# Signals and Backgrounds for the LHC SUSY

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# **SUSY** in hadron colliders

Inclusive signatures provide evidence up to  $2.5~{\rm TeV}$  for squarks and gluinos.

Everything is produced at once; squarks and gluinos have largest rates.

Production of Sparticles with only E-W couplings (e.g sleptons, Higgs) may be dominated by decays not direct production.

Must use a consistent model for simulation

cannot discuss one sparticle in isolation.

Makes studies somewhat complicated and general conclusions difficult to draw.

LHC Strategies different from Tevatron where weak gaugino production probably dominates

Studies shown here are not optimized

Large event rates are used to cut hard to get rid of standard model background.

Dominant backgrounds are combinatorial from SUSY events themselves.

Studies shown here are not optimized; large event rates are exploited to cut hard to get rid of standard model background.

Full program difficult to estimate, depends on masses and branching ratios



# **General remarks**

Huge number of theoretical models

Most general SUSY model has > 100 parameters Simulation has concentrated on cases that are qualitatively different; some examples were chosen in the expectation that they would be hard.

Model determines the masses, decays and signals. A Model must be used for simulation in order to understand the problems of reconstruction – Background to SUSY is SUSY itself



# **Characteristic SUSY signatures at LHC**

Not all present in all models

 $E_T$ High Multiplicity of large  $p_t$  jets Many isolated leptons Copious  $\boldsymbol{b}$  production Large Higgs production Isolated Photons Quasi-stable charged particles **N.B.**Production of heavy objects implies subset these signals Important for triggering considerations I show only examples, many cases have been studied (SUGRA, GMSB, broken R parity ....) Simplest models have few parameters;  $m_{1/2}$  and  $m_0$  determine gluino and slepton masses.



#### SUGRA Model

Grandaddy of SUSY models

Unification all scalar masses  $(m_0)$  at GUT scale

Unification all gaugino masses  $(m_{1/2})$  at GUT scale

Three more parameters  $aneta=v_1/v_2\;sign(\mu)$  (superpotential has  $\mu H_1H_2$ ) and

Trilinear term A, important only for  $3^{rd}$  generation

Full mass spectrum and decay table predicted

Gluino mass strongly correlates with  $m_{1/2}$ , slepton mass with  $m_0$ .

R parity good – neutral LSP stable – all events have 2 LSP's in them  $\Rightarrow$  missing  $E_T$ 

Gravitino has mass in TeV region: irrelevant to colliders

Can relax unification assumption – more parameters





Contours of fixed *gluino* and squark mass





Contours of fixed wino and slepton mass



#### Where to look



 $m_0$ 



## Nomenclature

- Bulk region: Masses low  $\widetilde{\chi}_1^0 \widetilde{\chi}_1^0 
  ightarrow e^+ e^-$
- $\bullet$  Connahilation region: Something almost degenerate with LSP:  $\widetilde{\chi}^0_1 \tilde{e_R} \to e \gamma$  enhanced
- Higgs pole region:  $M_A \sim 2 m_{\widetilde{\chi}^0_1}$
- Focus point region: Large  $m_0$ , LSP Higgsino, region is sensitive to  $m_t op$ , may not exist if  $m_{top} = 172$ . phenomenology of split SUSY is very similar to this region

Caveat emptor: Don't take this too seriously



#### Two typical spectra





Several SUGRA cases studied in detail In one case unification assumptions were relaxed to investigate how signals changed (New signals appeared, old ones stayed) Some cases were restudied assuming that R-Parity was broken  $\Rightarrow$  LSP decayed inside detector.

Note typically large rates

Point	$m_0$	$m_{1/2}$	$A_0$	aneta	$\mathrm{sgn} oldsymbol{\mu}$	$\sigma$
	(GeV)	(GeV)	(GeV)			(pb)
1	400	400	0	2.0	+	2.9
2	400	400	0	10.0	+	2.9
3	200	100	0	2.0	_	1300
4	800	200	0	10.0	+	28
5	100	300	300	2.1	+	15
6	200	200	0	45	—	99

Table 1: SUGRA parameters for the six LHC points.



