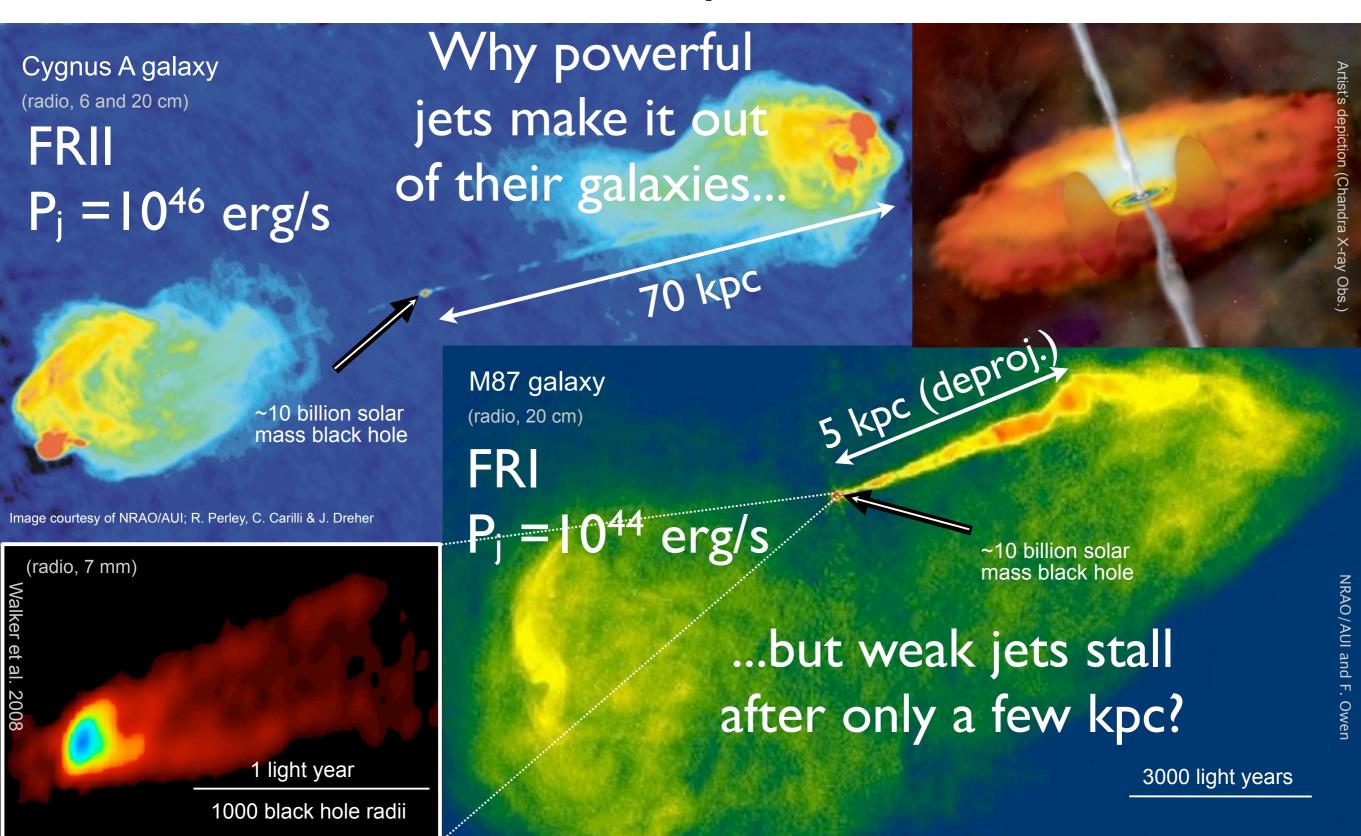


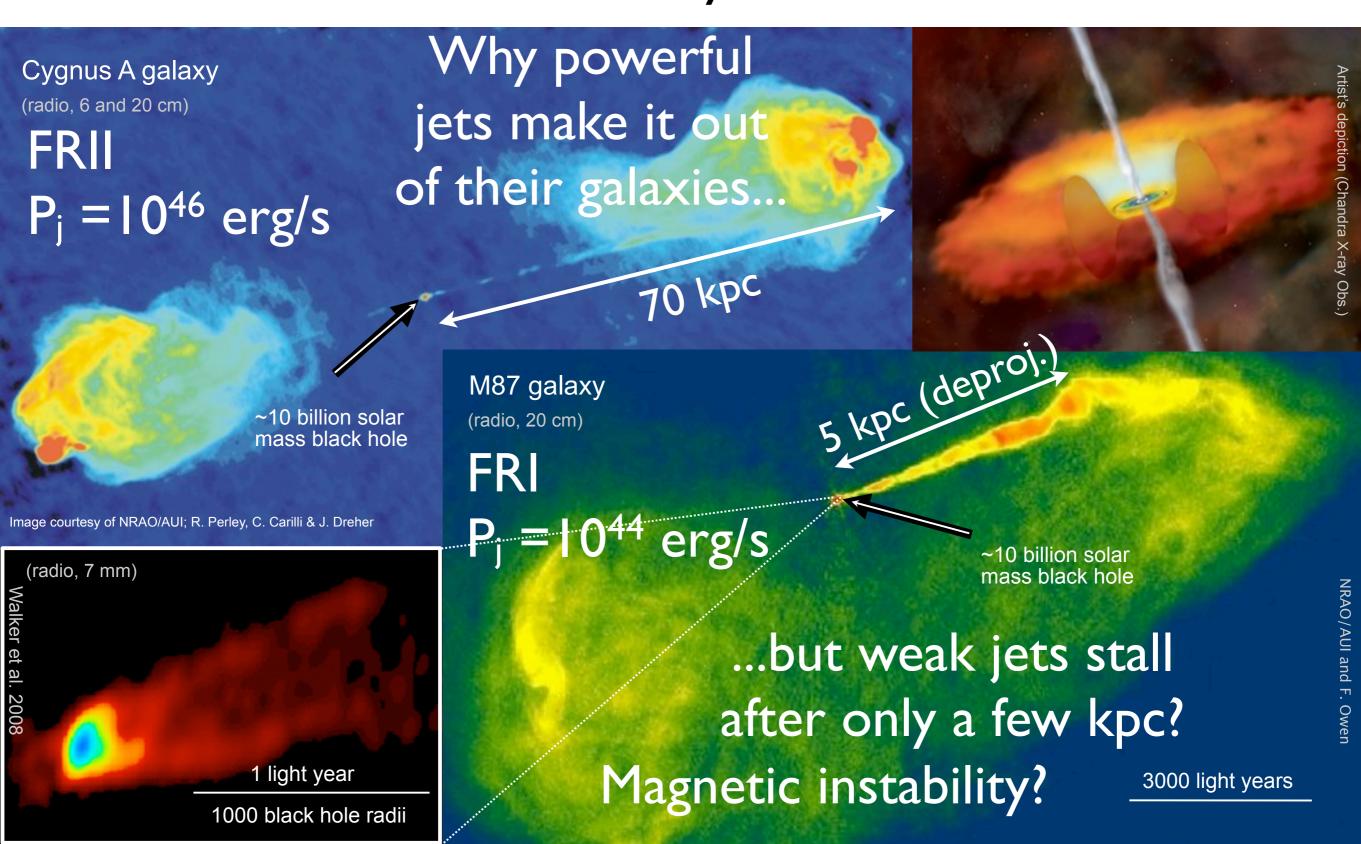
Recollimation → internal kink →

→ turbulence → reconnection → emission









# Instability of Magnetized Jets

 Kink instability