Point	1	2	3	4	5	6
~						
$\widetilde{\boldsymbol{g}}_{_{  }}$	1004	1009	298	582	767	540
$\widetilde{\chi}_{1}^{\pm}$	325	321	96	147	232	152
$\widetilde{\chi}^{\pm}_2$	764	537	272	315	518	307
$\widetilde{\chi}_1^0$	168	168	45	80	122	81
$\widetilde{\chi}^{0}_{2}$	326	321	97	148	233	152
$\widetilde{\chi}_3^0$	750	519	257	290	497	286
$\widetilde{\chi}_{4}^{0}$	766	538	273	315	521	304
$\widetilde{u}_L$	957	963	317	918	687	511
$\widetilde{u}_R^-$	925	933	313	910	664	498
$\widetilde{d}_L$	959	966	323	921	690	517
$\widetilde{d}_{R}$	921	930	314	910	662	498
$\widetilde{t_1}$	643	710	264	594	489	365
$\widetilde{t}_2$	924	933	329	805	717	517
$\widetilde{b}_1$	854	871	278	774	633	390
$\widetilde{b}_2$	922	930	314	903	663	480
$\widetilde{e}_L$	490	491	216	814	239	250
$\widetilde{e}_{oldsymbol{R}}$	430	431	207	805	157	219
$\widetilde{ u}_{m{e}}$	486	485	207	810	230	237
$\widetilde{ au}_1$	430	425	206	797	157	132
$\widetilde{ au}_2$	490	491	216	811	239	259
$\widetilde{\nu}_{\tau}$	486	483	207	806	230	218
$h^0_{}$	111	125	68	117	104	112
$H^0_{0}$	1046	737	379	858	638	157
$A^0$	1044	737	371	859	634	157
$H^{\pm}$	1046	741	378	862	638	182



#### **General Features**

In general  $m_{squark} > m_{slepton}$ ,  $m_{gluino} > m_{\widetilde{W}}$ Splitting between  $m_{\widetilde{e}_l}$  and  $m_{\widetilde{e}_r}$ Stop is usually lightest squark Lightest SUSY particle (LSP) stable if R-parity good. LSP must be neutral if stable SUSY particles produced in pairs even if R-parity broken. SUSY production is dominated by gluinos and squarks. Not necessarily true for Tevatron. Stable LSP  $\Rightarrow$  Missing  $E_T$ Background for SUSY usually other SUSY, not Standard Model.



#### Gauge Mediated Model

Aims to solve FCNC problem by using gauge interactions instead of Gravity to transmit SUSY breaking

Messenger Sector consists of some particles (X) that have SM interactions and are aware of SUSY breaking.

 $M_i^2 = M^2 \pm F_A$  Simplest X is complete SU(5) **5** or **10** to preserve GUT

Fundamental SUSY breaking scale  $F > F_A$ , but  $\sqrt{F} \leq 10^{10}$  GeV or SUGRA breaking will dominate Gaugino masses at 1-loop

 $M_{\widetilde{g}} \sim lpha_s N_X \Lambda$ 

Squark and Slepton masses at 2-loop

$$M_{\widetilde{e}}\sim lpha_W\sqrt{N_X}\Lambda$$

True LSP is a (almost) massless Gravitino Sparticles decay as in SUGRA, then "NLSP" decays to  $\tilde{G}$  lifetime model dependent NLSP does not have to be neutral



#### 6 parameters

 $\Lambda$ , M,  $N_5$ , aneta,

#### $sign\mu$

 $10 \,\mathrm{TeV} \lesssim \Lambda \equiv F_A/M \lesssim 400 \,\mathrm{TeV}$ : Scale for SUSY masses.

 $M > \Lambda$ : Messenger mass scale.

 $N_5 \geq 1$ : Number of equivalent  $5+ar{5}$  messenger fields.

 $1 \lesssim aneta \lesssim m_t/m_b$ : Usual ratio of Higgs VEV's.

 $\operatorname{sgn} \mu = \pm 1$ : Usual sign of  $\mu$  parameter.

 $C_{ ext{grav}} \geq 1$ : Ratio of  $M_{\widetilde{G}}$  to value from  $F_A$ , controls lifetime of NLSP.

Point	Λ	$M_{m}$	$N_5$	aneta	$\mathrm{sgn} oldsymbol{\mu}$	$C_{ m grav} \geq 1$	$\sigma$
	(TeV)	(TeV)					(pb)
G1a	90	500	1	5.0	+	1.0	7.6
G1b	90	500	1	5.0	+	$10^3$	7.6
G2a	30	250	3	5.0	+	1.0	23
G2b	30	250	3	5.0	+	$5 imes 10^3$	23



Sparticle	G1	G2	Sparticle	G1	G2
$\widetilde{\boldsymbol{g}}_{_{1}}$	747	713			
$\widetilde{\chi}_1^{\pm}$	223	201	$\widetilde{\chi}_2^{\pm}$	469	346
$\widetilde{\chi}_1^{0}$	119	116	$\widetilde{\chi}_2^0$	224	204
$\widetilde{\chi}_3^{ar{0}}$	451	305	$\widetilde{\chi}_4^{m 0}$	470	348
$\widetilde{u}_L^{"}$	986	672	$\widetilde{u}_{R}^{+}$	942	649
$\widetilde{d}_L$	989	676	$\widetilde{d}_{R}$	939	648
$\widetilde{t_1}$	846	584	$\widetilde{t_2}$	962	684
$\widetilde{b}_1$	935	643	$\widetilde{b}_2$	945	652
$\widetilde{e}_{L}^{-}$	326	204	$\widetilde{e}_{R}$	164	103
$\widetilde{\nu_e}$	317	189	$ ilde{ au}_2^{ ilde{ au}}$	326	204
$\widetilde{oldsymbol{ au}}_1$	163	102	$\widetilde{ u}_{oldsymbol{ au}}^{-}$	316	189
$h^{\overline{0}}$	110	107	$H^0$	557	360
$A^0$	555	358	$H^{\pm}$	562	367

Mass spectrum more spread out than in SUGRA m(squark)/m(slepton) bigger



#### Anomaly mediated model

Superconformal anomaly always present

predicts sparticle masses in terms of  $m_{3/2}$ Simplest version predicts tachyonic sleptons!

Randall, Sundrum, Luty, Giudice, Wells, Murayama

Some other SUSY breaking mechanism must be present to get realistic spectrum Add universal squark masses (mAMSB) or new very heavy fields (DAMSB) (similar to gauge mediated), both variants are in ISAJET.

AMSB only – Most important feature  $M_3 > M_1 > M_2 \Rightarrow$  LSP is a  $\tilde{W^0}$  and almost degenerate with  $\tilde{\chi_1^+} \chi_1^+ \rightarrow \tilde{\chi}_1^0 \pi^+$  with  $c\tau < 10$  cm DAMSB has very short lifetime and bigger mass difference wells, Paige

Sleptons are lighter than squarks  $\tilde{q_r} \to \tilde{\chi}_2^0 q$  and  $\tilde{q_l} \to \tilde{\chi}_1^0 q$ , *i.e.* opposite to SUGRA and GMSB.

Gravitino mass is  $\sim$  TeV, irrelevant to LHC.



AMSB has 4 parameters  $m_0$ ,  $m_{3/2}$ , aneta,  $sign\mu$ 

DAMSB has 5 parameters  $M_0$ ,  $m_{3/2}$ ,  $n \tan \beta$ ,  $sign \mu$ . n is the number of new fields at mass M.

 $m_{3/2}$  is the gravitino mass

Sparticle	AMSB	DAMSB	Sparticle	AMSB	DAMSB
$\widetilde{\boldsymbol{g}}_{_{ }}$	815	500			
$\widetilde{\chi}_1^{\pm}$	101	145	$\widetilde{\chi}_2^{\pm}$	658	481
$\widetilde{\chi}_1^{\overline{0}}$	101	136	$\widetilde{\chi}_2^{\overline{0}}$	322	152
$\widetilde{\chi}_3^0$	652	462	$\widetilde{\chi}_4^0$	657	483
$\widetilde{u}_L^{\circ}$	754	432	$\widetilde{u}_{R}^{+}$	758	384
$\widetilde{d}_L$	757	439	$\widetilde{d}_R$	763	371
$\widetilde{t}_1$	516	306	$\widetilde{t}_2$	745	454
$\widetilde{b}_1$	670	371	$\widetilde{b}_2$	763	406
$\widetilde{e}_L$	155	257	$\widetilde{e}_{oldsymbol{R}}$	153	190
$\widetilde{oldsymbol{ u}}_{oldsymbol{e}}$	137	246	$\widetilde{oldsymbol{ au}}_2$	166	257
$\widetilde{oldsymbol{ au}}_1$	140	190	$\widetilde{oldsymbol{ u}}_{oldsymbol{ au}}$	137	246
$h^0$	107	98	$H^0_{\perp}$	699	297
$A^0$	697	293	$H^{\pm}$	701	303



# Start with SUGRA

- Look for characteristic signals
- jets and missing energy
- these should work anywhere



#### **Inclusive** analysis



Will determine gluino/squark masses to  $\sim 15\%$ 







#### Can also add lepton(s): More channels more robustness



mSUGRA reach in various final states for 100 fb<sup>-1</sup>



Energy of LHC is most crucial: reach increases slowly with luminosity







#### How robust is this?

Backgrounds based on showering MC may underestimate multi-jet final states.