growth timescale controlled by the magnetic pitch (high-mag., mildly relativistic):

$$t_{\rm kink} \simeq \frac{2\pi R_{\rm j}}{c} \frac{B_p}{B_\phi}$$

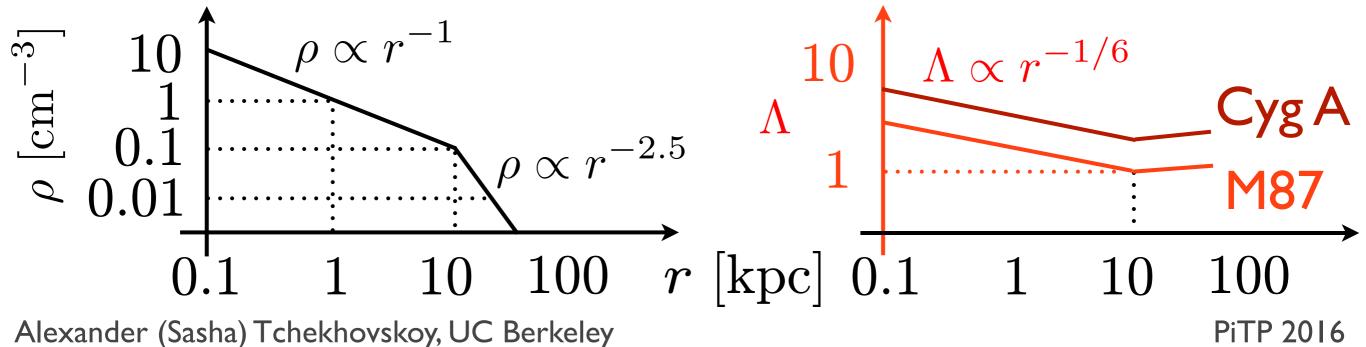
(Appl et al. 2001)