May lower reach slightly Plot shows all jet state Signal in Lepton+jets+missin more robust

#### Asai et al



Peak in  $M_{
m eff}$  distribution correlates well with SUSY mass scale



 $M_{\mathrm{SUSY}} = \min(M_{\widetilde{u}}, M_{\widetilde{g}})$ 

Use this and similar global distributions to establish that new physics exists and determine its mass scale Method is slightly model dependent



## **Generalizations to other models**

Similar method works in GMSB and MSSM In MSSM, 15 parameters were varied Events selected to have no isolated leptons, at least 4 jets, large missing  $E_T$ More global variables were used; best is

Error is bigger in MSSM





Model	Var	$ar{m{x}}$	$\sigma$	$\sigma/ar{x}$	Prec. (%)
mSUGRA	1	1.585	0.049	0.031	2.9
	2	0.991	0.039	0.039	3.8
	3	1.700	0.043	0.026	2.1
	4	1.089	0.030	0.028	2.5
	5	1.168	0.029	0.025	2.1
MSSM	1	1.657	0.386	0.233	23.1
	2	0.998	0.214	0.215	21.1
	3	1.722	0.227	0.132	12.8
	4	1.092	0.143	0.131	12.8
	5	1.156	0.176	0.152	14.8
GMSB	1	1.660	0.149	0.090	8.1
	2	1.095	0.085	0.077	6.6
	3	1.832	0.176	0.096	9.0
	4	1.235	0.091	0.074	6.1
	5	1.273	0.109	0.086	7.9

 $\sigma(M_{susy} < 13\%)$ 

Not optimized

Leptonic channels not used

More work on "global signatures" needed



What about  $1 fb^{-1}$ 



Now for examples of specific final states...



# **Digression on Simulations**

Varying degrees of sophistication

Run an event generator: Makes a list of particles (e.g Pythia)

#### • "Full simulation"

Full material description of detector
Track particle: model its interactions, follow all the secondaries with energy above some threshold
Translate "hits" into simulated electronic signals
Now event looks like a real event: reconstruct it using same software
Advantage: Full description
Disadvantage: Slow (~15 mins per event), can only be done by experimenter.
Problems, Low energy hadronic interactions, geometry is hard

#### • "Fast simulation (theorists version)"

Assume perfect detector: Apply jet finding algorithm Smear, electron, muon, jet momenta and missing ET. using some resolution function



Advantage. Very fast (comparable to generation time), can by done by theorist Disadvantage: Only as good as resolution function: problems in "tails"

• "Fast simulation (experimenters version)"

Response of individual particles is parametrized. Might use "full" for some particles (*e.g.* Photons) and parametrized for others (pions) May use full reconstruction.

In practice all are used: Theorists version is often good enough for evaluating a model



# **Characteristic SUSY Decays**

Illustrate techniques by choosing examples from case studies.

Both  $\widetilde{q}$  and  $\widetilde{g}$  produced; one decays to the other

Weak gauginos (  $\widetilde{\chi_i^0}, \widetilde{\chi_i^\pm}$  ) then produced in their decay.  $e.g. \ \widetilde{q_L} \to \widetilde{\chi}_2^0 q_L$ 

Two generic features  $\chi_2^0 \rightarrow \chi_1^0 h$  or  $\chi_2^0 \rightarrow \chi_1^0 \ell^+ \ell^-$  possibly via intermediate slepton  $\chi_2^0 \rightarrow \widetilde{\ell^+} \ell^- \rightarrow \chi_1^0 \ell^+ \ell^-$ Former tends to dominate if kinematically allowed.

Use these characteristic decays as a starting point for mass measurements

Many SUSY particles can then be identified by adding more jets/leptons



# **Decays to Higgs bosons**

If  $\chi^0_2 \to \chi^0_1 h$  exists then this final state followed by  $h \to b\overline{b}$  results in discovery of Higgs at LHC.

In these cases  $\sim 20\%$  of SUSY events contain h 
ightarrow bb





# **Generally applicable**



Over rest of parameter space, leptons are the key...



#### Starting with Leptons

Isolated leptons indicate presence of t, W, Z, weak gauginos or sleptons Key decays are  $\tilde{\chi}_2 \rightarrow \tilde{\ell}^+ \ell^-$  and  $\tilde{\chi}_2 \rightarrow \tilde{\chi}_1 \ell^+ \ell^-$ Mass of opposite sign same flavor leptons is constrained by decay



 $\begin{array}{ll} & {}^{\rm M_{I\!I}(GeV)} \\ \mbox{Decay via real slepton: } \widetilde{\chi}_2 \rightarrow \widetilde{\ell}^+ \ell^- & \mbox{Decay} \\ \mbox{Plot shows } e^+ e^- + \mu^+ \mu^- - e^\pm \mu^\mp & \mbox{and } Z \end{array}$ 



Decay via virtual slepton:  $\widetilde{\chi}_2 \rightarrow \widetilde{\chi}_1 \ell^+ \ell^$ and Z from other SUSY particles



# **Building on Leptons**

Decay  $ilde q_L o q \widetilde \chi^0_2 o q \widetilde \ell \ell o q \ell \ell \widetilde \chi^0_1$ 

Identify and measure decay chain

- ullet 2 isolated opposite sign leptons;  $p_t > 10~{
  m GeV}$
- $\bullet \geq 4$  jets; one has  $p_t > 100~GeV$ , rest  $p_t > 50~{
  m GeV}$
- $E_T > max(100, 0.2M_{eff})$



Mass of  $q\ell\ell$  system has max at

$$M_{\ell\ell q}^{\max} = [rac{(M_{\widetilde{q}_L}^2 - M_{\widetilde{\chi}_2^0}^2)(M_{\widetilde{\chi}_2^0}^2 - M_{\widetilde{\chi}_1^0}^2)}{M_{\widetilde{\chi}_2^0}^2}]^{1/2} = 552.4\,{
m GeV}$$

and min at 271  ${\rm GeV}$ 





smallest mass of possible  $\ell\ell jet$  combinations

Kinematic structure clearly seen Can also exploit  $\ell jet$  mass




Can now solve for the masses. Note that no model is needed

Very naive analysis has 4 constraints from  $lq, llq_{upper}, llq_{lower}, ll$  masses 4 Unknowns,  $m_{\tilde{q_L}}, m_{\tilde{e_R}}, m_{\tilde{\chi}^0_2}, m_{\tilde{\chi}^0_1}$ 







# **Right squarks**

#### s-tranverse mass. Definition

- Select Events:
  - $pp \to X + \tilde{q}_R \tilde{q}_R \to X + q \bar{q} \tilde{\chi}_1^0 \tilde{\chi}_1^0$ 
    - 2 Jets with  $P_T$  >200GeV
    - ∆(j1-j2)>1
    - Missing  $E_T$  >400GeV



• Partition  $\vec{E}_T = \vec{E}_{T,1} + \vec{E}_{T,2}$  in all possible ways and compute:

$$M_T^2 = \min_{\vec{k}_{T,1}, \vec{k}_{T,2}} \left[ \max\{m_T^2(P_{T,j1}, \vec{k}_{T,1}, M_{\tilde{\chi}_1^0}), m_T^2(P_{T,j2}, \vec{k}_{T,2}, M_{\tilde{\chi}_1^0}) \} \right]$$

 $\bullet~M_T^2$  depends on the choice of  $M(\tilde{\chi}_1^0)$ 



# s-tranverse mass for the SU5 point

Fit the end-point with a straight line and extrapolate to the *x*-axis Use "true" value of  $M(\tilde{\chi}_1^0)$ 

#### (Ola Kristoff Oye)



Again Limited by LSP mass uncertainty



#### Final states with taus

Large an eta implies that  $m( ilde{ au}) < m( ilde{\mu})$ Taus may be the only produced leptons in gaugino decay. Leptonic tau decays are of limited use – where did lepton come from? Use Hadronic tau decays, using jet shape and multiplicity for ID and jet rejection. Full simulation study used to estimate efficiency and rejection Rely on Jet and  $E_t(miss)$  cuts to get rid of SM background Measure "visible" tau energy Event selection  $\geq 4$  jets, one has  $p_t > 100$  GeV, rest  $p_t > 50$  GeV No isolated leptons with  $p_t > 10 \text{ GeV}$  $E_T > max(100, 0.2M_{eff})$ 







Real signal visible above fakes (dashed) and SM (solid)

Can use peak position to infer end point in decay  $\tilde{\chi}_2^0 \rightarrow \tau \tau \tilde{\chi}_1^0$  (61 GeV) Estimate 5% error

Large  $\tan \beta \Rightarrow$  light sbottom – Look for these



 $ilde{g} 
ightarrow b ilde{b} 
ightarrow b b au^\pm au^\mp \widetilde{\chi}_1^0$ 

Previous sample with 2  $b-{
m jets}$  having  $p_t>25$  GeV

Lots of missing  $E_T$ : tau decays and  $\widetilde{\chi}_1^0$ 's

Select  $40 < m_{ au au} < 60$  GeV

Combine with b jets

Look at au au bb and au au: should approximate gluino and sbottom use partial reconstruction technique assuming mass of  $\widetilde{\chi}^0_1$ 