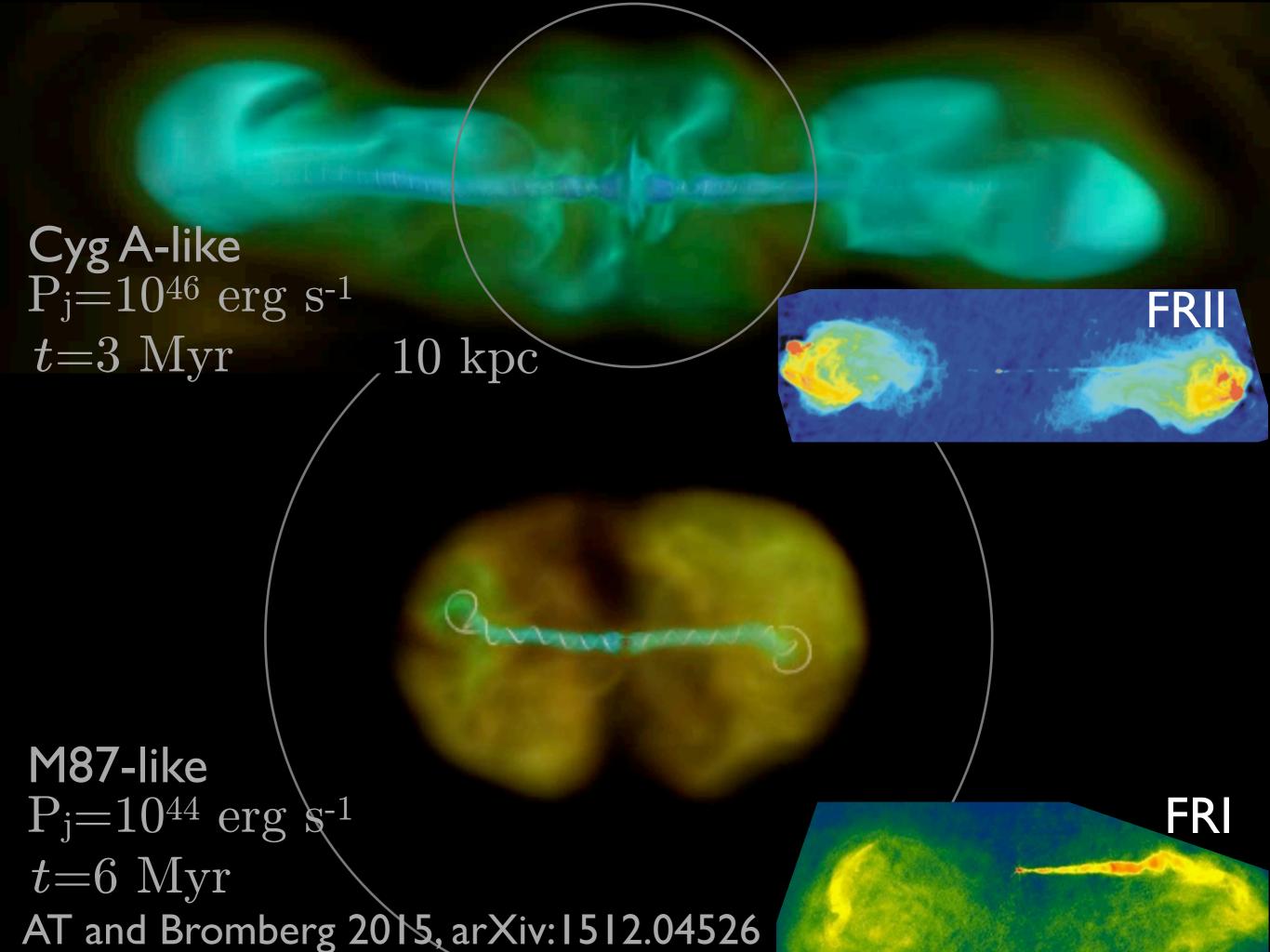
• Jets are unstable if  $5t_{\rm kink} \lesssim t_{\rm travel}$ , or

$$\Lambda \simeq 10 \left(\frac{L_{\rm j}}{\rho r^2 c^3}\right)^{1/6} \lesssim 1$$

(Bromberg & AT 2016)

Cartoon galaxy density profile:





## Summary

- Dynamically important magnetic fields wide spread:
  - Jets are robust and happy to feed on anything
  - lacktriangle MADs give the upper envelope of jet power for given M
  - BUT: strong jets benefit from disk rotation
  - Problem: How do you get rid of large-scale magnetic flux?
- Large-scale poloidal field dynamo is now a reality
  - No need for large-scale poloidal flux: toroidal would do too
- Jet morphology is set by 3D external kink and controlled by jet power and ambient density:
  - low-power jets are unstable and get stalled inside galaxies
  - FRI/FRII dichotomy likely mediated by magnetic instabilities