Peaks are low; should be expected due to missing energy



$$m(\widetilde{\chi}^0_2 bb)$$
 vs  $m(\widetilde{\chi}^0_2 bb)$  –



#### Projections







Denegri, Majerotto, Rurua



## Explicit flavor violation is also possible

Neutrino oscillations imply lepton number is violated Atmospheric muon neutrino deficit implies  $\nu_{mu} \leftrightarrow \nu_{\tau}$  with maximal mixing In a SUSY model, expect significant flavor violation in slepton sector Simplest model of lepton number violation involves addition of right handed neutrino N with SUSY conserving Majorana mass mNN and coupling to lepton left doublet and Higgs of the form LNHIncluding only  $\mu \leftrightarrow \tau$  mixing gives

$$M_{\widetilde{\ell}\widetilde{\ell}}^2 = \left[ egin{array}{cccccccc} M_L^2 + D_L & 0 & 0 & 0 & 0 & 0 \ 0 & M_L^2 + D_L & M_{\mu au}^2 & 0 & 0 & 0 \ 0 & M_{\mu au}^2 & M_{ au_L}^2 + D_L & 0 & 0 & m_ auar{A}_ au \ 0 & 0 & 0 & M_R^2 + D_R & 0 & 0 \ 0 & 0 & 0 & 0 & M_R^2 + D_R & 0 \ 0 & 0 & m_ auar{A}_ au & 0 & 0 & M_R^2 + D_R \end{array} 
ight]$$

Atmospheric neutrinos suggest maximal mixing *i.e.*  $\delta = O(1)$ 



Two types of flavor violation production (  $\widetilde{\chi}^0_2 \to \tilde{\tau} \mu$ ) and decay ( $\tilde{\tau} \to \widetilde{\chi}^0_1 \mu$ ).



#### Signal for lepton number violation comes by comparing $\mu au_h$ and $e au_h$ final states





 $\ell^{\pm} \tau_{h}^{\mp}$  signal (red),  $\ell^{\pm} \tau_{h}^{\pm}$  signal (blue),  $\mu^{\pm} \tau_{h}^{\mp}$  from LFV decays with BR = 10%(magenta), and Standard Model  $\ell^{\pm} \tau_{h}^{\mp}$ 

Lepton number violating decay  $\widetilde{\chi}_2^0 \to \mu \tau_h \widetilde{\chi}_1^0$  give harder  $\mu \tau$  mass distribution than that from  $\widetilde{\chi}_2^0 \to \tau \tau \widetilde{\chi}_1^0 \to \mu \tau_h \widetilde{\chi}_1^0$ 



#### Subtraction removes background



$$\ell^{\pm}\tau_{h}^{\mp} - \ell^{\pm}\tau_{h}^{\pm}$$
 (red ) and  $\mu^{\pm}\tau_{h}^{\mp}$  from  
LFV decays with  $BR = 10\%$  (magenta)  
Signal is established from  $E =$   
 $N(\mu^{\pm}\tau_{h}^{\mp}) - N(e^{\pm}\tau_{h}^{\mp})$   
10 fb<sup>-1</sup> and  $5\sigma$  implies BR=2.3% or  $\delta \sim$   
0.1 well within value needed for neutrino  
data

Sensitive provided that  $\widetilde{\chi}^0_2$  production is large enough (large fraction of parameter space More sensitive than  $\mu \to e\gamma$ 



#### **R-parity broken**

Implies either Lepton number or Baryon number is violated and LSP decays Either  $\tilde{\chi}_1^0 \rightarrow qqq$ , or  $\tilde{\chi}_1^0 \rightarrow q\overline{q}\ell$  or  $\tilde{\chi}_1^0 \rightarrow \ell^+\ell^-\nu$ First two have no  $\not\!\!\!E_T$ , last 2 have more leptons and are straightforward First case is hardest, Global S/B is worse due to less  $\not\!\!\!E_T$  Example, SUGRA with  $\tilde{\chi}_1^0 \rightarrow qqq$  Leptons are essential to get rid of QCD background  $\geq 8$  jets with  $p_t > 50$  GeV 2 OSSF isolated leptons.  $S_T > 0.2$ , selects "ball like" events  $\Sigma_{jets+leptons}E_T > 1$  TeV





As nothing is lost, should be possible to reconstruct  $\widetilde{\chi}_1^0$ 

Difficult because jet multiplicity is very high and  $\widetilde{\chi}^0_1$  mass is usually small, so jets are soft





Nominal mass 122 GeV



#### Can cut around peak and combine with either leptons or quarks



Note that tight cuts imply low event rate (analysis not optimized)



# New signals in GMSB

Lightest superpartner is unstable and decays to Gravitino  $( ilde{G})$ Either neutral

 $\chi_1^0 \rightarrow \gamma \tilde{G} : c\tau \sim C^2 (100 \text{ GeV}/M_{\chi_1^0})^5 (\Lambda/180 \text{TeV})^2 (M_M/180 \text{TeV})^2 \text{mm}$  $\Rightarrow$  extra photons or similar signals to SUGRA depending on lifetime Or charged

Almost always slepton:  $ilde{e_R} 
ightarrow e\widetilde{G}$ 

No Missing  $E_T$  if c au large, events have a pair of massive stable charged particles ("G2b")

Large lepton multiplicity if c au small ("G2a").

Discovery and measurement in these cases is trivial

In case "G2b", every decay product can be measured

In case "G1a"  $\widetilde{G}$  momenta can be inferred and events fully reconstructed.



GMSB case 1a: Event selection (not optimized) Decay  $\widetilde{\chi}_2^0 \rightarrow \ell^+ \ell^- \widetilde{\chi}_1^0 \rightarrow \ell^+ \ell^- \gamma \tilde{G}$  is key Lifetime of  $\widetilde{\chi}_1^0$  is short

Find jets

$$M_{\rm eff} \equiv E_T + p_{T,1} + p_{T,2} + p_{T,3} + p_{T,4}$$
.

Require

$$\widetilde{\chi}^0_2 o \widetilde{\ell}^\pm \ell^\mp o \widetilde{\chi}^0_1 \ell^\pm \ell^\mp o \widetilde{G} \gamma \ell^\pm \ell^\mp \,,$$

Electrons and photons :  $p_T > 20 \, {
m GeV}$ Muons :  $p_T > 5 \, {
m GeV}$ .

Require at least 2 photons and two leptons.



Dilepton mass distribution, flavor subtracted  $e^+e^-+\mu^+\mu^--e^\pm\mu^\mp$ 



End is at

$$M_{\widetilde{\chi}^0_2}\sqrt{1-\left(rac{M_{\widetilde{\ell}_R}}{M_{\widetilde{\chi}^0_2}}
ight)^2}\sqrt{1-\left(rac{M_{\widetilde{\chi}^0_1}}{M_{\widetilde{\ell}_R}}
ight)^2}=105.1$$





Form  $\ell^+\ell^-\gamma$  mass and take smallest combination. Linear vanishing at

$$\sqrt{M_{\widetilde{\chi}^0_2}^2 - M_{\chi^0_1}^2} = 189.7\,{
m GeV}\,,$$





Form  $\ell^{\pm}\gamma$  mass also. Two structures at

$$\sqrt{M_{ ilde{\ell}_R}^2 - M_{\chi_1^0}^2} = 112.7\,{
m GeV}$$

 $\quad \text{and} \quad$ 

$$\sqrt{M_{\chi^0_2}^2 - M_{ ilde{\ell}_R}^2} = 152.6\,{
m GeV}$$



These four measurements are sufficient to determine the masses of the particles  $(\tilde{\chi}_2^0, \tilde{\ell}_R, \text{ and } \tilde{\chi}_1^0)$  in this decay chain without assuming any model of SUSY breaking.

Now use this to reconstruct the decay chain and measure the  $\widetilde{G}$  momenta despite the fact that there are two in each event and both are invisible!



Full reconstruction of SUSY events



Know masses  $\Rightarrow$  can calculate p assuming  $p^2 = 0$ :

$$egin{array}{rll} 2p_0k_0-2ec p\cdotec k &=& M_{\widetilde{\chi}_1^0}^2\ 2p_0l_0-2ec p\cdotec l &=& M_{\widetilde{\ell}_R}^2-M_{\widetilde{\chi}_1^0}^2-2k\cdot l\ 2p_0k_0-2ec p\cdotec q &=& M_{\widetilde{\chi}_2^0}^2-M_{\widetilde{\ell}_R}^2-2(k+l)\cdot q \end{array}$$

0C fit with  $2 \times 2$  solutions.

Event has two of these decays so require 4 leptons and 2 gammas



Calculate missing  $E_T$ Form a  $\chi^2$  using measured missing  $E_T$  to resolve ambiguities

$$\chi^2 = \left(rac{E_x - p_{1x} - p_{2x}}{\Delta E_x}
ight)^2 + \left(rac{E_y - p_{1y} - p_{2y}}{\Delta E_y}
ight)^2 \,.$$

use  $\Delta E_x = \Delta E_x = 0.6 \sqrt{E_T} + 0.03 E_T$ .



Compare to generated  $\widetilde{G}$  momenta Plot shows all solutions with  $\chi^2 < 10$  $\Delta \vec{p} = \vec{p}_{\tilde{G}} - \vec{p}_{reconst}$  $\Delta |\vec{p}|/|\vec{p}| \sim 10\%$ 



## Squark and Gluino Masses

Use measured  $\widetilde{\chi}_2^0$  momenta and combine with jets  $\widetilde{q} \to \widetilde{g}q \to \widetilde{\chi}_2^0 \overline{q}qq$ Require at least 4 jets with  $p_T > 75 \,\mathrm{GeV}$ 



Figure shows mass of  $\tilde{\chi}_2^0+2$  jets; peak is below gluino mass (747 GeV); no correction applied for small jet cone.

Much easier than the SUGRA cases; masses measured directly



#### Measuring the fundamental scale of SUSY breaking

Lifetime of  $\widetilde{\chi}^0_1 o ilde{G}$  is important as it measures the fundamental scale of SUSY breaking Measure lifetime of  $\chi^0_1 \ ( o \widetilde{G}\gamma)$  using Dalitz decay  $\chi^0_1 o e^+ e^- \gamma \widetilde{G}$ Works for short lived  $\widetilde{\chi}_1^0$ Statistics limited ( $\sim$ few-K events) Measure lifetime of  $\chi_1^0 (\to \widetilde{G}\gamma)$ : photon pointing. Angular resolution of photons from primary vertex (ATLAS)  $\Delta \theta \sim 60 mr/\sqrt{E}$  Detailed study of efficiency for non-pointing photons Important for long lived  $\widetilde{\chi}_1^0$ Decays are uniformly distributed in the detector Cross check from time delay of decay Failure to see photons  $\Rightarrow c au > 100$  km or  $\sqrt{F} > 10^4$  TeV



#### Mass measurement of quasi-stable sleptons – ATLAS

Sleptons are produced at the end of decay chains  $\Rightarrow$  large velocity Most of these will pass the Muon Trigger Measure the velocity using TOF in Muon system, then infer mass Time resolution  $\sim 65$  ns

 $\Rightarrow \Delta M/M \sim 3\%$  for  $M=100~{
m GeV}$ 





#### How well does this work in CMS?

Three cases studied  $m_{stau} = 104, 303, 636 \; {
m GeV}$ 







#### **AMSB**

Look at paper of Wells and Paige



## Odd Ball: R-hadron (split susy)

Long lived or quasi stable gluino produces "cannon ball" that charge exchanges as it PYTHIA R-hadron event from ATLSIM



passes through detector



## Spin measurements at LHC?

Conventional wisdom says that you need LC for this but...



Angle between q and  $e^-$  in  $\tilde{\chi}_2^0$  rest frame is sensitive to spin correlations. But effect washes out if we do not know which lepton comes out first: Use kinematics But effect washes out if average over q and  $\overline{q}$ : LHC is a pp machine: more  $\tilde{q}$  than  $\tilde{\overline{q}}$ 



Form an asymmetry from invariant mass distribution of lepton and jet

$$A = rac{(l^+q) - (l^-q)}{(l^+q) + (l^-q)}$$

Green:spin correlation off Yellow: No detector ( $\times 0.6$ ) Needs at least 100 fb<sup>-1</sup>



Barr



#### **Difficult cases I: Small mass gaps**

Co-annihilation region: Near degeneracy between LSP and sleptons.Soft leptons and more messy decays.

 $\widetilde{q_L} \to q\widetilde{\chi}_2^0$  then  $\widetilde{\chi}_2^0(260) \to \widetilde{\ell}_R(153)\ell \to \ell\ell\widetilde{\chi}_1^0(136)$  and  $\widetilde{\chi}_2^0 \to \widetilde{\ell}_L(255)\ell \to \ell\ell\widetilde{\chi}_1^0(136)$ 

200 אפא / אוי ביווניא אין אוי אין אוי אין אפ<mark>א א</mark> 150 100 Leptons can still be found despite small mass gaps 50 50 100 150 200 0  $M_{\parallel}$  (GeV) **ATLAS** 



# Conclusions

• Serious thinking has started about what might be done at 10<sup>35</sup> and what machine and detector upgrades are needed.



## Discovery cannot be far off

"The train is already late" (Altarelli): Fine tuning is a problem already. We expected gauginos in the LEP range Tevatron "window" is small but low masses are more likely



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An era is about to end Low energy SUSY has provided employment for > 20 years It will be discovered or die in the next 6 years.

